

Highlights of particle physics

- Tracking in particle physics experiments
- Challenges for the Sensors
 - Speed
 - Granularity
 - Radiation hardness
- Status of the R&D
 - From a Liverpool perspective

High energy physics experiments

- Tracking is essentially measuring the path of minimum ionising particles
- Very high multiplicity
- High repetition rate (the present interaction frequency at the LHC at CERN is 40MHz)

High energy physics experiments

**The Large Hadron Collider
under the French-Swiss
border near Geneva**



HEP Experiments are big!

1232 superconducting (1.9 K) dipoles are needed to bend the beams around the 27 km circumference of the LHC.

At 7 TeV these magnets have to produce a vertical B field of 8.4 Tesla at a current of 11,700 A to bend the beam round via the Lorentz force.

The magnets have two apertures, one for each of the counter-rotating beams.

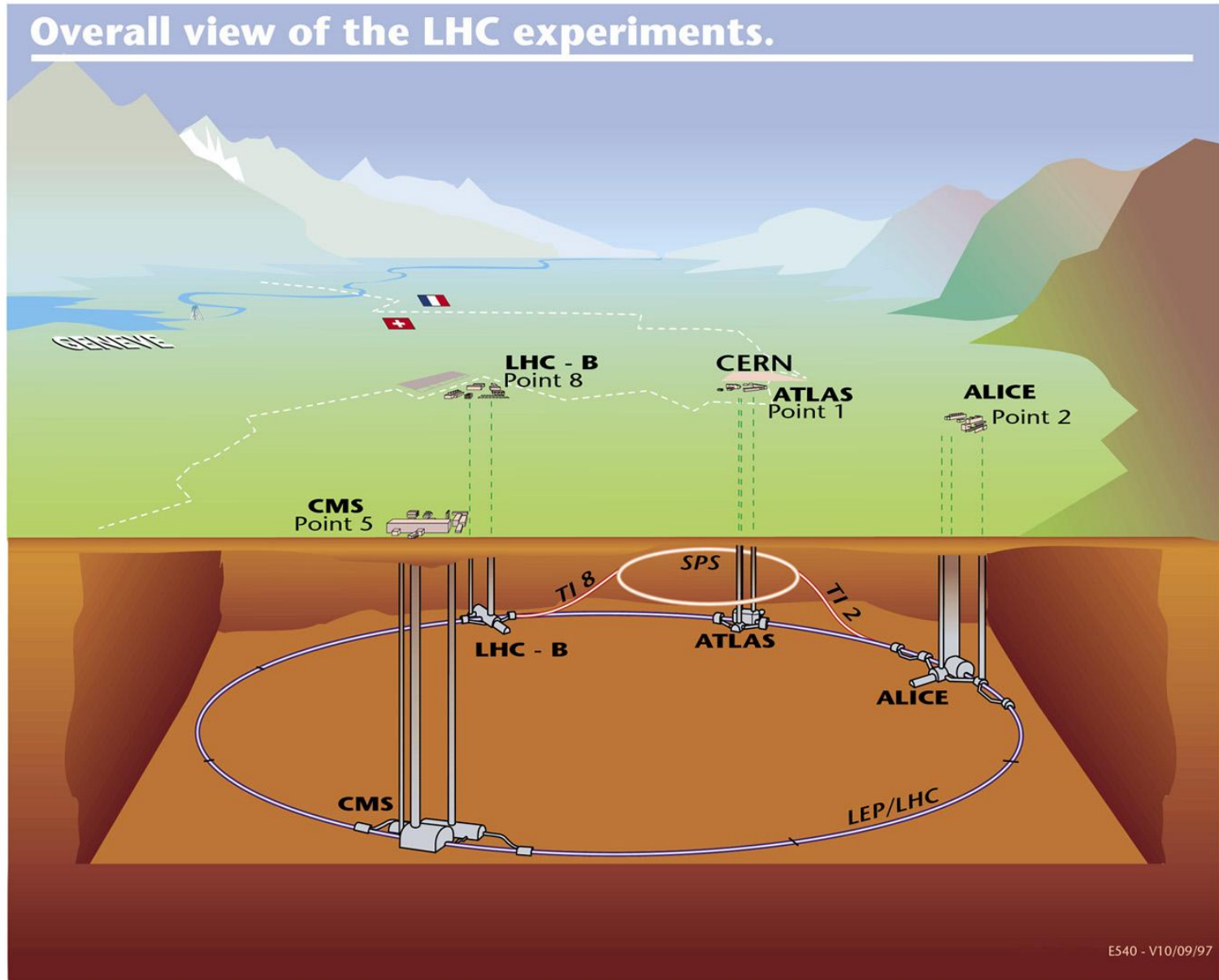
Each one is 14.3 metres long, weighs 35 tonnes and costs 0.5M€

Quads etc are also needed to keep the beam focused and the motion stable

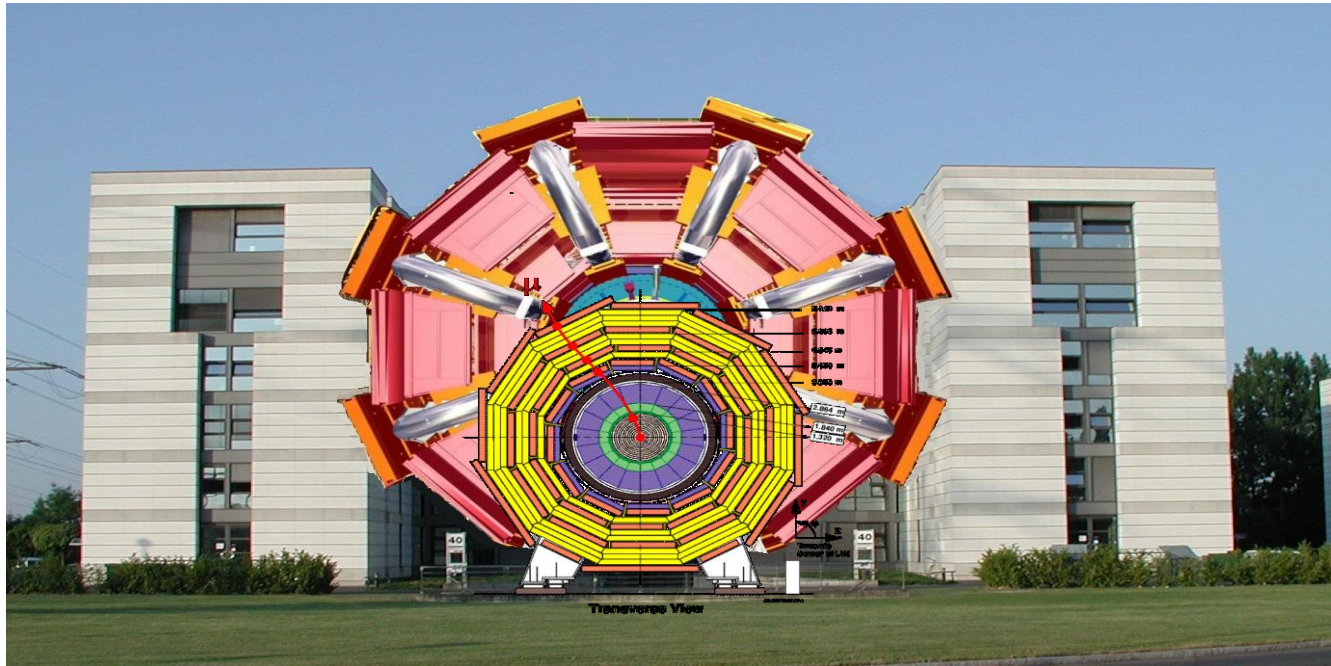
The total stored magnetic energy in the LHC is 11,000,000,000 Joules

With 2808 bunches in the LHC, the stored kinetic energy in the beam is 350,000,000 Joules

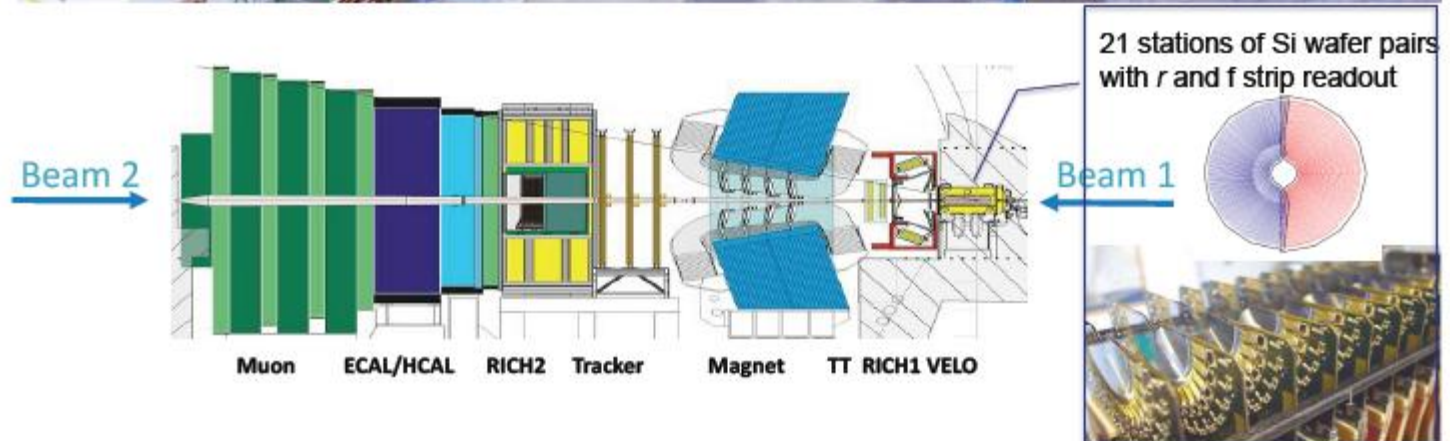
High energy physics experiments



High energy physics experiments

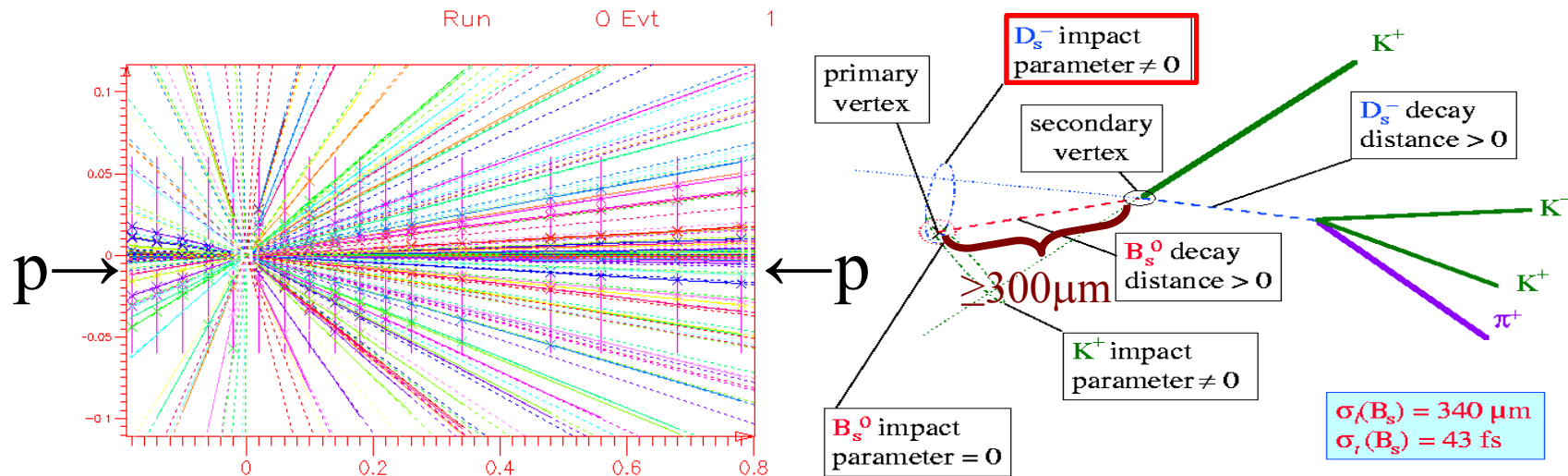


- **ATLAS, CMS, ALICE and LHCb**
- **Detector Technologies**
 - Noble gases, scintillators, crystals, Cherenkov, ...
 - **Silicon Micro-strip Tracking Detectors**



High energy physics experiments

- Nearly all early applications of silicon micro-strip detectors were to detect and measure particles with pico-second (10^{-12}) lifetimes such that (taking account of special relativity) $\beta\gamma c\tau \geq 300\mu\text{m}$
- This meant the primary goal was to locate primary (collision) and secondary vertices (as is the case in LHCb)



Side on view of 7 TeV on 7 TeV proton collision meson decay

Zoom in showing just particles from B_s^-

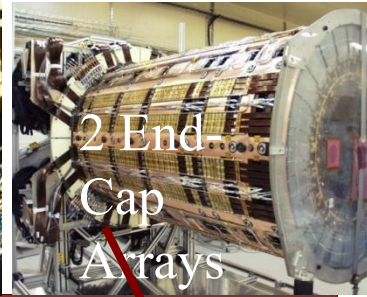
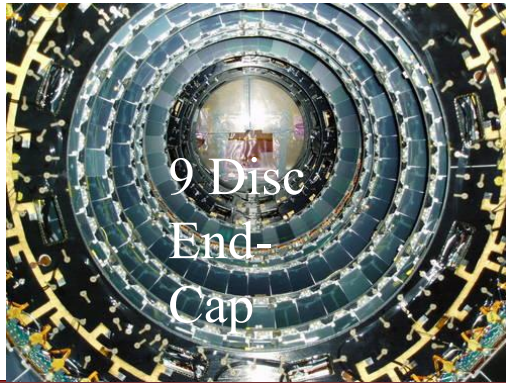
The LHCb-VeLo detector

All built in the Liverpool Semiconductor Detector Centre!



The ATLAS experiment

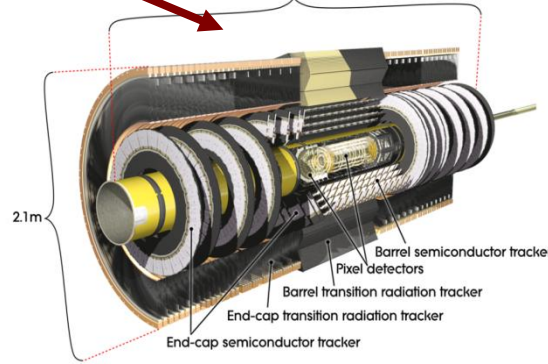
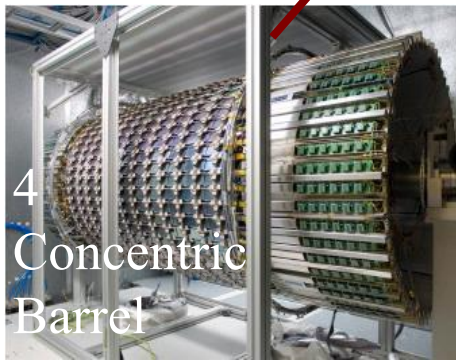
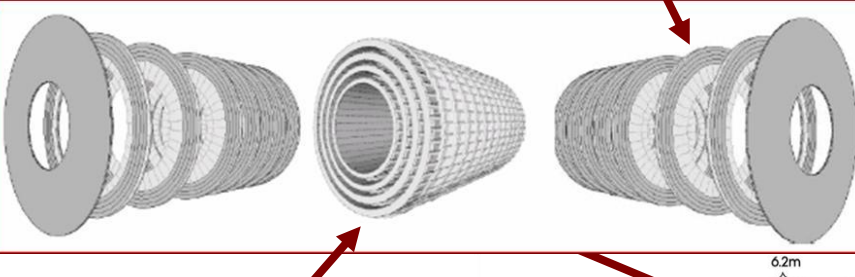
Liverpool and NIKHEF



2112 Barrel and 1976 End-Cap Double-Sided Modules
 61m² of silicon micro-strip detectors
 ~20,000 separate sensors ordered

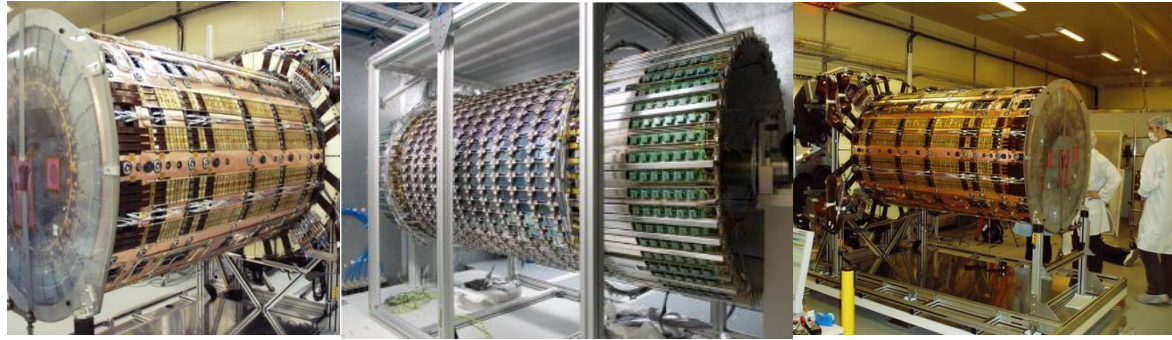


ATLAS

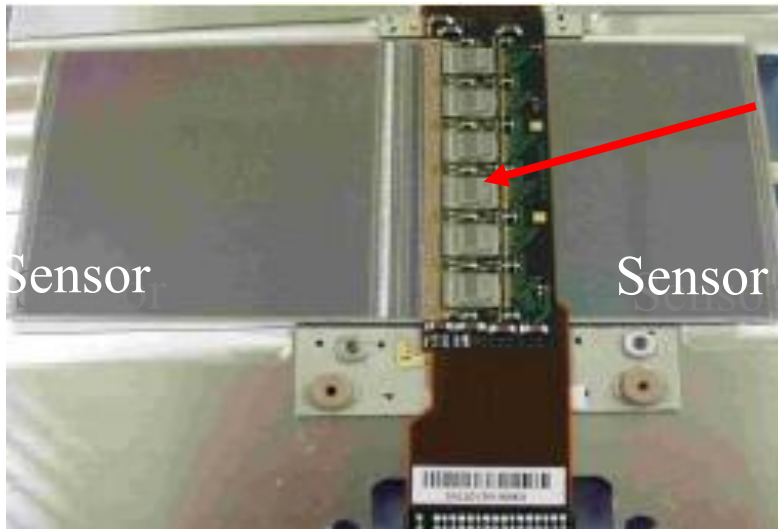


Oxford and CERN

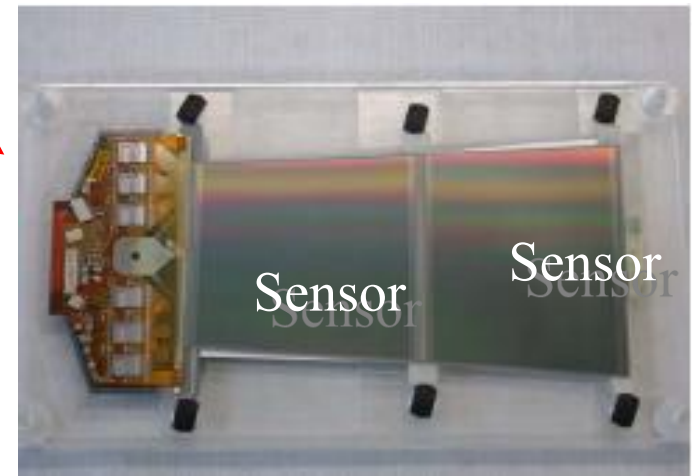
ATLAS Tracker Based on Barrel and Disc Supports



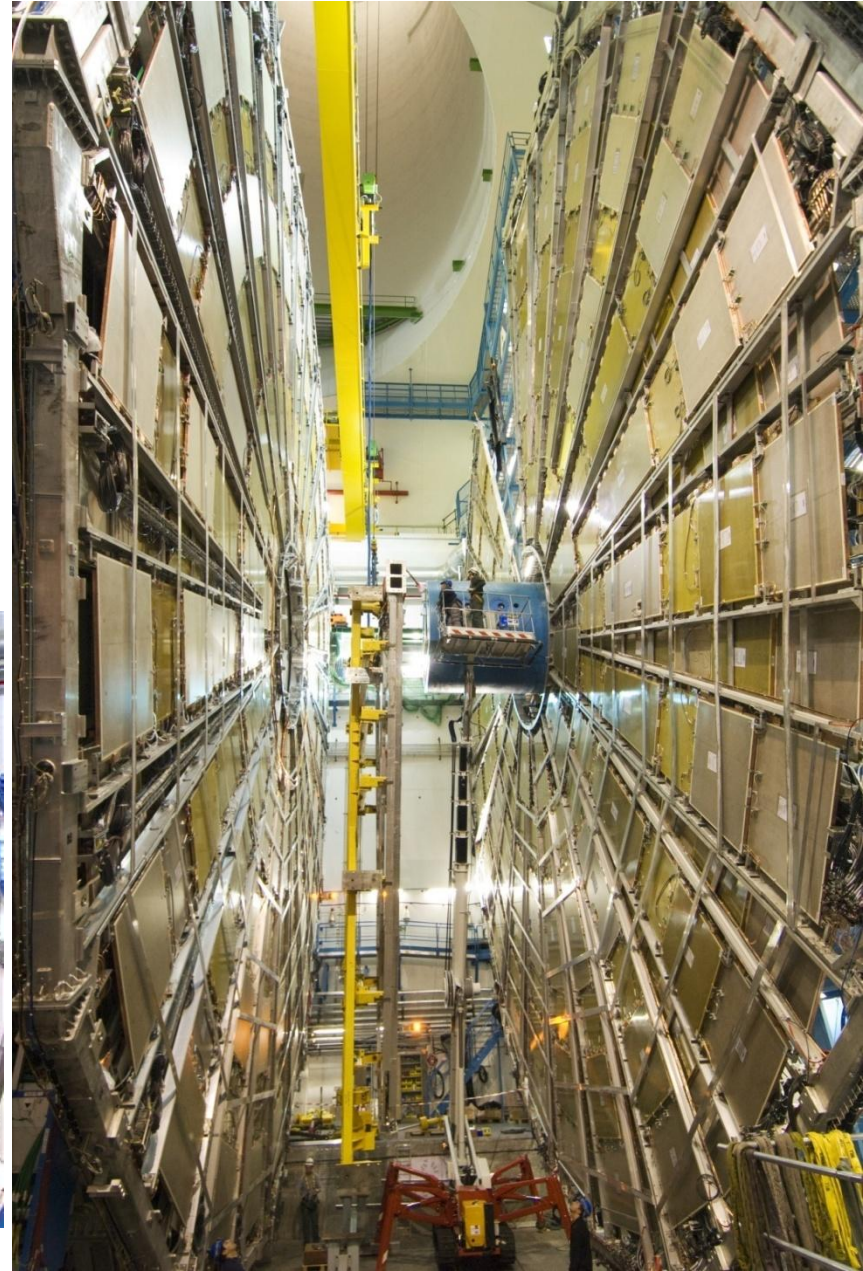
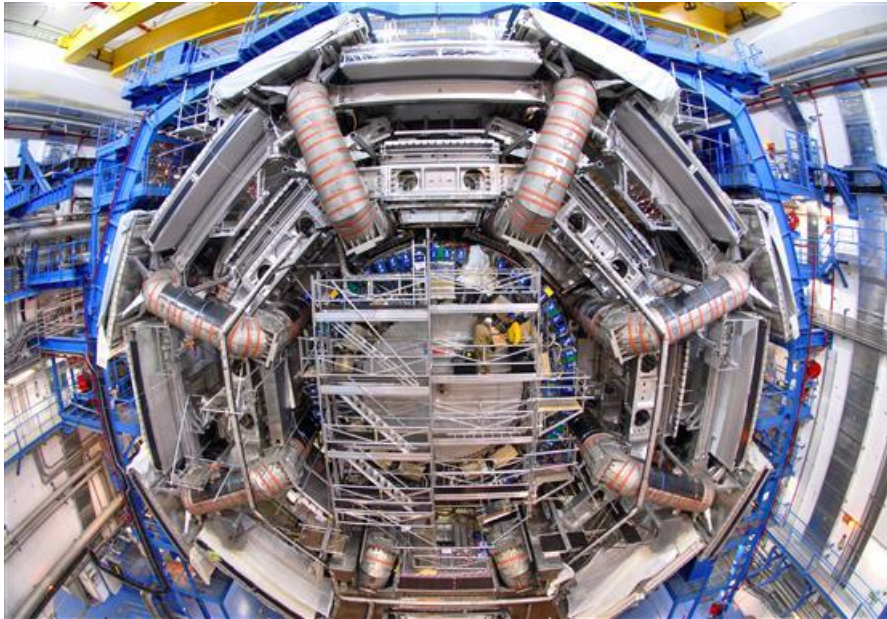
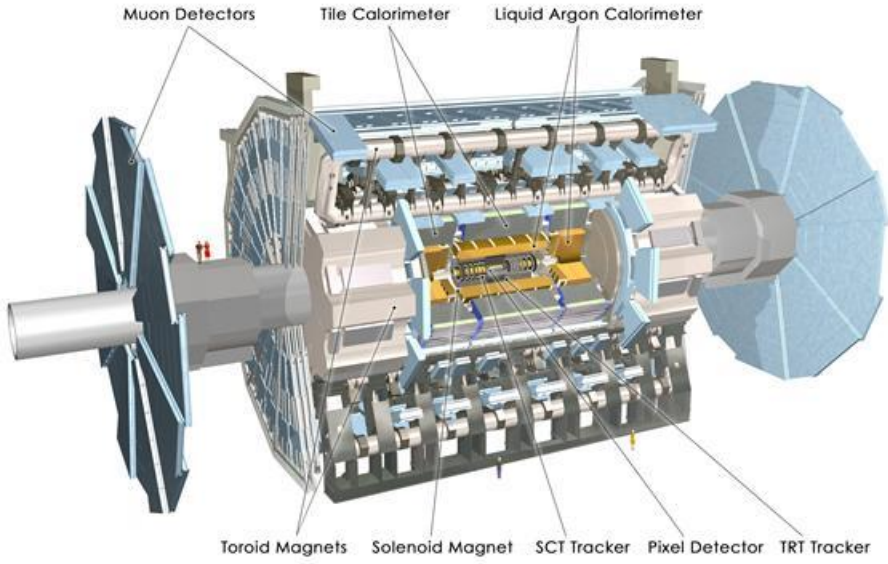
Effectively two styles of double-sided modules (2×6 cm long)
each sensor ~ 6 cm wide (768 strips of $80 \mu\text{m}$ pitch per side)



Hybrid
cards
carrying
read- out
chips and
multilayer
interconnect
circuit

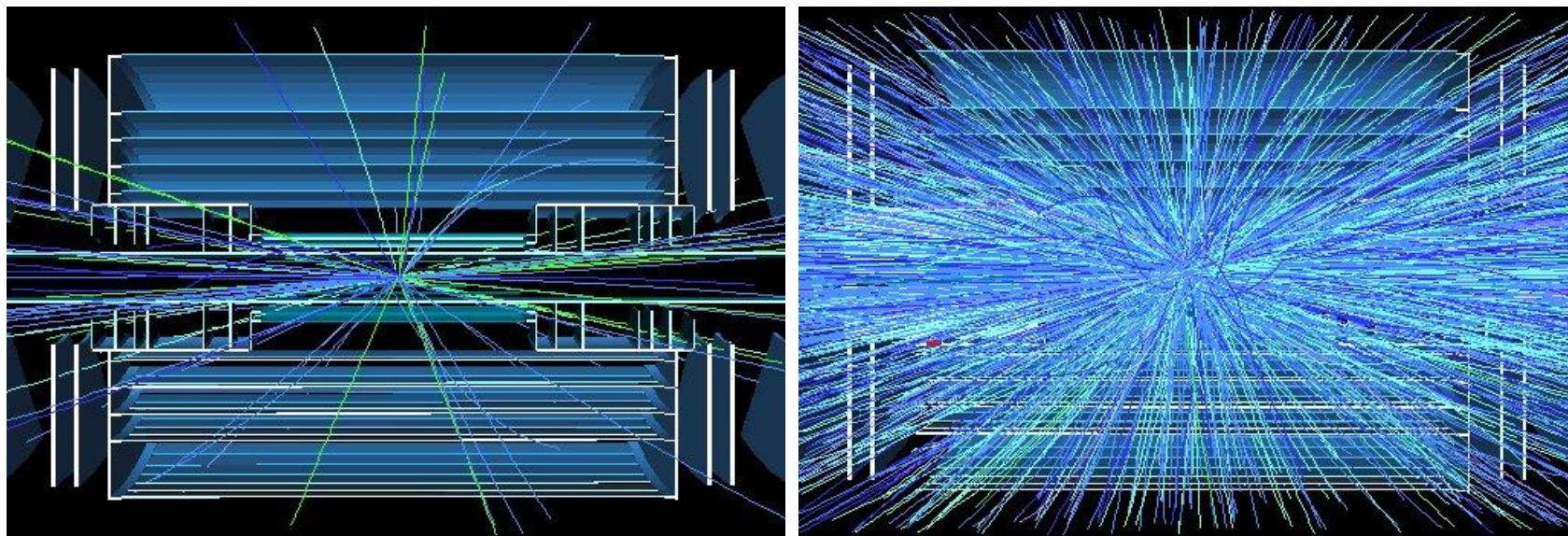


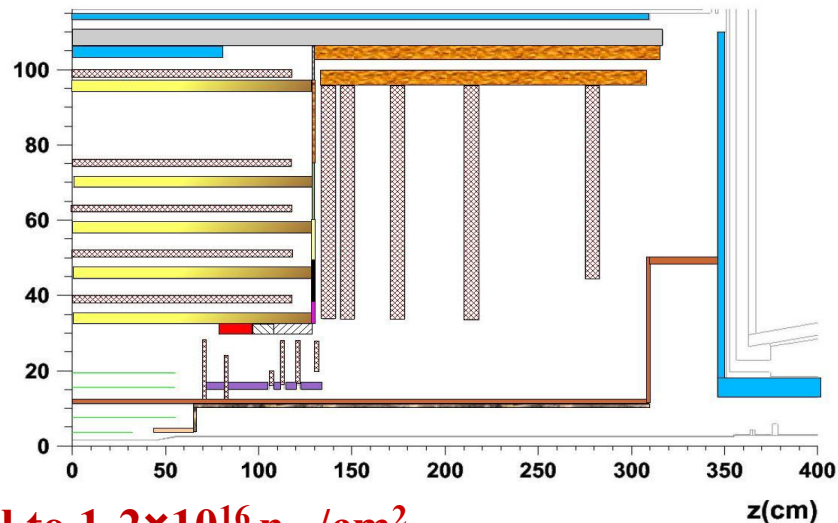
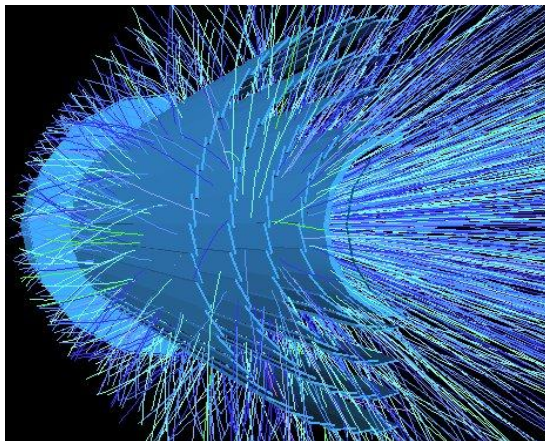
The ATLAS experiment



Upgrading the experiments: HL-LHC

The number of pile-up events per bunch crossing will be between 250-400! Extremely high multiplicity and radiation!





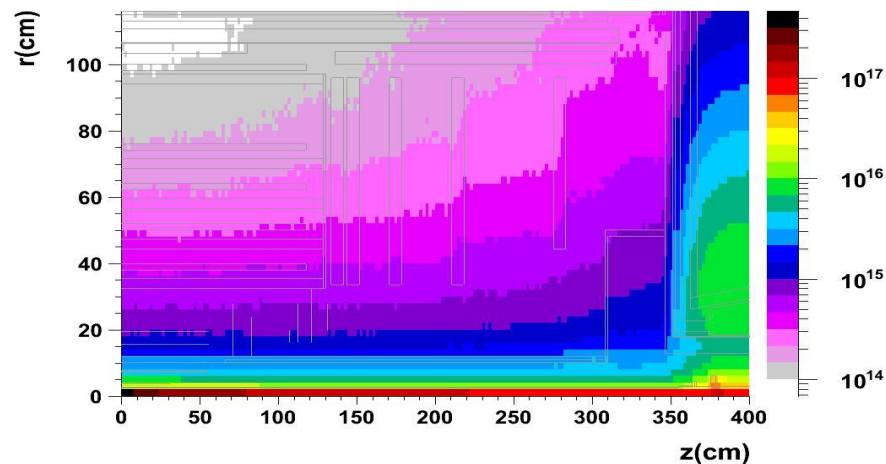
At inner pixel radii - target survival to $1-2 \times 10^{16} n_{eq}/cm^2$

For strips $\sim 1.3 \times 10^{15} n_{eq}/cm^2$

Numbers obtained 9/10/09 (corresponding to new layout) assuming 3000^{ph-1} and 84.5mb

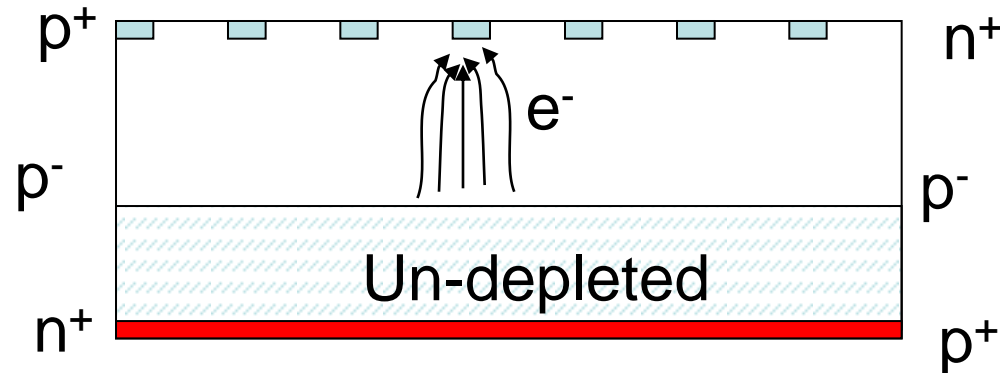
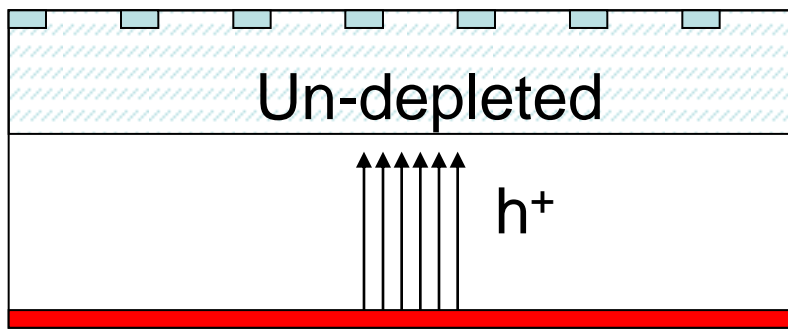
		1 MeV neutron eq fluence
Strip barrel 1 (SS)	($r=38cm$; $z=0cm$)	4.4×10^{14}
	($r=38cm$; $z=117cm$)	4.9×10^{14}
Strip barrel 4 (LS)	($r=74.3cm$; $z=0.0cm$)	1.6×10^{14}
	($r=74.3cm$; $z=117cm$)	1.8×10^{14}

Strip Disc 1 ($z=137.1$, Rinner=33.6)	6.0×10^{14}
Strip Disc 2 ($z=147.6$, Rinner=33.6)	6.2×10^{14}
Strip Disc 3 ($z=174.4$, Rinner=33.6)	5.8×10^{14}
Strip Disc 4 ($z=214.1$, Rinner=33.6)	6.1×10^{14}
Strip Disc 5 ($z=279.1$, Rinner=44.4)	5.8×10^{14}
Strip Disc 5 ($z=279.1$, Rinner=54.1)	4.4×10^{14}
Strip Disc 5 ($z=279.1$, Rinner=61.7)	3.9×10^{14}
new	
Strip Disc 5 ($z=279.1$, Rinner=73.6)	3.0×10^{14}
Strip Disc 5 ($z=279.1$, Rinner=84.9)	2.7×10^{14}



P-strip vs. N-strip Readout

“Standard” p-in-n geometry (after “New” n-in-p geometry type inversion)



Type inversion turns lightly doped material to “p” type

- Holes collected
- Deposited charge cannot reach electrode
 - Charge spread over many strips
 - Lower signal
- Electron collected
 - Higher mobility and $\sim 33\%$ smaller trapping constant
- Deposited charge can reach electrode

N-side read-out can make planar segmented Si detectors suitable for tracking in extreme (SLHC levels: $1-2 \times 10^{16} \text{ cm}^{-2}$) radiation environments.

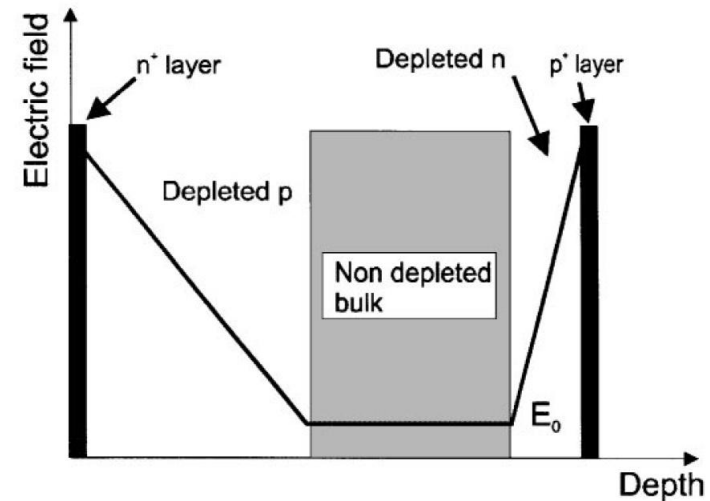
Schematic changes of Electric field after irradiation

Effect of trapping on the Charge Collection Efficiency (CCE)

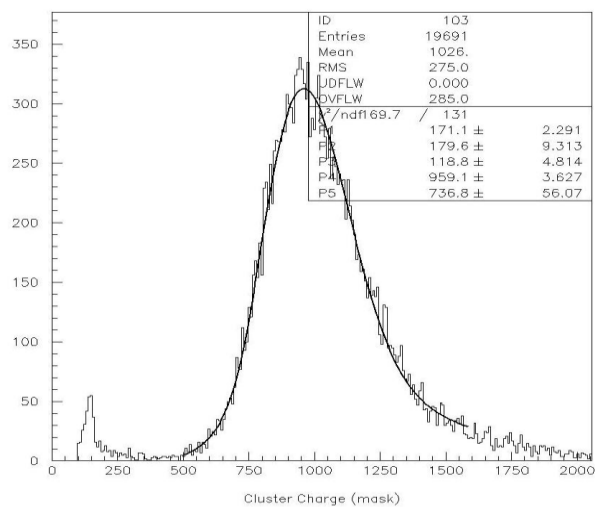
$$Q_{tc} \cong Q_0 \exp(-t_c / \tau_{tr}), \quad 1/\tau_{tr} = \beta \Phi.$$

Collecting electrons provide a sensitive advantage with respect to holes due to a much shorter t_c . P-type detectors are the most natural solution for e collection on the segmented side.

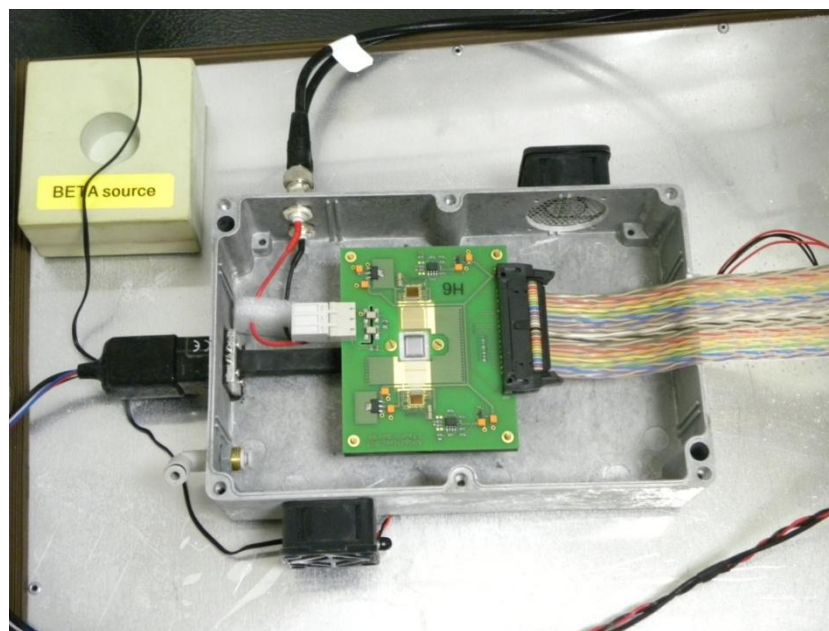
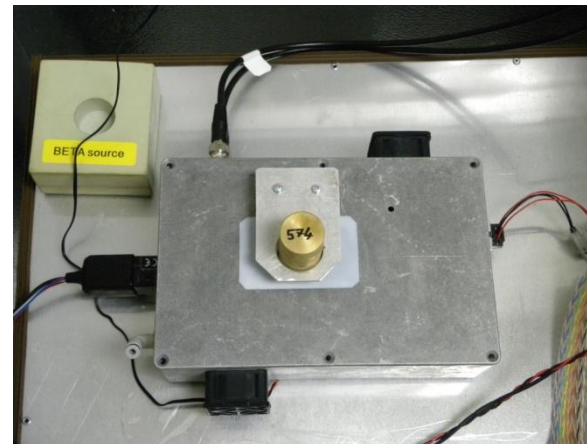
N-side read out to keep lower t_c



Mip signal from
 ^{90}Sr source



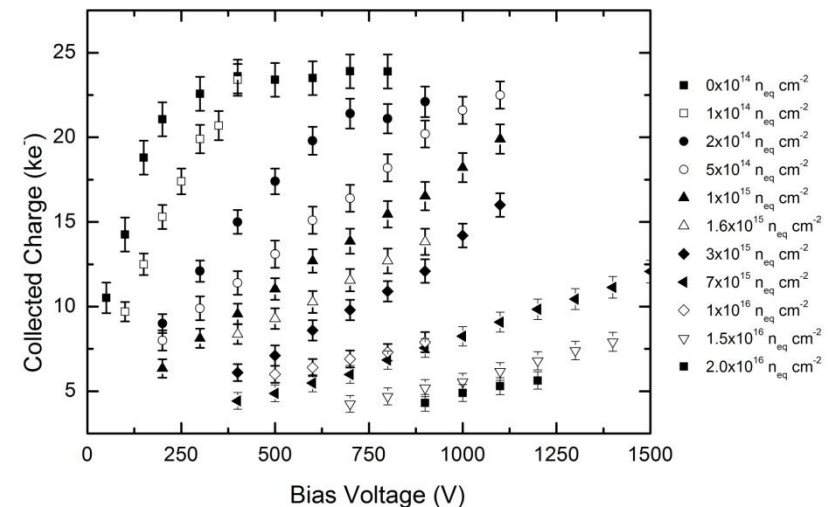
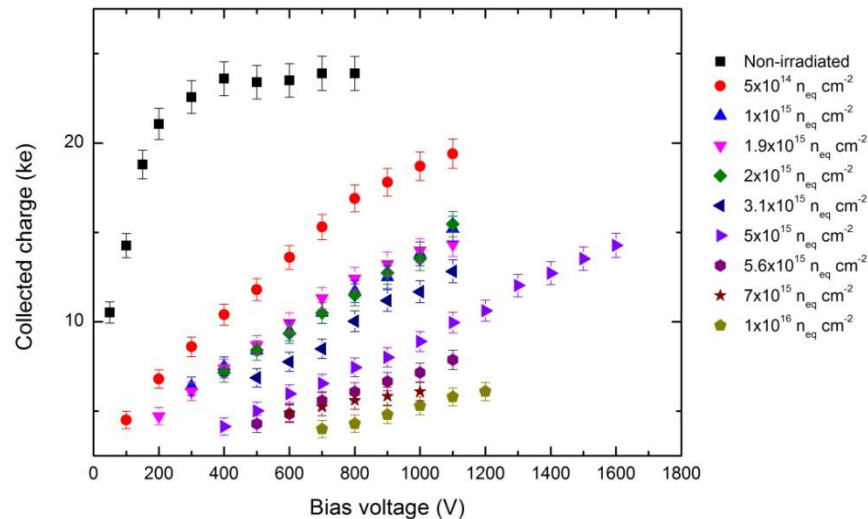
Analogue information
from LHC speed
analogue electronics

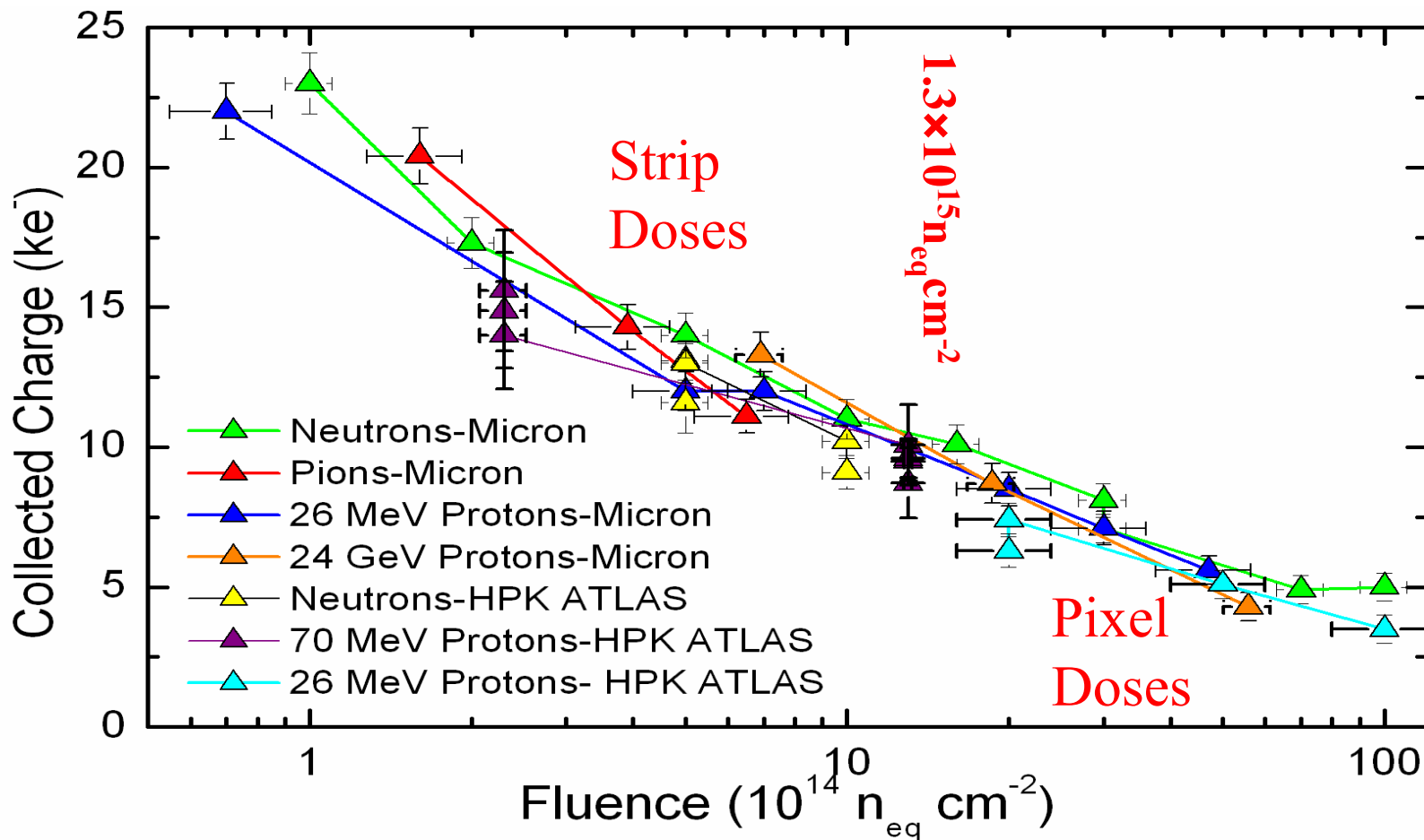


Using this method, the charge collection properties of n-in-p silicon detectors have been studied to the very highest fluences anticipated for the future HL-LHC at CERN.

Degradation of the charge collection with 23GeV and 26MeV proton fluence of 300 μm thick silicon detectors.

Degradation of the charge collection with reactor neutrons of 300 μm thick silicon detectors.





Effect of trapping on the Charge Collection Distance

After heavy irradiation the charge collection distance (CCD) of thin detectors should have a similar (better?) charge collection efficiency (CCE) as thicker ones.

The reverse current is proportional to the depleted volume in irradiated detectors. Do thin sensors offer an advantage in term of reduced reverse current compared to thicker ones (this aspect is particularly important for the inner layer detectors of SLHC, where significant contribution to power consumption is expected from the sensors themselves)?

$$Q_{tc} \cong Q_0 \exp(-t_c/\tau_{tr}), \quad 1/\tau_{tr} = \beta\Phi.$$

$$v_{sat,e} \times \tau_{tr} = \lambda_{av}$$

$$\beta_e = 4.2E-16 \text{ cm}^{-2}/\text{ns}$$

$$\beta_h = 6.1E-16 \text{ cm}^{-2}/\text{ns}$$

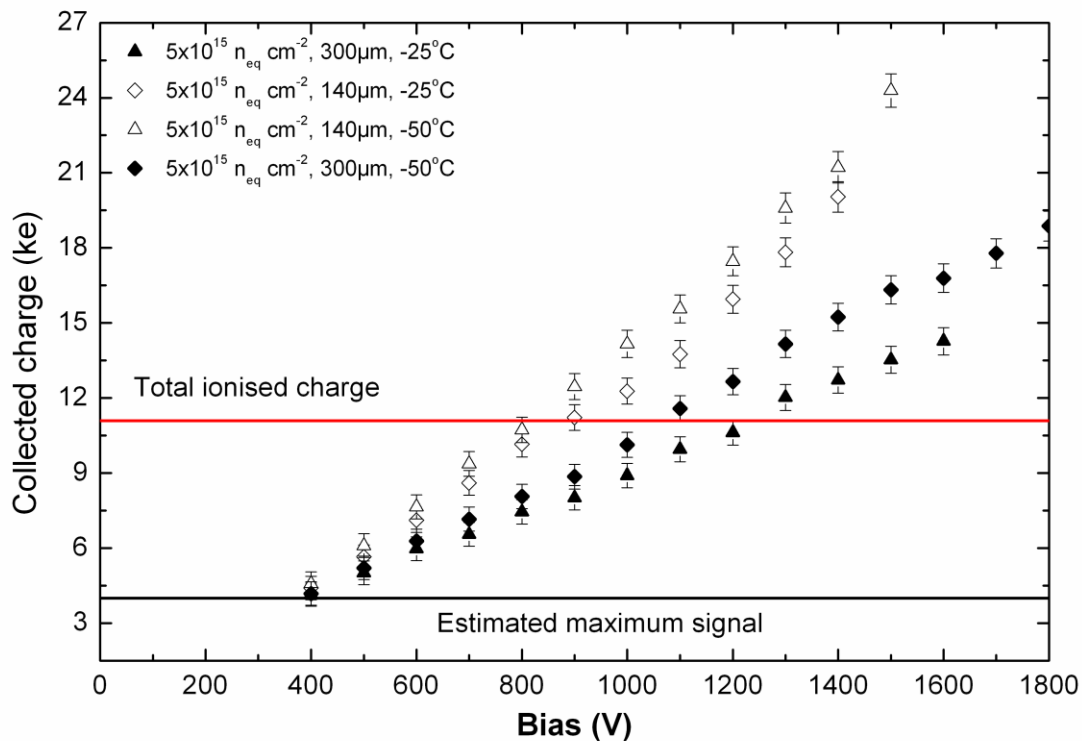
$$\lambda_{Max,n} (\Phi=1e14) \cong 2400\mu\text{m}$$

$$\lambda_{Max,n} (\Phi=1e16) \cong 24\mu\text{m}$$

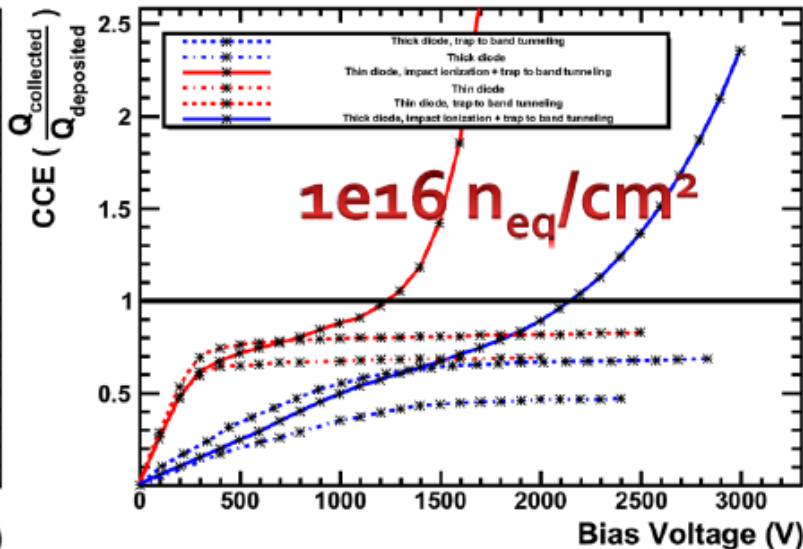
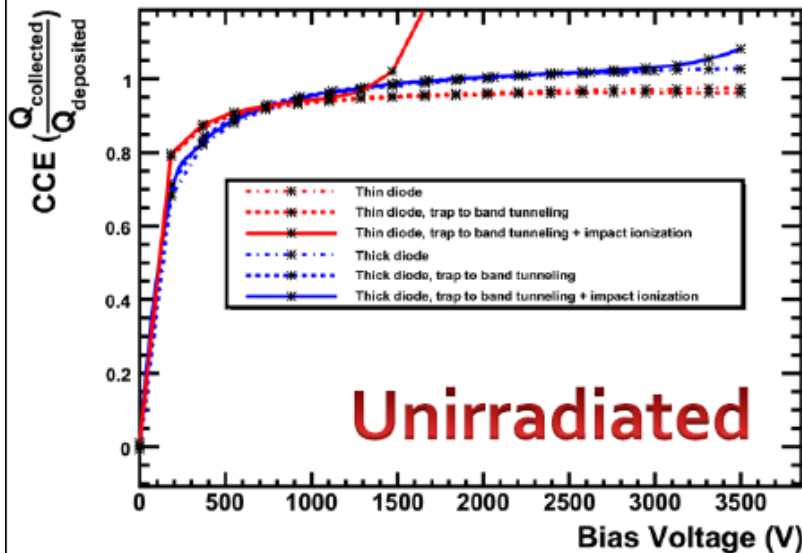
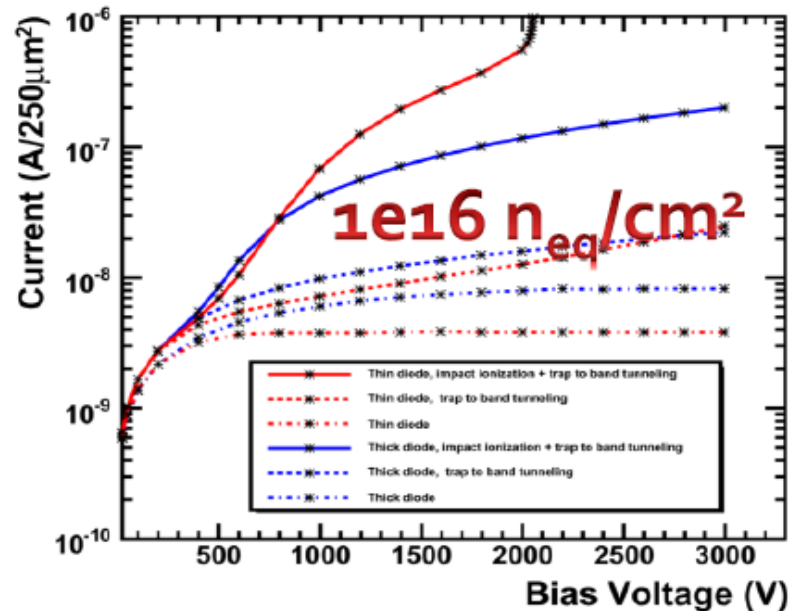
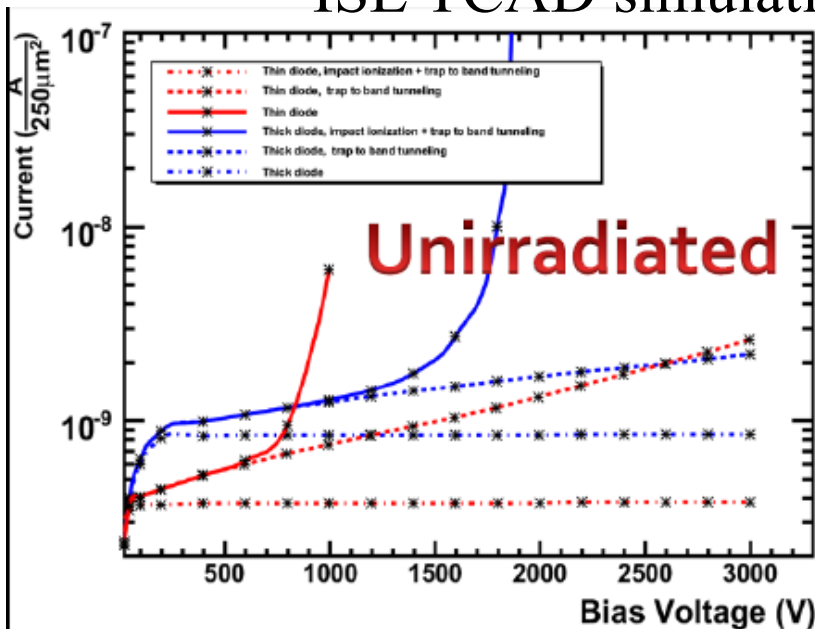
$$\lambda_{Max,p} (\Phi=1e14) \cong 1600\mu\text{m}$$

$$\lambda_{Max,p} (\Phi=1e16) \cong 16\mu\text{m}$$

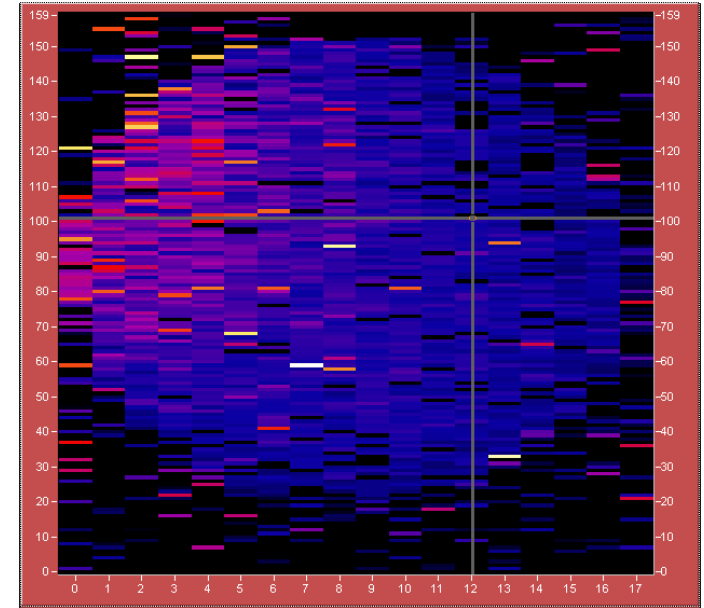
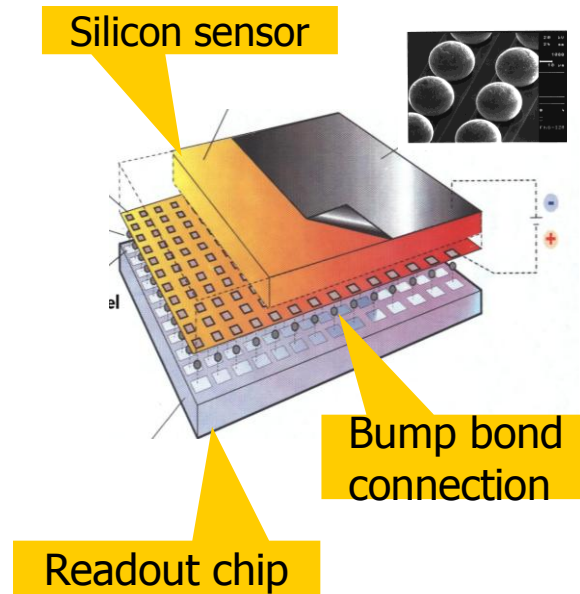
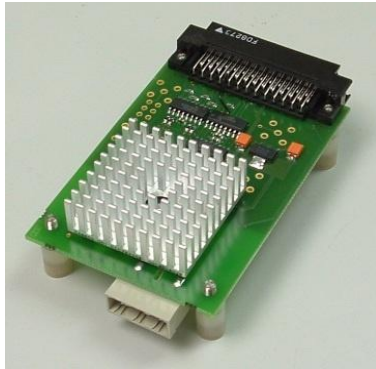
Charge multiplication by impact ionisation increases the signal after high doses!



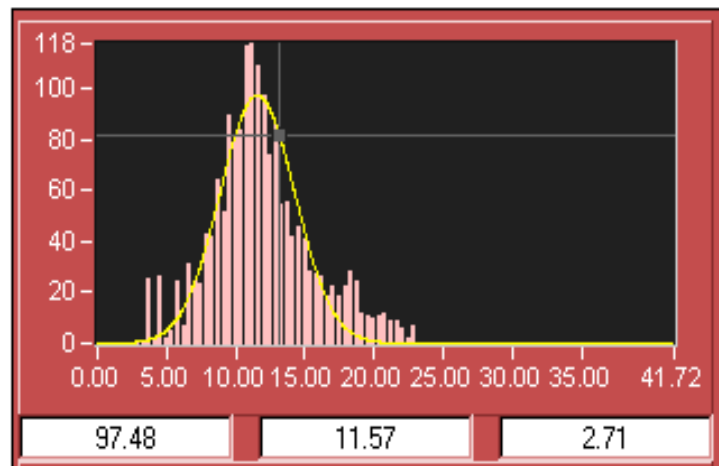
ISE TCAD simulations



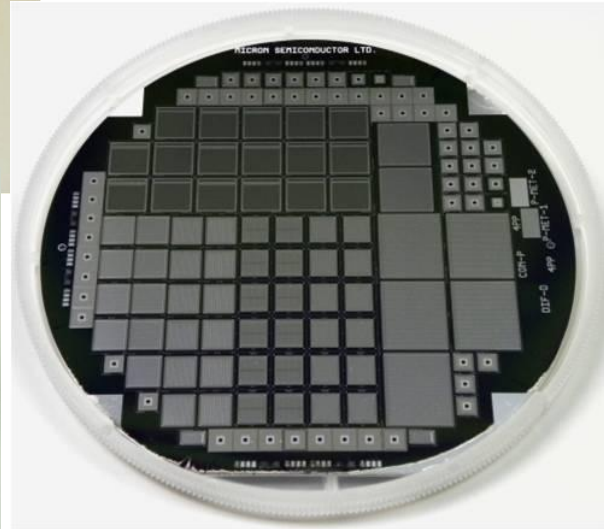
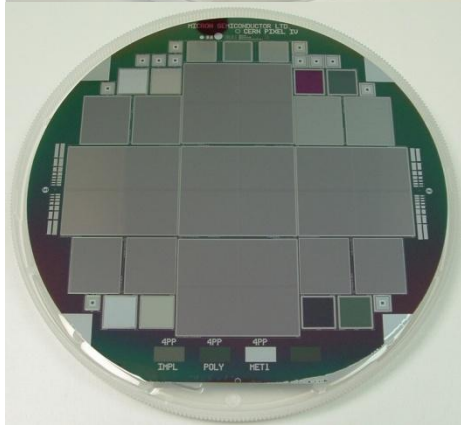
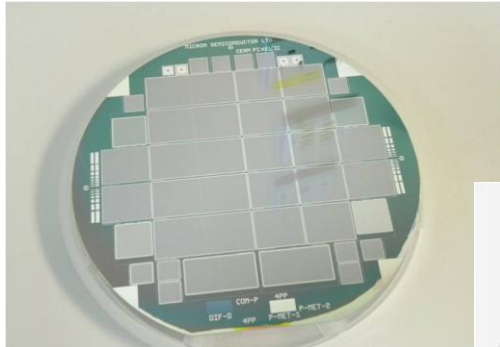
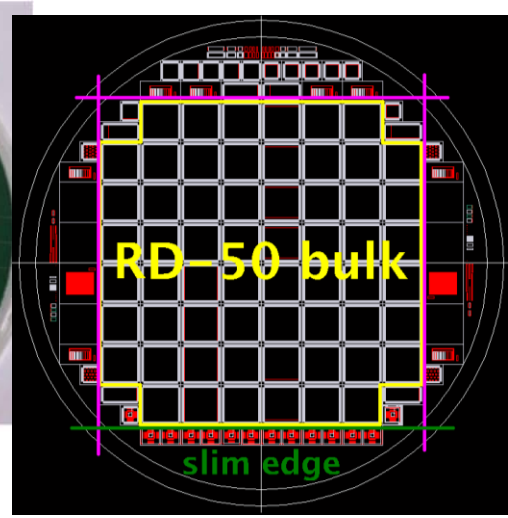
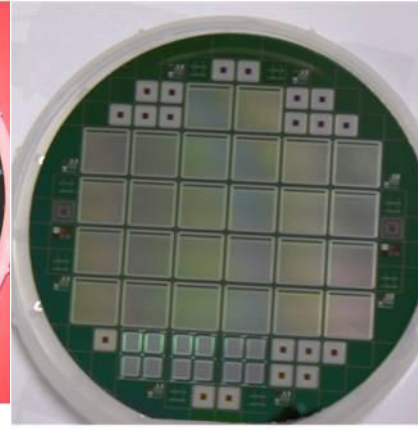
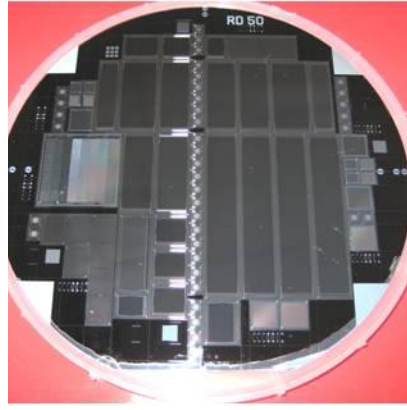
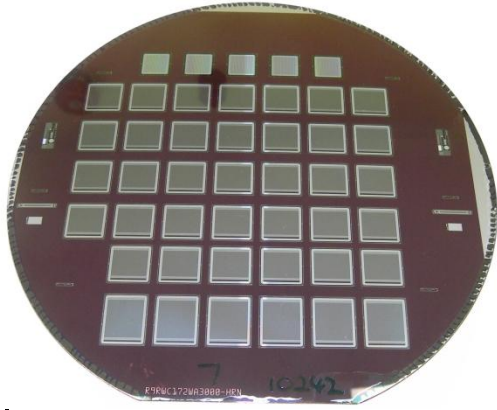
Radiation hardness and fine granularity



After irradiation to
 $1.5 \times 10^{16} \text{ p/cm}^2$
 At CERN PS
 $(9 \times 10^{15} \text{ n}_{\text{eq}} \text{ cm}^{-2})$
 peak charge at
 500V is $\sim 4000\text{e}$
 (Threshold 3500e
 -26°C , I_b 44 μA)



R&D for future experiments



Knowledge of production procedures and technologies. Most extreme radiation tolerant devices in the world. Collaboration with different producers.