



UNIVERSITY OF  
LIVERPOOL

Enhanced efficiency of segmented  
silicon detectors of various  
thicknesses after hadron irradiations  
up to  $2 \times 10^{16} n_{eq} \text{ cm}^2$ .

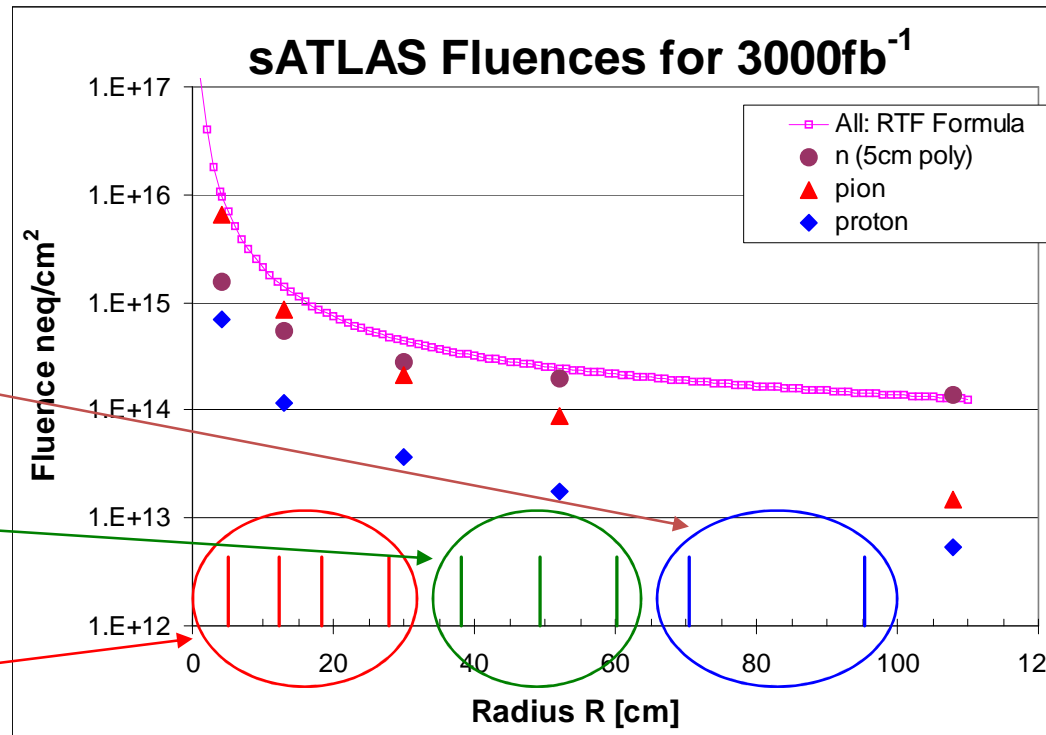
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H. Brown, M. Wormald

# OUTLINE

- Comparison of thin and thick sensors after neutron and proton irradiation
- Studies at low temperature (-50°C)
- Non-linearity of  $N_{\text{eff}}$  vs  $\Phi$
- Conclusions

# Fluence in Proposed sATLAS Tracker

Strip length and segmentation determined by occupancy < 2%



Mix of **neutrons**, **protons**, **pions** depending on radius **R**

Long and short strips damage largely due to **neutrons**

**Pixels** damage due to **neutrons** and **pions**

ATLAS Radiation Taskforce [http://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/RADIATION/RadiationTF\\_document.html](http://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/RADIATION/RadiationTF_document.html)

## Design fluences for sensors (includes 2x safety factor) :

Innermost Pixel Layer:	$1-1.6 \cdot 10^{16} \text{ n}_{\text{eq}}/\text{cm}^2 = 500 \text{ Mrad}$
Outer Pixel Layers:	$3 \cdot 10^{15} \text{ n}_{\text{eq}}/\text{cm}^2 = 150 \text{ Mrad}$
Short strips:	$1 \cdot 10^{15} \text{ n}_{\text{eq}}/\text{cm}^2 = 50 \text{ Mrad}$
Long strips:	$4 \cdot 10^{14} \text{ n}_{\text{eq}}/\text{cm}^2 = 20 \text{ Mrad}$

Need to study response to both neutral (neutrons) and charged (proton) particle irradiations

# Silicon miniature microstrip detectors and irradiation

[RD50 mask set](http://rd50.web.cern.ch/rd50/) (see: <http://rd50.web.cern.ch/rd50/>)

Miniature sensors,  $\sim 1 \times 1 \text{ cm}^2$ , 128 strips, 80  $\mu\text{m}$  pitch, designed by Liverpool and produced by Micron on 300 $\mu\text{m}$  and 140 $\mu\text{m}$  thick wafers.

Irradiation and dosimetry:

**Neutron:**

TRIGA Mark II research reactor

Reactor Centre of the

[Jozef Stefan Institute, Ljubljana](#), Slovenia, thanks to

V. Cindro

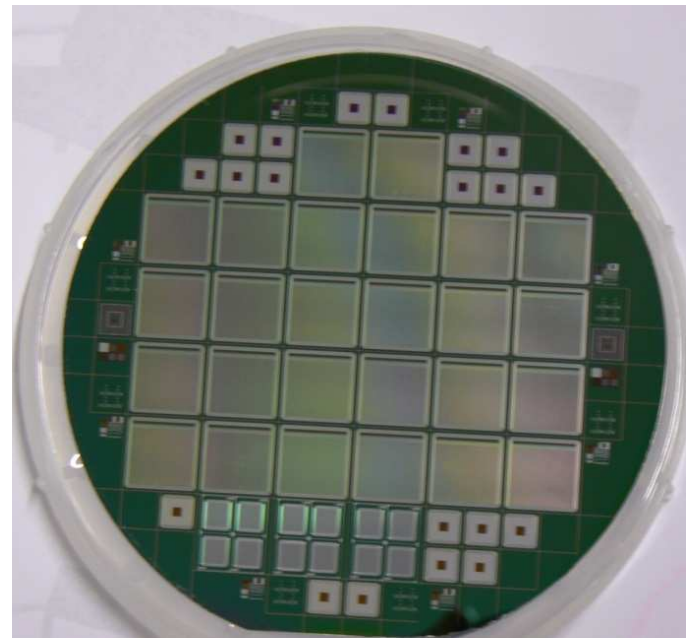
**24 GeV protons:**

[CERN-PS Irrad1](#), irradiated at area temperature ( $\sim 30^\circ\text{C}$ ) and cold ( $< 5^\circ\text{C}$ ). thanks to M.

Glaser.

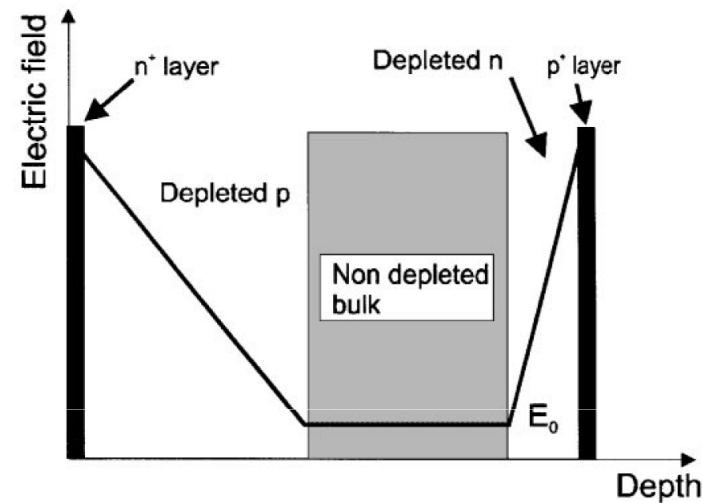
**26 MeV protons:**

Compact Cyclotron of the University of Karlsruhe, thanks to A. Dierlamm and W. de Boer.



N-side read-out can make planar segmented Si detectors suitable for tracking in extreme (SLHC levels:  $1-2 \times 10^{16} \text{ cm}^{-2}$ ) radiation environments.

Schematic changes of Electric field after irradiation



Effect of trapping on the Charge Collection Efficiency (CCE)

$$Q_{tc} \cong Q_0 \exp(-t_c / \tau_{tr}), \quad 1/\tau_{tr} = \beta \Phi.$$

Collecting electrons provide a sensitive advantage with respect to holes due to a much shorter  $t_c$ . P-type detectors are the most natural solution for  $e$  collection on the segmented side.

N-side read out to keep lower  $t_c$

## Effect of trapping on the Charge Collection Distance

After heavy irradiation the charge collection distance (CCD) of thin detectors should have a similar (better?) charge collection efficiency (CCE) as thicker ones.

$$Q_{tc} \cong Q_0 \exp(-t_c/\tau_{tr}), 1/\tau_{tr} = \beta\Phi.$$

$$v_{sat,e} \times \tau_{tr} = \lambda_{av}$$

$$\beta_e = 4.2E-16 \text{ cm}^{-2}/\text{ns} \quad \text{G. Kramberger et al., NIMA 476(2002), 645-651.}$$

$$\beta_h = 6.1E-16 \text{ cm}^{-2}/\text{ns}$$

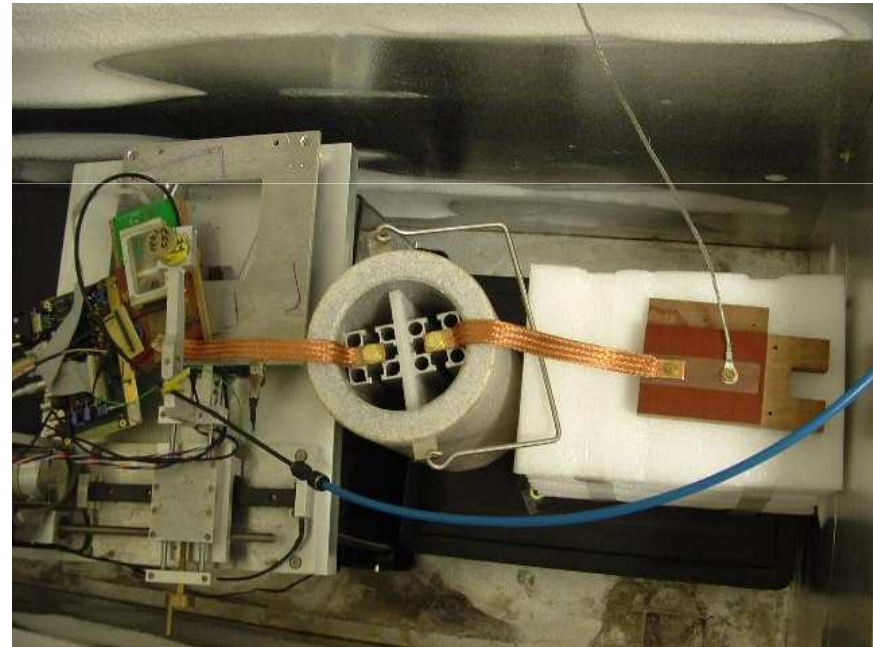
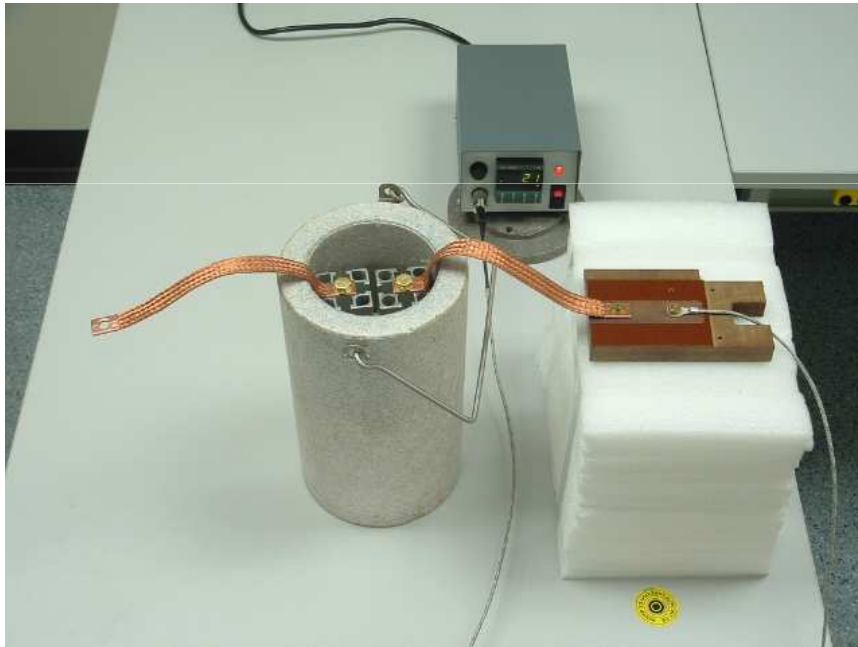
$$\lambda_{av} (\Phi=1e14) \cong 2400\mu\text{m}$$

$$\lambda_{av} (\Phi=1e16) \cong 24\mu\text{m}$$

The reverse current is proportional to the depleted volume in irradiated detectors. Do thin sensors offer an advantage in term of reduced reverse current compared to thicker ones (this aspect is particularly important for the inner layer detectors of SLHC, where significant contribution to power consumption is expected from the sensors themselves)?

## MEASUREMENTS: at -25 and -45/50 °C.

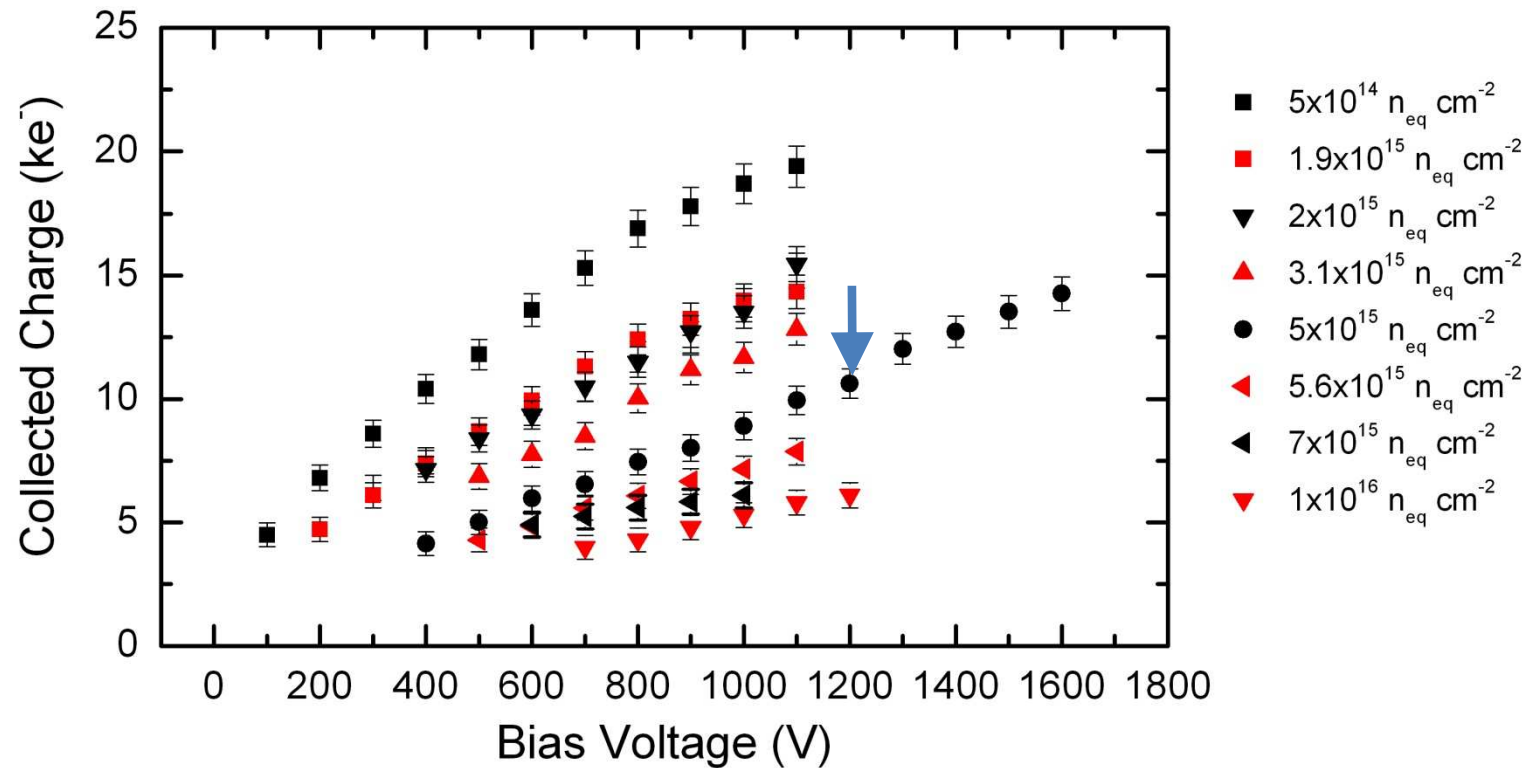
Cooling: “cold finger” system in liquid nitrogen bath. 40MHz analogue readout.  $^{90}\text{Sr}$  with plastic scintillator trigger, with plastic filter to remove the lower energy electrons (mip's). MP value is measured, calibration with a 300 $\mu\text{m}$  non irradiated sensor.



# Results with proton irradiated 300 $\mu\text{m}$ n-in-p Micron sensors (up to $1 \times 10^{16} \text{ n}_{\text{eq}} \text{ cm}^{-2}$ )

RED: irradiated with  
24GeV/c protons  
Other: 26MeV protons

24GeV/c protons irradiated COLD,  
all but the  $3.1 \text{E}15 \text{ cm}^{-2}$  series,  
which had  $\sim 10$  days @  $33^\circ\text{C}$

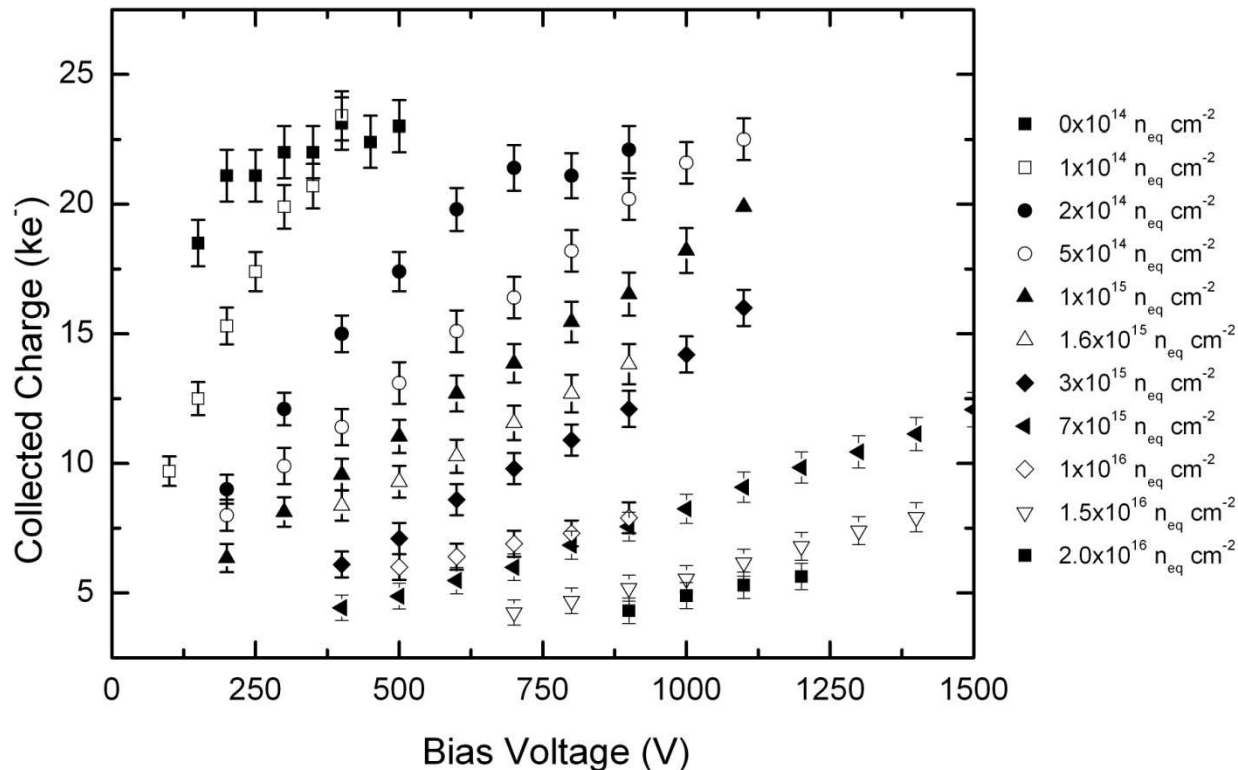




# Results with neutron irradiated 300 $\mu\text{m}$ n-in-p Micron sensors (up to $2 \times 10^{16} n_{\text{eq}} \text{ cm}^{-2}$ )

Short irradiation time,  
practically no annealing

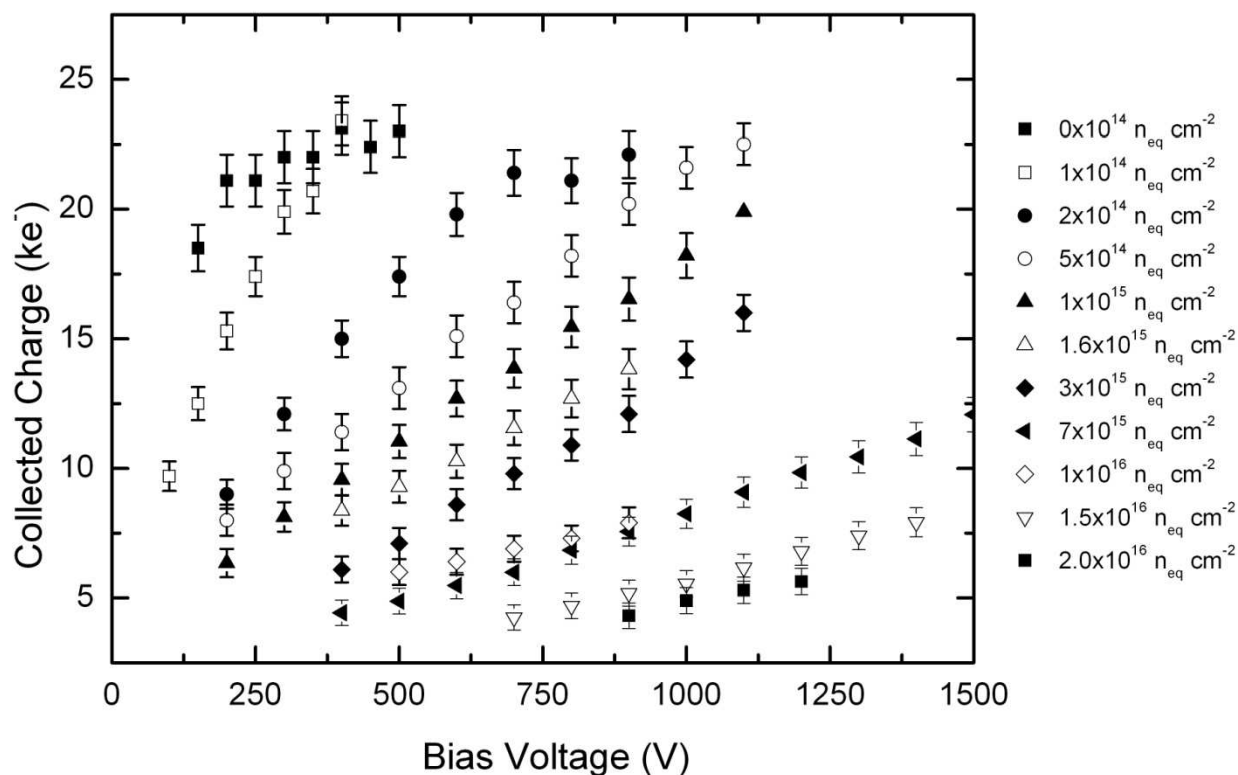
After both p and n irradiations, the  
collected charge is much higher than  
expected from charge trapping  
measurements ( $< 2000e$  after  $1 \times 10^{16} \text{ cm}^{-2}$ ).



# Results with neutron irradiated 300 $\mu\text{m}$ n-in-p Micron sensors (up to $2 \times 10^{16} \text{ n}_{\text{eq}} \text{ cm}^{-2}$ )

Short irradiation time,  
practically no annealing

Is there a mechanism to explain this  
enhanced signal, at high bias voltages and  
after heavy irradiation?

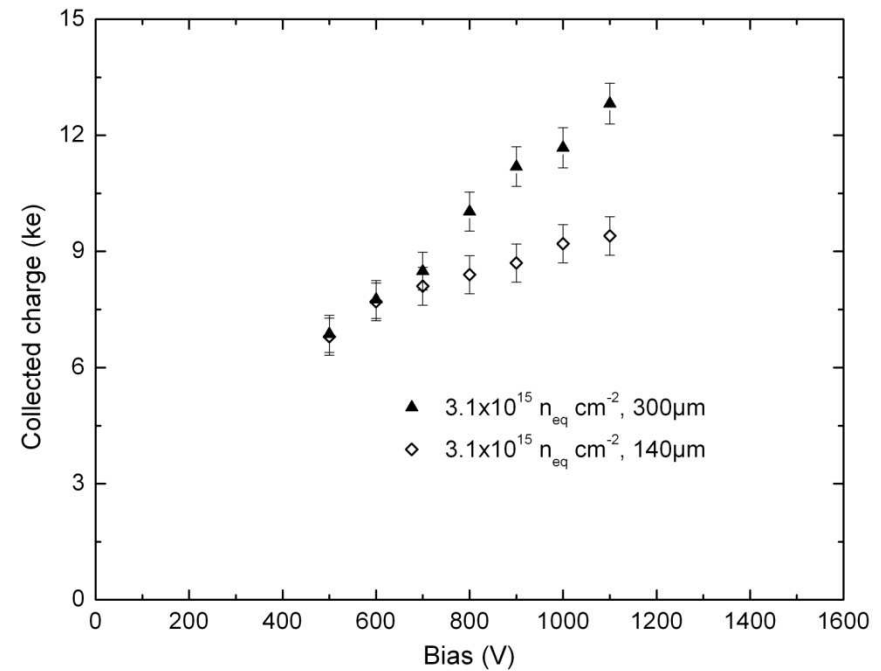
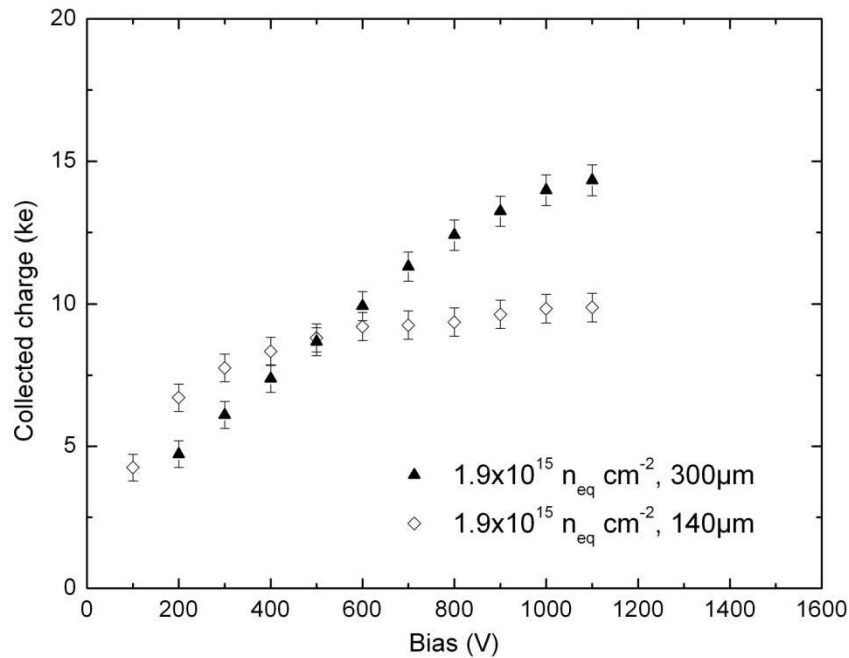


This was achieved with 300 $\mu$ m thick sensors (High Energy Physics standard). The collected charge after extreme doses is much higher than expected from the dependence of  $\tau_{tr}$  with fluence. Moreover, can thin devices be even better?

# 140 and 300 $\mu\text{m}$ n-in-p Micron sensors after 1.9 and $3.1 \times 10^{15} n_{\text{eq}}$ 24GeV/c p

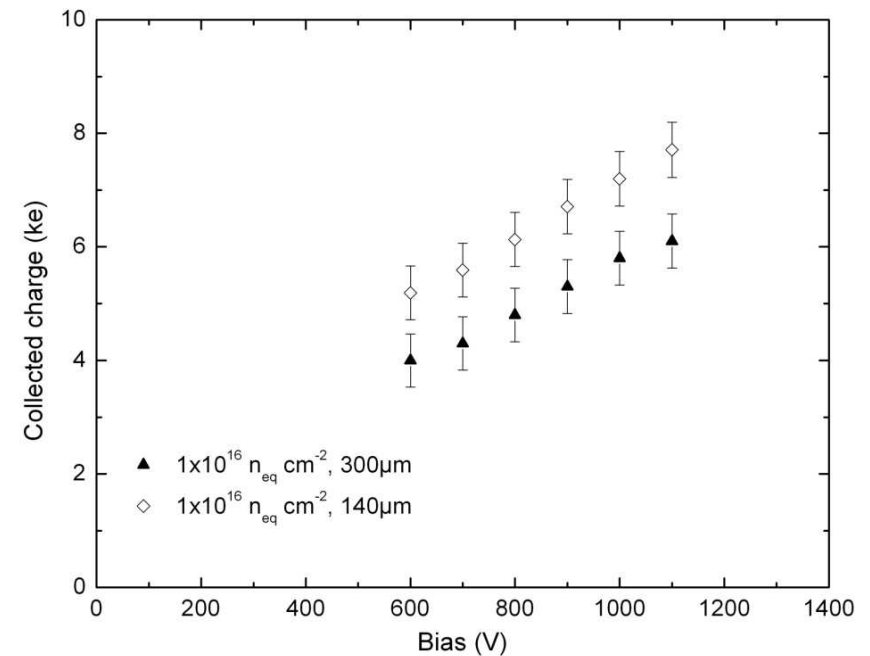
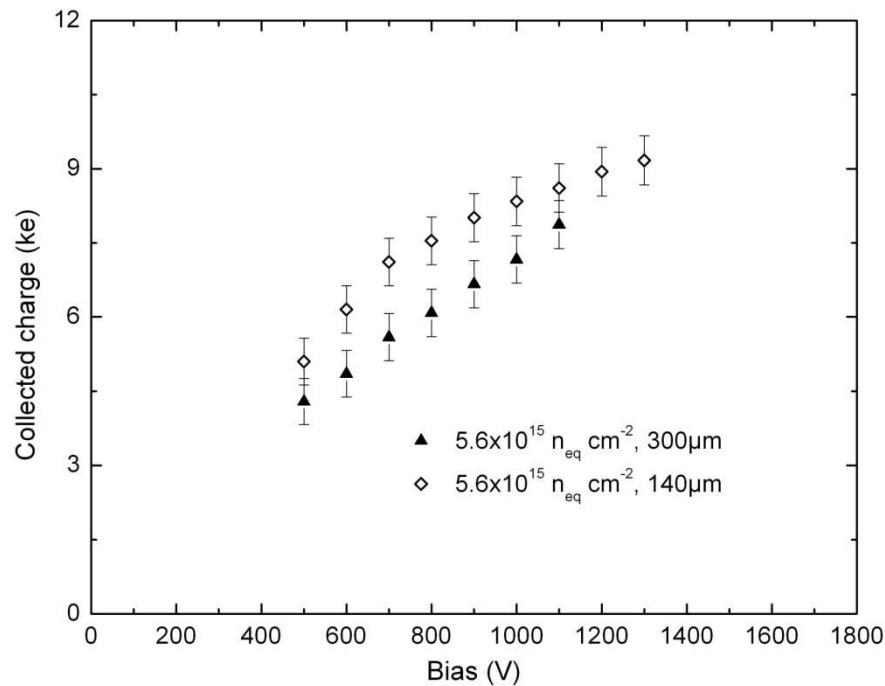
Room T (29-31  $^{\circ}\text{C}$ )  
irradiation

Cold(0-5  $^{\circ}\text{C}$ ) irradiation



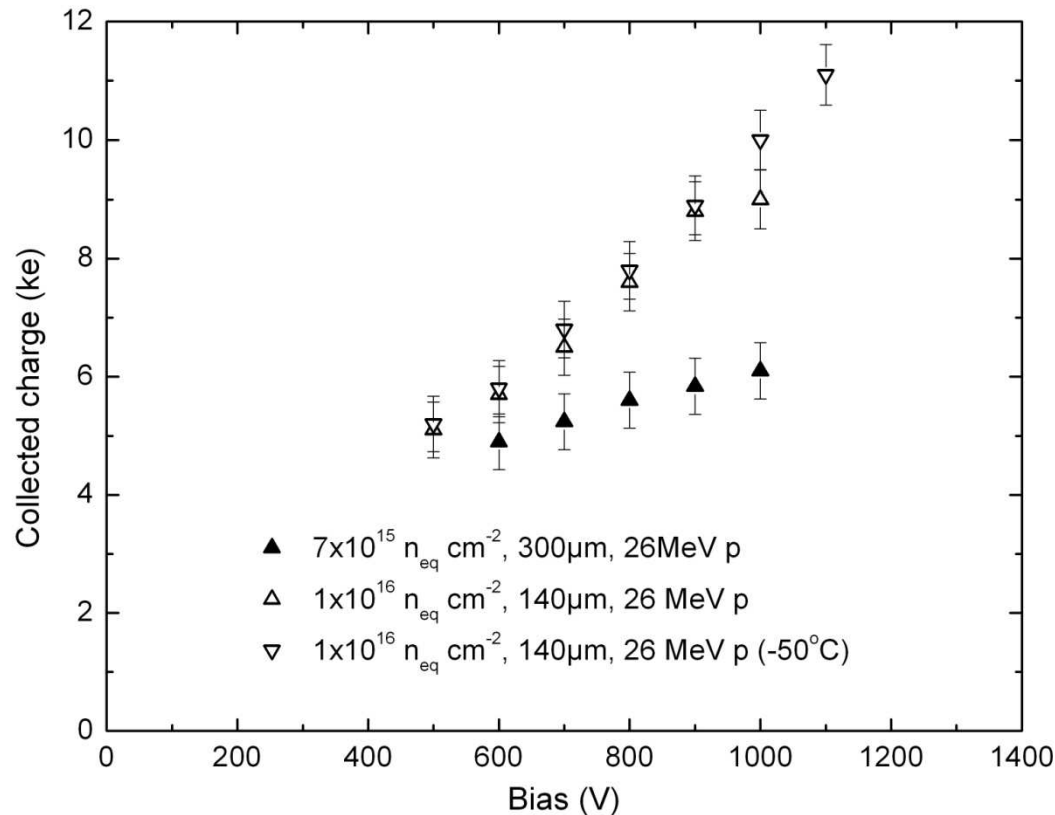
# 140 and 300 $\mu\text{m}$ n-in-p Micron sensors after 5.6 and $1 \times 10^{16} \text{ n}_{\text{eq}} \text{ 24GeV/c p}$

Cold(0-5 °C) irradiation



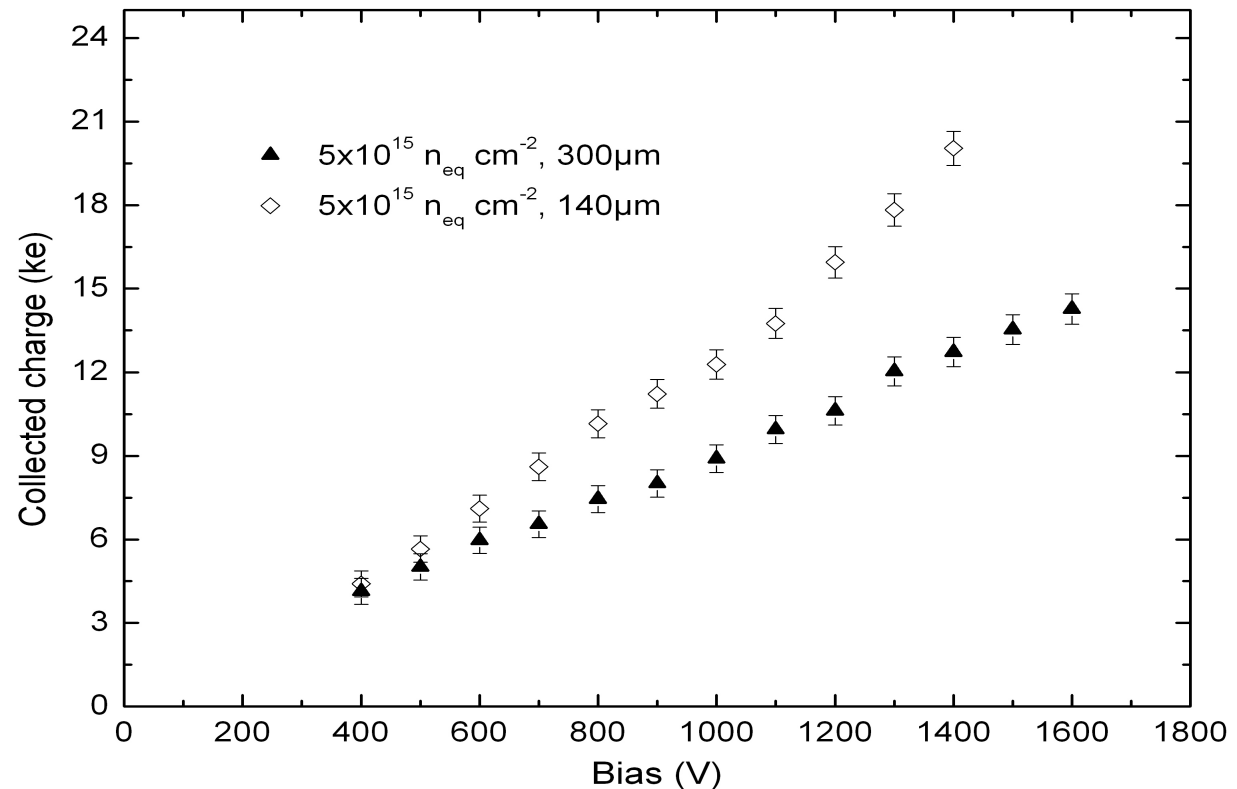
# 140 and 300 $\mu\text{m}$ n-in-p Micron sensors after $1 \times 10^{16}$ $n_{\text{eq}}$ 26MeV p

Even after heavy irradiation it is possible to recover the whole ionised charge (the 140 $\mu\text{m}$  thick sensors collect here  $\sim 12000e$ ).



# 140 and 300 $\mu\text{m}$ n-in-p Micron sensors after $5 \times 10^{15}$ $n_{\text{eq}}$ 26MeV p

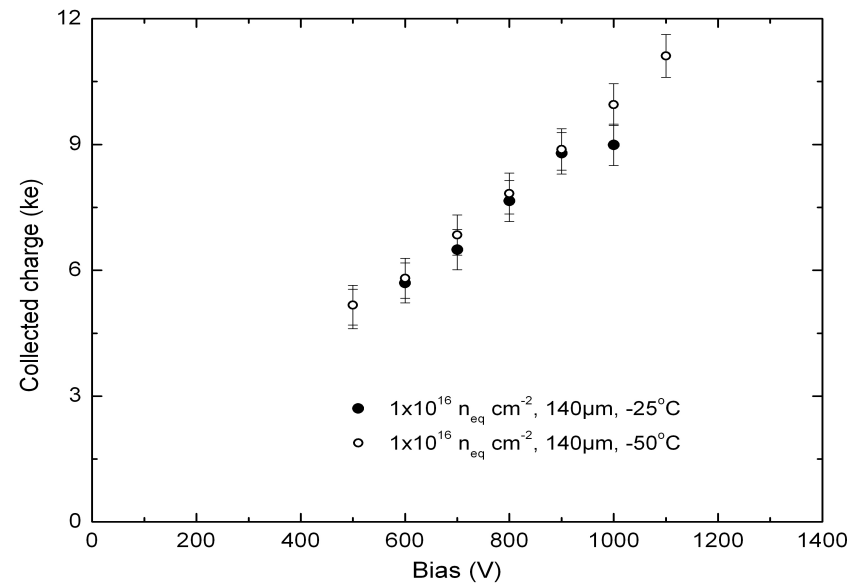
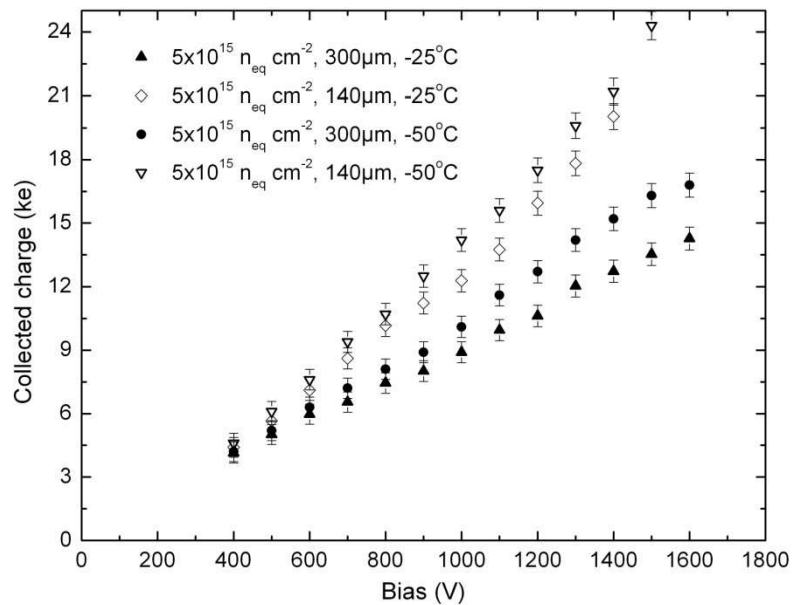
Evidence of a charge multiplication effect: not only the whole charge is recovered, but increased by  $f = 1.75$



# 140 and 300 $\mu\text{m}$ n-in-p Micron sensors after 5 (26MeV p) and $10 \times 10^{15} n_{\text{eq}}$ (24GeV/c p) at low ( $-50^\circ\text{C}$ ) T.

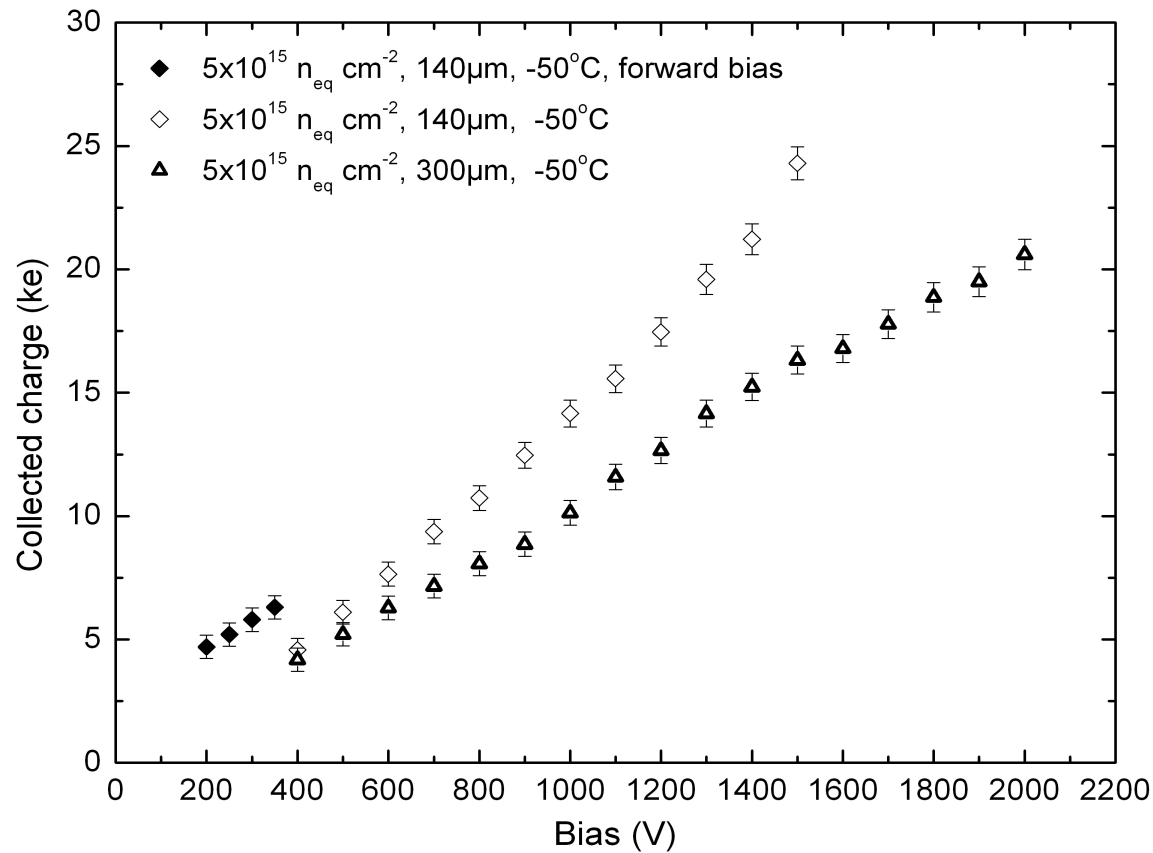
**F = 2.1 at  $-50^\circ\text{C}$ !**

**Pre-irradiation charge recovered after  $1 \times 10^{16} n_{\text{eq}}$ .**



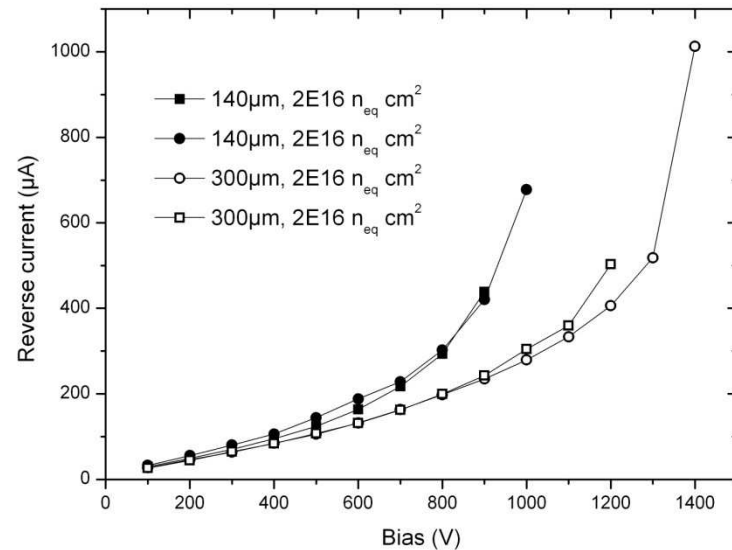
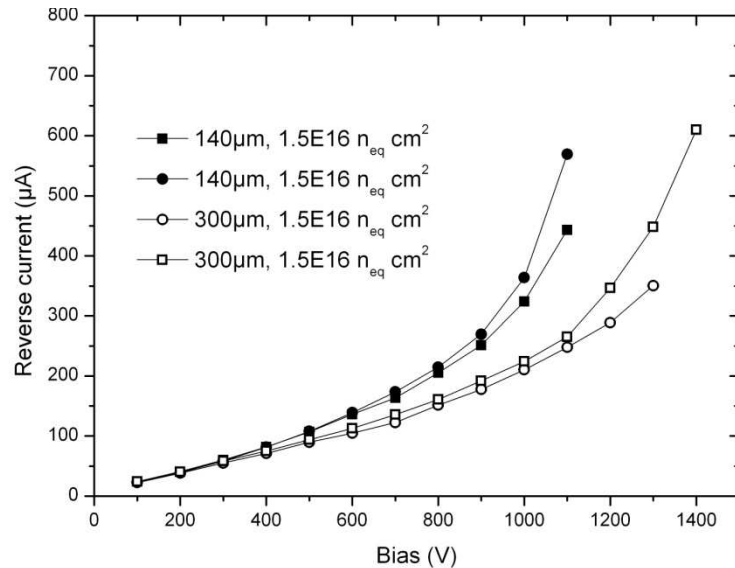
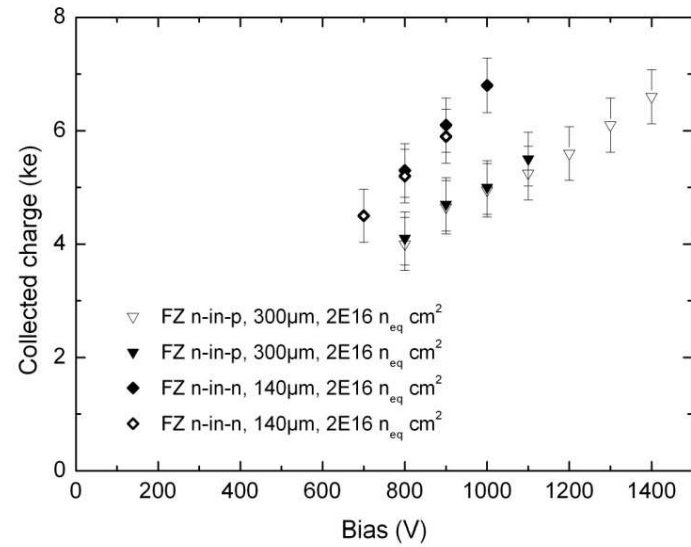
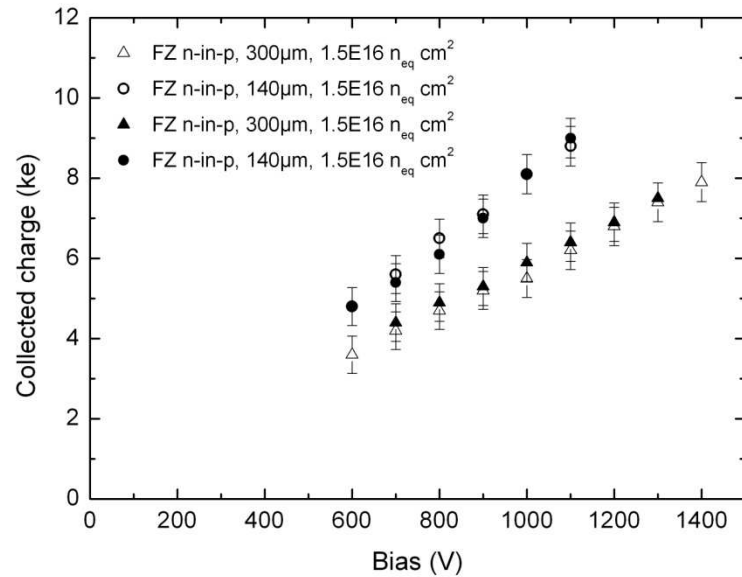


## Special effects: forward bias



The  $140 \mu\text{m}$  detectors could also be operated in forward bias (unlike the thicker devices). The CCE was surprisingly higher than at the corresponding reverse bias voltage, up to the maximum applicable bias before thermal runaway.

# CCE and currents after neutron irradiations after $1.5$ and $2 \times 10^{16} \text{ n}_{\text{eq}} \text{ cm}^{-2}$



# CONCLUSIONS

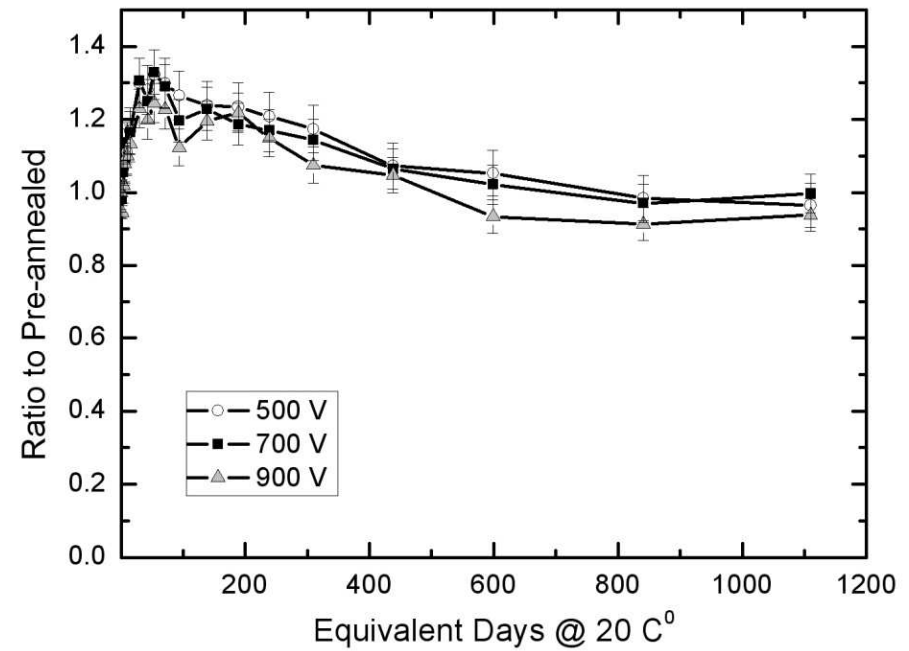
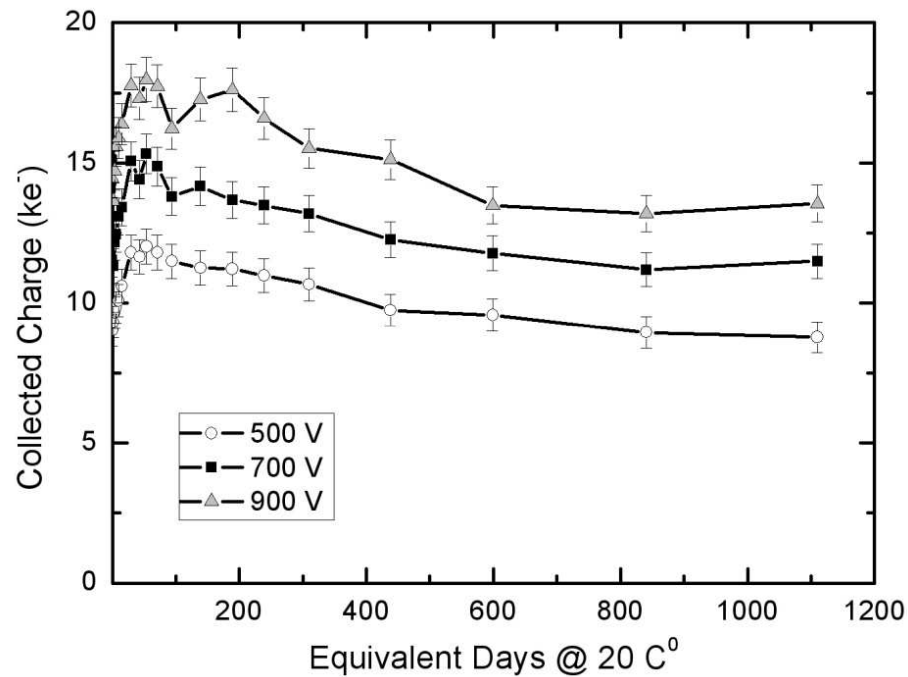
A much enhanced CCE is measured with proton and neutron irradiated diodes after extreme doses (equivalent to the anticipated doses for the innermost pixel layers of SLHC).

Even considering a non-linear dependence on fluence of the so called  $V_{FD}$ , and under the hypothesis of reduced trapping (non linear dependence on fluence? Field enhanced fast de-trapping?) the charge collected exceed the expectations. There is also evidence of collection of more charge than the ionised amount.

A controlled multiplication is taking place with heavily irradiated microstrip detectors. This opens the possibility of engineering the geometrical parameters of the detectors to increase the radiation tolerance of segmented detectors for applications where extreme hadron radiation damage is expected.

# Spare slides

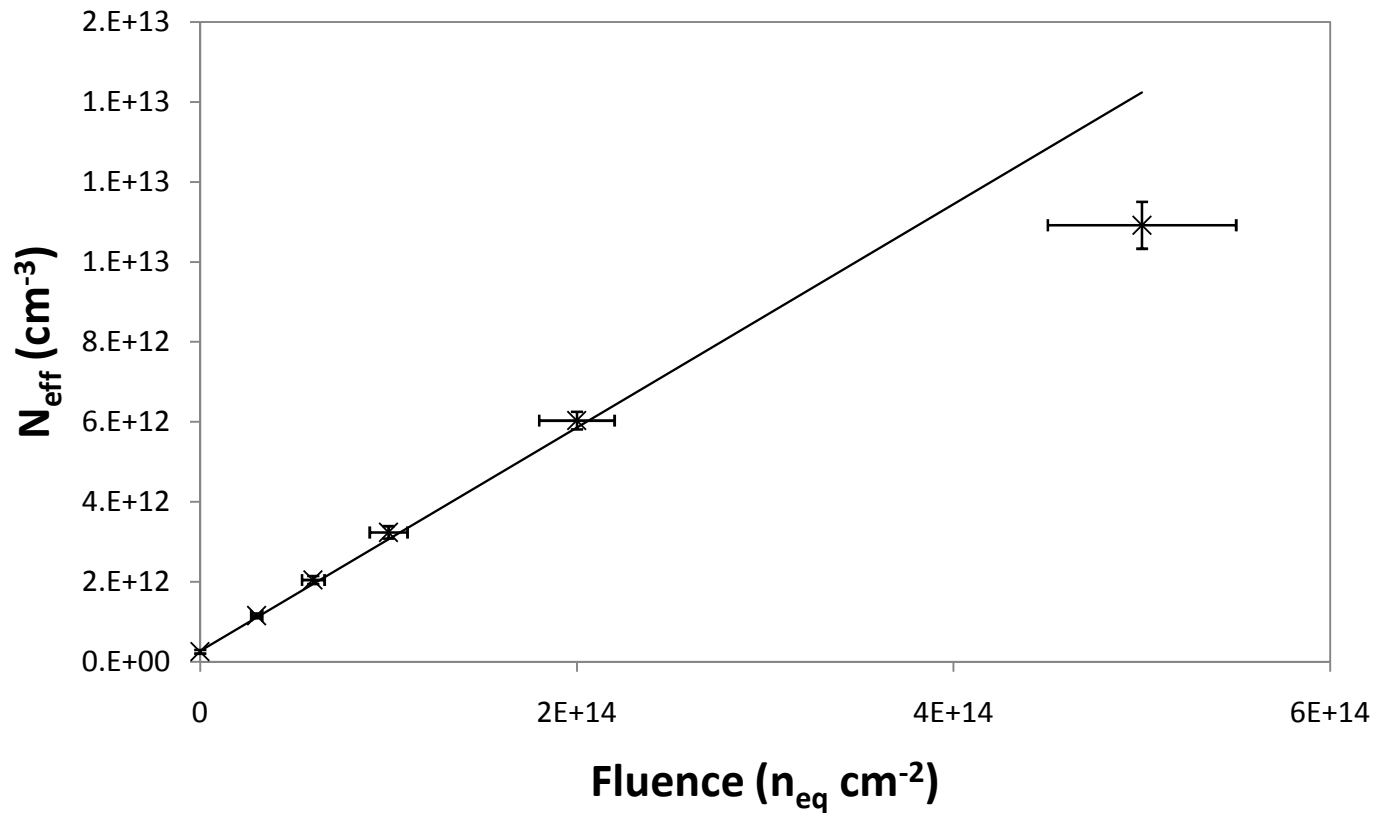
# A word on the annealing of the CCE (Neutrons)



# $N_{\text{eff}}$ vs $\phi$ measured with the CV characteristic of 300 $\mu\text{m}$ n-in-p Micron sensors after 80 min 60°C annealing time

$$N_{\text{eff}} = N_0(\exp(-c\phi)) - \beta\phi$$

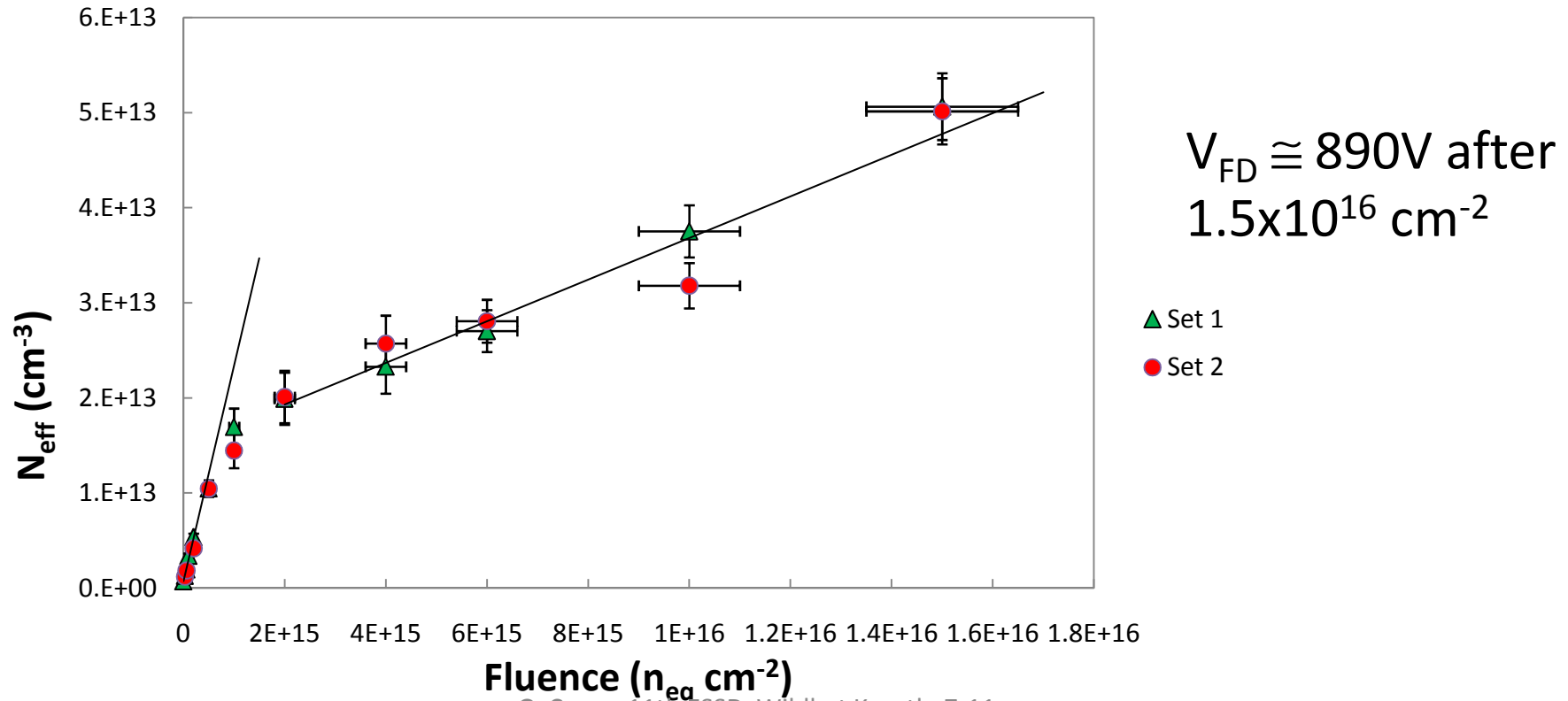
$$\beta = 0.028 \pm 0.002 \text{ cm}^{-1}$$



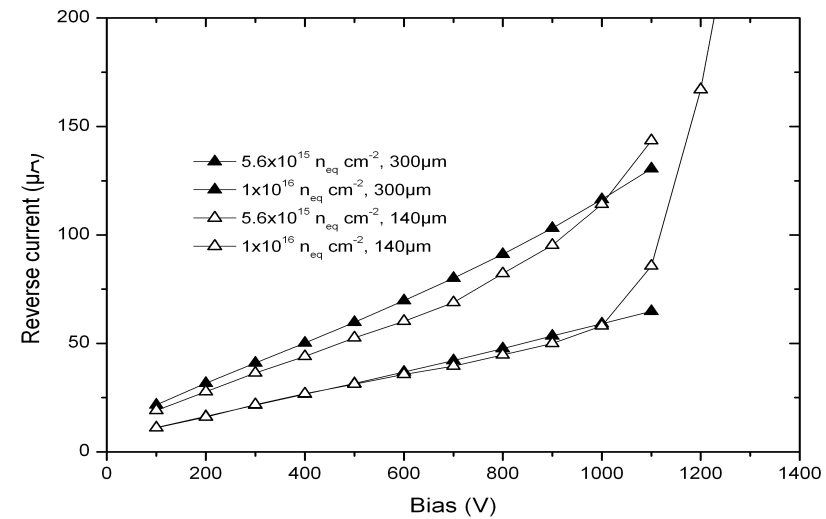
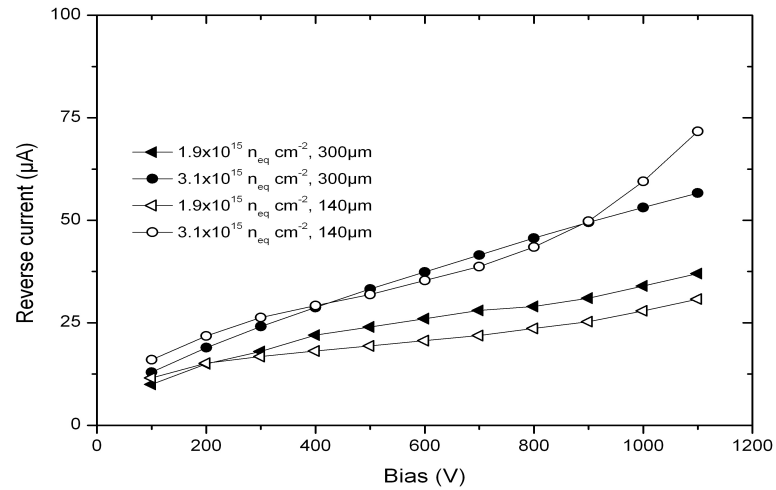
# $N_{\text{eff}}$ vs $\phi$ measured with the CV characteristic of 140 $\mu\text{m}$ n-in-p Micron sensors after 80 min 60°C annealing time

$\beta$  calculated for dose up to  $5 \times 10^{14} \text{ n}_{\text{eq}} \text{ cm}^{-2}$   
 $\beta = 0.023 \pm 0.002 \text{ cm}^{-1}$

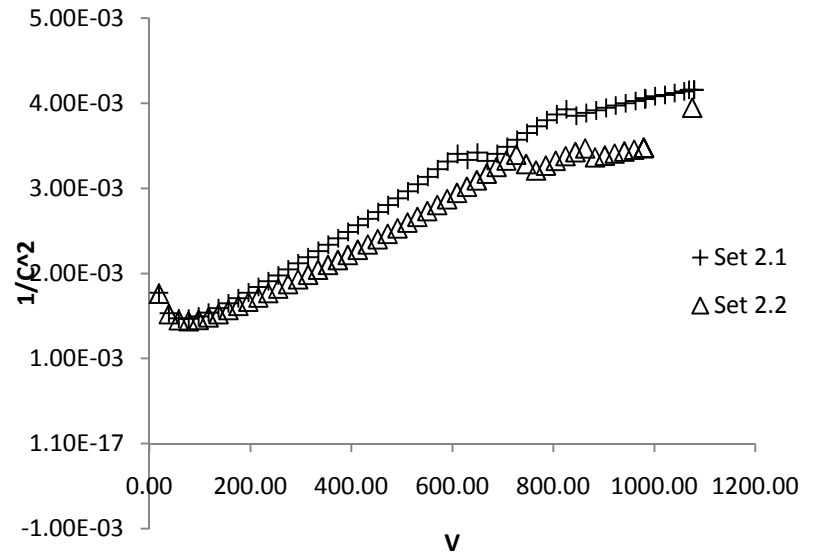
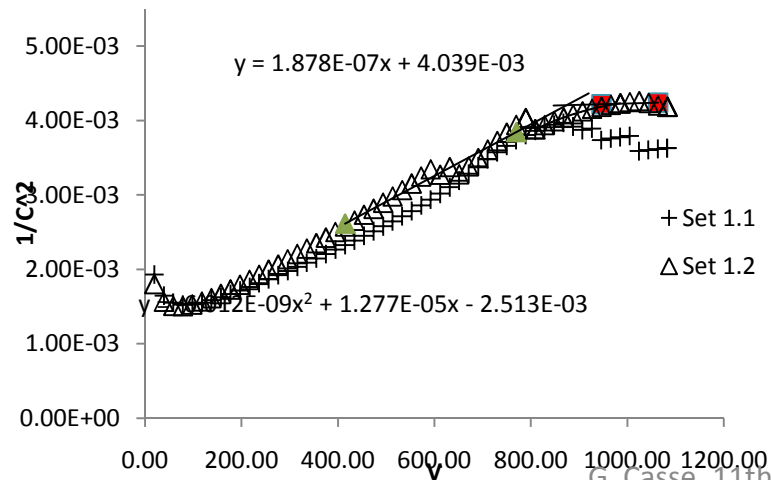
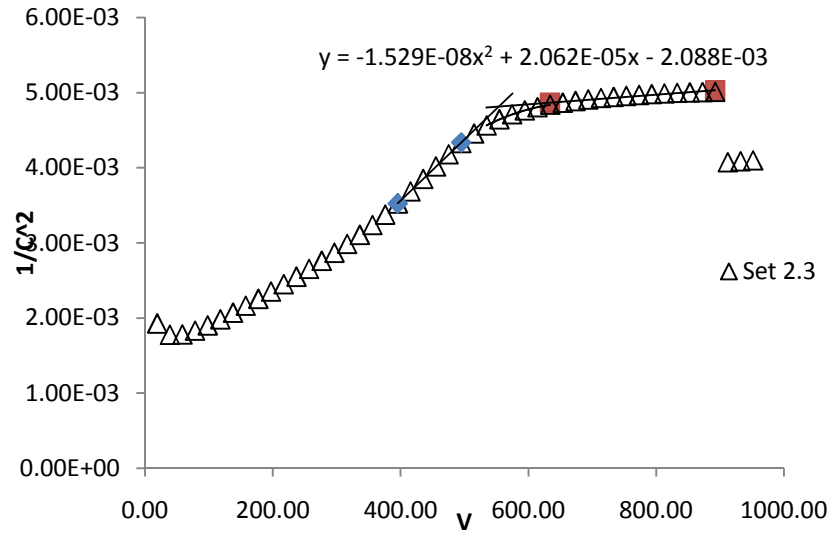
$\beta$  calculated for dose up to  $5 \times 10^{14} \text{ n}_{\text{eq}} \text{ cm}^{-2}$   
 $\beta = 0.0022 \pm 0.0003 \text{ cm}^{-1}$



# Reverse currents of 140 and 300 $\mu\text{m}$ n-in-p Micron sensors after various proton doses









$\alpha$  (300 $\mu\text{m}$  thick) =  $3.2 \times 10^{-17}$  A  $\text{cm}^{-2}$ .

$\alpha$  (140 $\mu\text{m}$  thick) =  $2.7 \times 10^{-17}$  A  $\text{cm}^{-2}$ .

