

# Collider Probes of Axion-like Particles

Andrea Thamm  
CERN

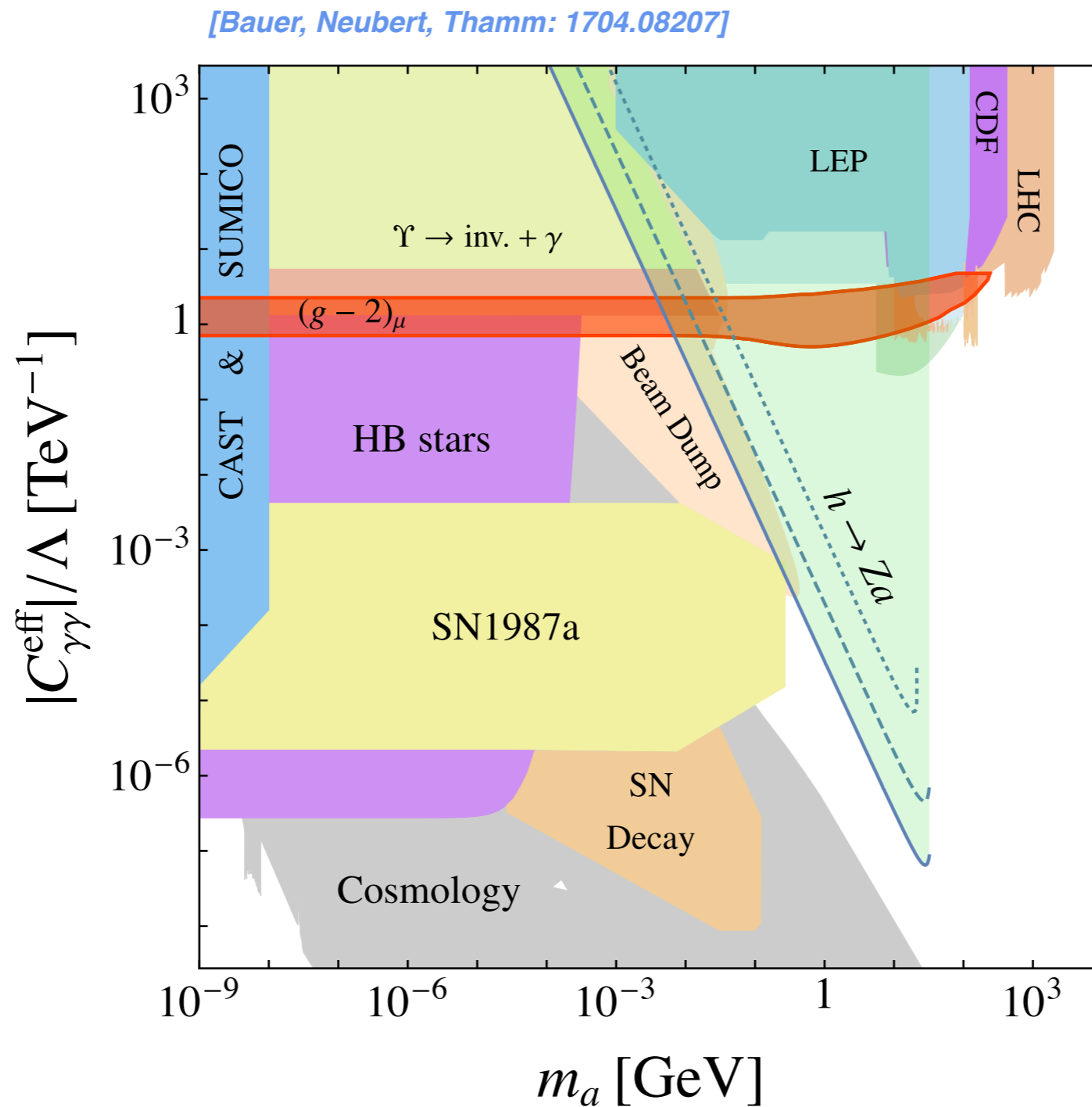
with Martin Bauer and Matthias Neubert

based on arXiv:1610.00009, 1704.08207, 1708.00443  
and work in progress

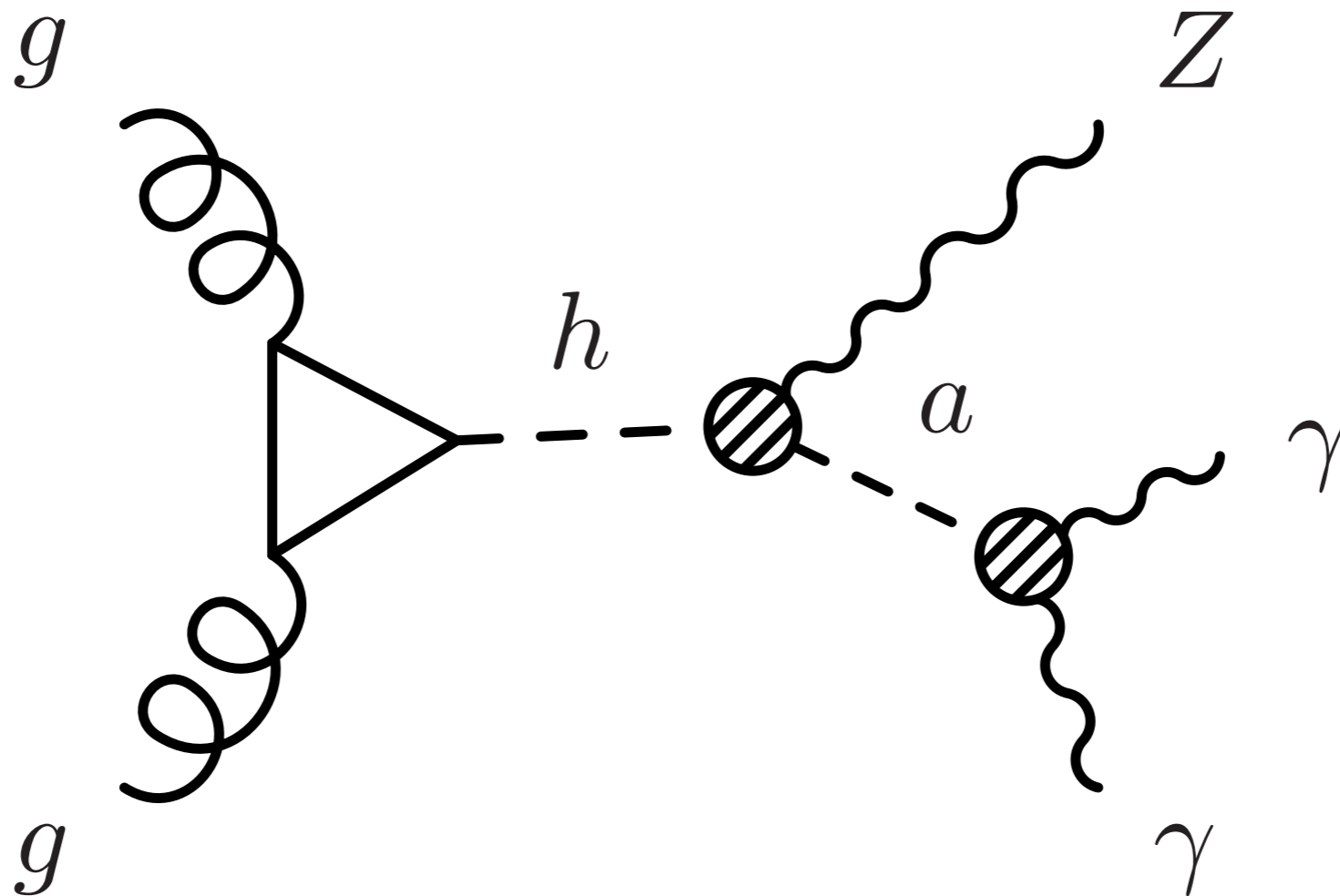


4 April 2018  
Liverpool

# Our main result

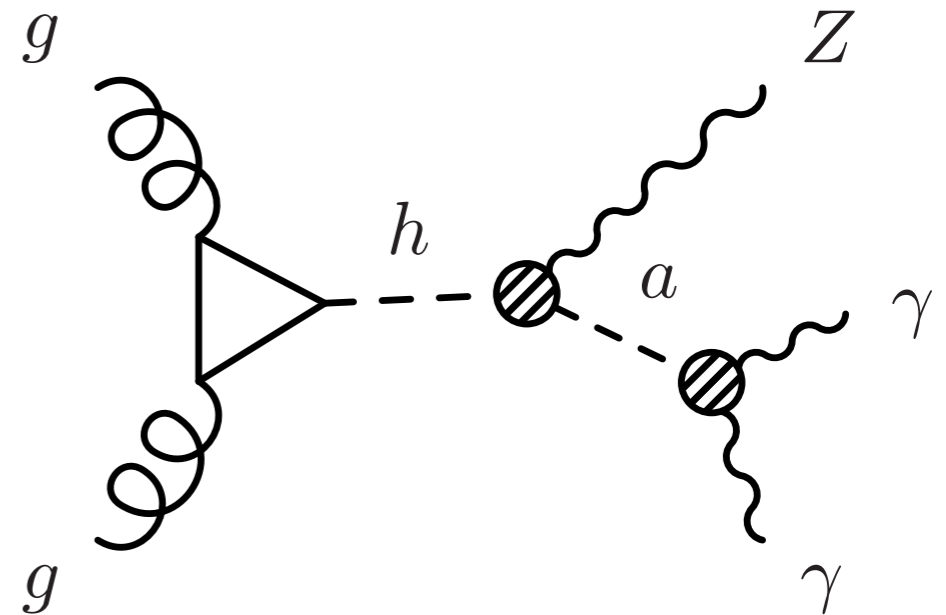


# Example process



# Outline

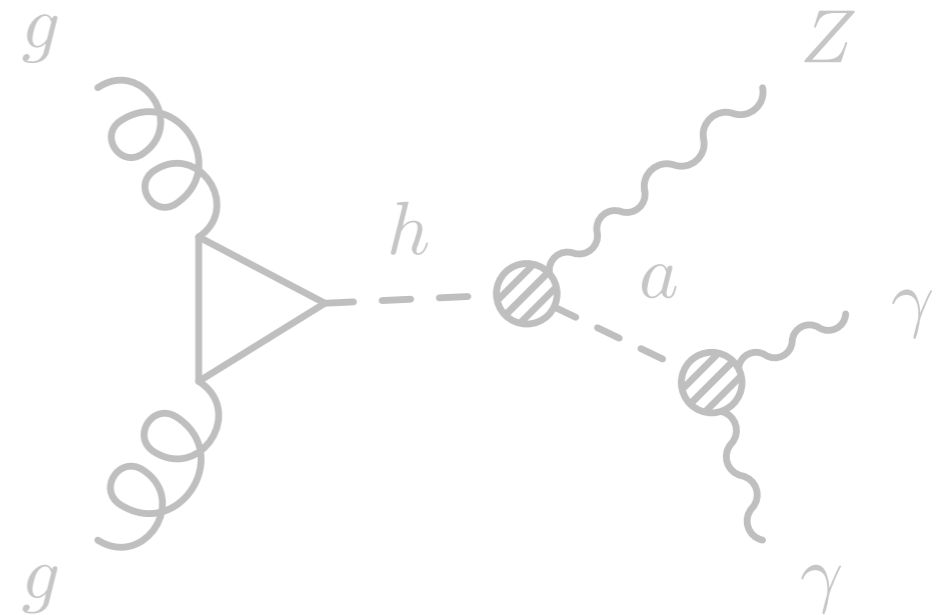
- Motivation
- ALPs and collider probes
  - ♦ Effective Lagrangian
  - ♦ Exotic Higgs decays
  - ♦ ALP Decays
  - ♦ Probing the ALP parameter space
  - ♦ Muon  $(g - 2)_\mu$
  - ♦ Future Colliders
- Conclusions and Outlook





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# Motivation

- Pseudo-scalars in many extensions of the SM
  - ♦ QCD axion - solution to strong CP-problem
  - ♦ Nambu-Goldstone bosons of a broken symmetry
  - ♦ mediators to the dark sector
  - ♦ explanations of various anomalies
- Good reason to study them!
- Large regions of parameter space already probed by many different experiments
- We add a region that can be probed through exotic Higgs decays in run 2 of LHC

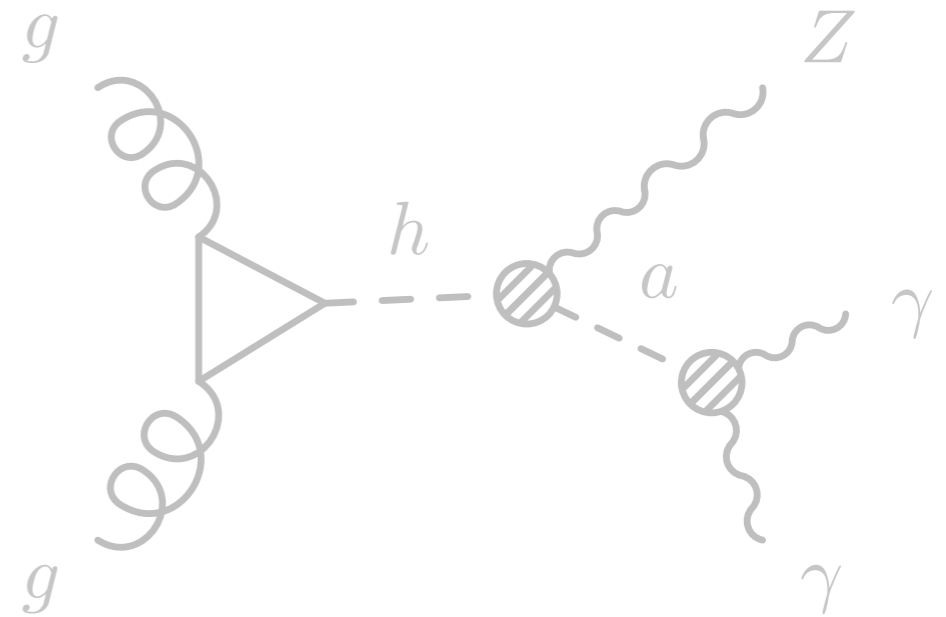
# Motivation

- Consider a singlet:  $(1,1,0)$  under  $SU(3)_C \times SU(2)_L \times U(1)_Y$
- Pseudoscalar and light
- Shift symmetry protects mass  $a \rightarrow a + c$
- Mass obtained through explicit soft breaking  
or non-perturbative dynamics

*[Weinberg: PRL 40 (1978) 223]  
[Wilczek: PRL 40 (1978) 279]*

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# Effective Lagrangian

- Interactions at dimension-5

[Georgi, Kaplan, Randall: Phys. Lett. 169 B (1986)]

$$\begin{aligned} \mathcal{L}_{\text{eff}}^{D \leq 5} = & \frac{1}{2} (\partial_\mu a)(\partial^\mu a) + \frac{1}{2} m_a^2 a^2 + \sum_f \frac{c_{ff}}{2} \frac{\partial^\mu a}{\Lambda} \bar{f} \gamma_\mu \gamma_5 f \\ & + g_s^2 C_{GG} \frac{a}{\Lambda} G_{\mu\nu}^A \tilde{G}^{\mu\nu, A} + g^2 C_{WW} \frac{a}{\Lambda} W_{\mu\nu}^A \tilde{W}^{\mu\nu, A} + g'^2 C_{BB} \frac{a}{\Lambda} B_{\mu\nu} \tilde{B}^{\mu\nu} \end{aligned}$$

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$$+ g_s^2 C_{GG} \frac{a}{\Lambda} G_{\mu\nu}^A \tilde{G}^{\mu\nu,A} + g^2 C_{WW} \frac{a}{\Lambda} W_{\mu\nu}^A \tilde{W}^{\mu\nu,A} + g'^2 C_{BB} \frac{a}{\Lambda} B_{\mu\nu} \tilde{B}^{\mu\nu}$$

- After EWSB

$$\mathcal{L}_{\text{eff}}^{D \leq 5} \ni e^2 C_{\gamma\gamma} \frac{a}{\Lambda} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{2e^2}{s_w c_w} C_{\gamma Z} \frac{a}{\Lambda} F_{\mu\nu} \tilde{Z}^{\mu\nu} + \frac{e^2}{s_w^2 c_w^2} C_{ZZ} \frac{a}{\Lambda} Z_{\mu\nu} \tilde{Z}^{\mu\nu}$$

$$C_{\gamma\gamma} = C_{WW} + C_{BB}, \quad C_{\gamma Z} = c_w^2 C_{WW} - s_w^2 C_{BB} \quad C_{ZZ} = c_w^4 C_{WW} + s_w^4 C_{BB}$$

# Effective Lagrangian

- Vanishes through equations of motion

$$\frac{(\partial^\mu a)}{\Lambda} (\phi^\dagger iD_\mu \phi + \text{h.c.})$$

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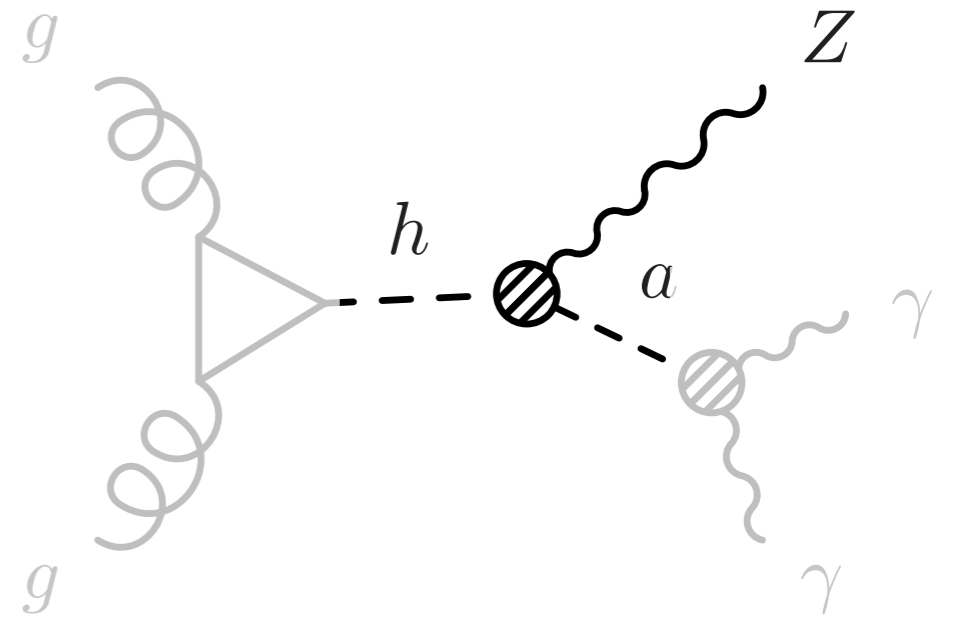
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[Bauer, Neubert, Thamm: 1607.01016]

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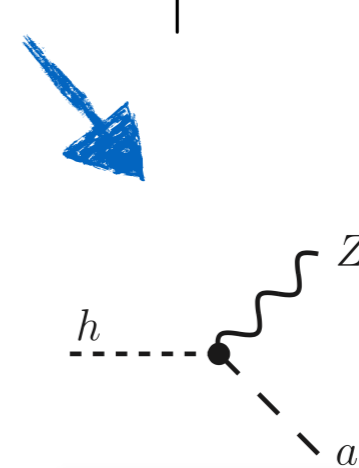
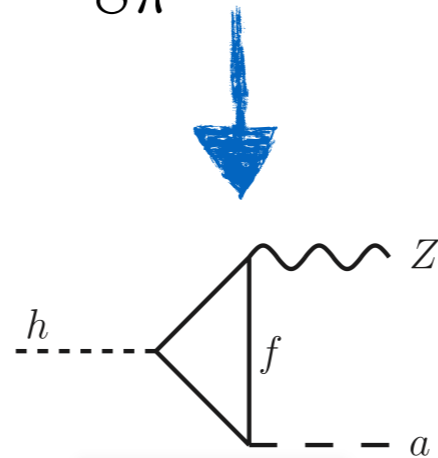
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# Exotic Higgs Decays $h \rightarrow Za$

- Contributions

$$\Gamma(h \rightarrow Za) = \frac{m_h^3}{16\pi\Lambda^2} \left| C_{Zh}^{(5)} - \frac{N_c y_t^2}{8\pi^2} T_3^t c_{tt} F + \frac{v^2}{2\Lambda^2} C_{Zh}^{(7)} \right|^2 \lambda^{3/2} \left( \frac{m_Z^2}{m_h^2}, \frac{m_a^2}{m_h^2} \right)$$



[Bauer, Neubert, Thamm:1610.00009]

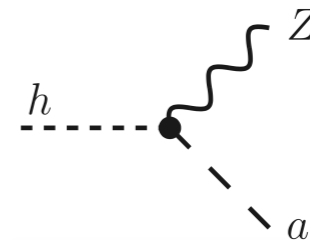
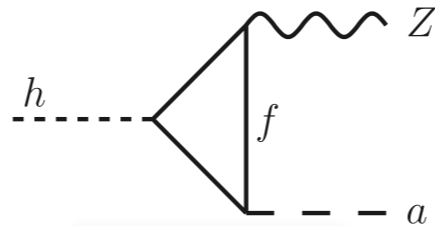
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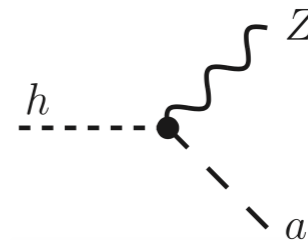
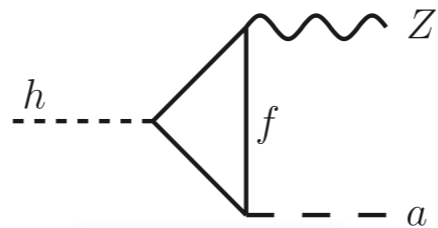
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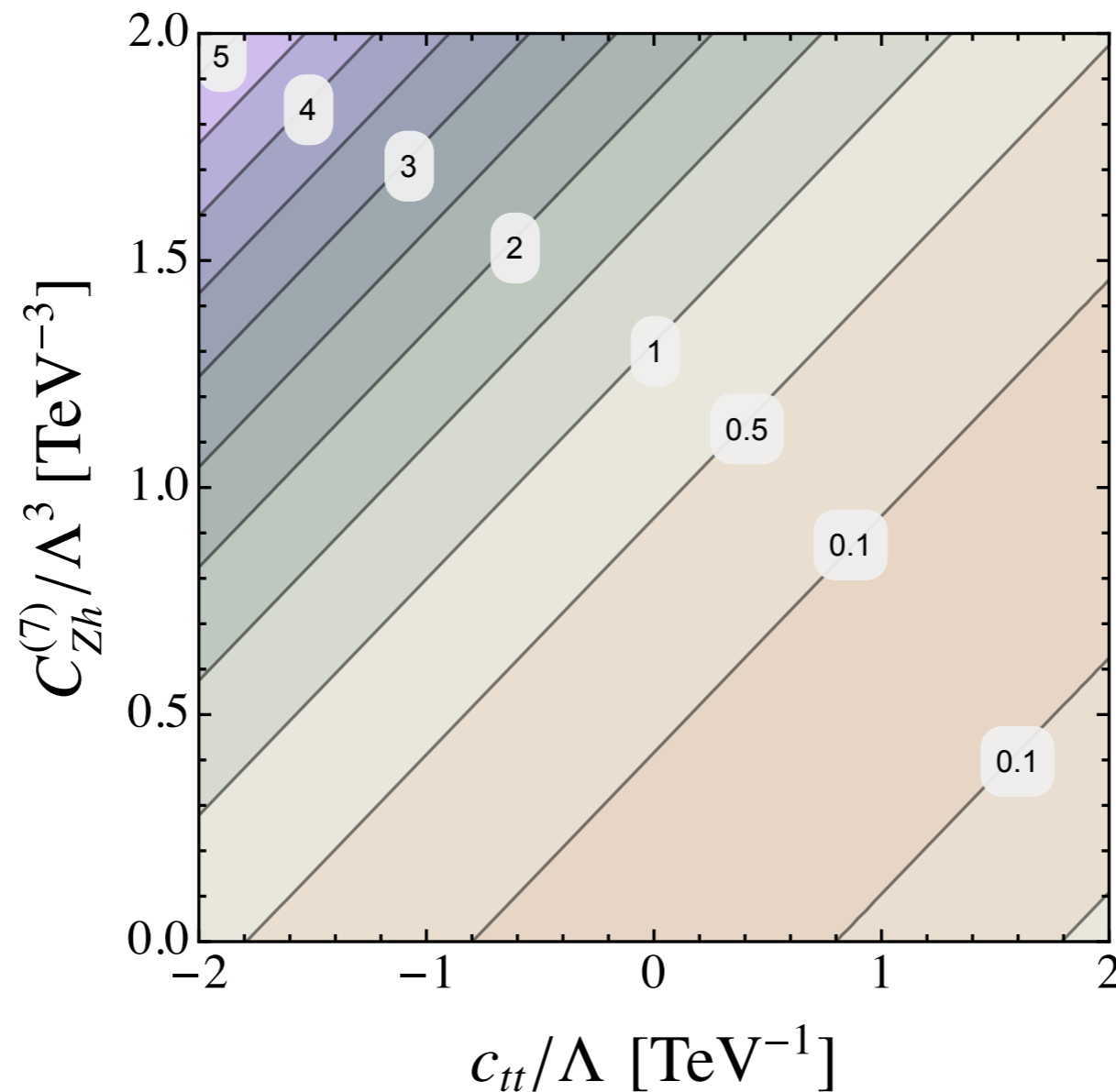
- Numerically

$$C_{Zh}^{\text{eff}} \approx C_{Zh}^{(5)} - 0.016 c_{tt} + 0.030 C_{Zh}^{(7)} \left[ \frac{1 \text{ TeV}}{\Lambda} \right]^2$$

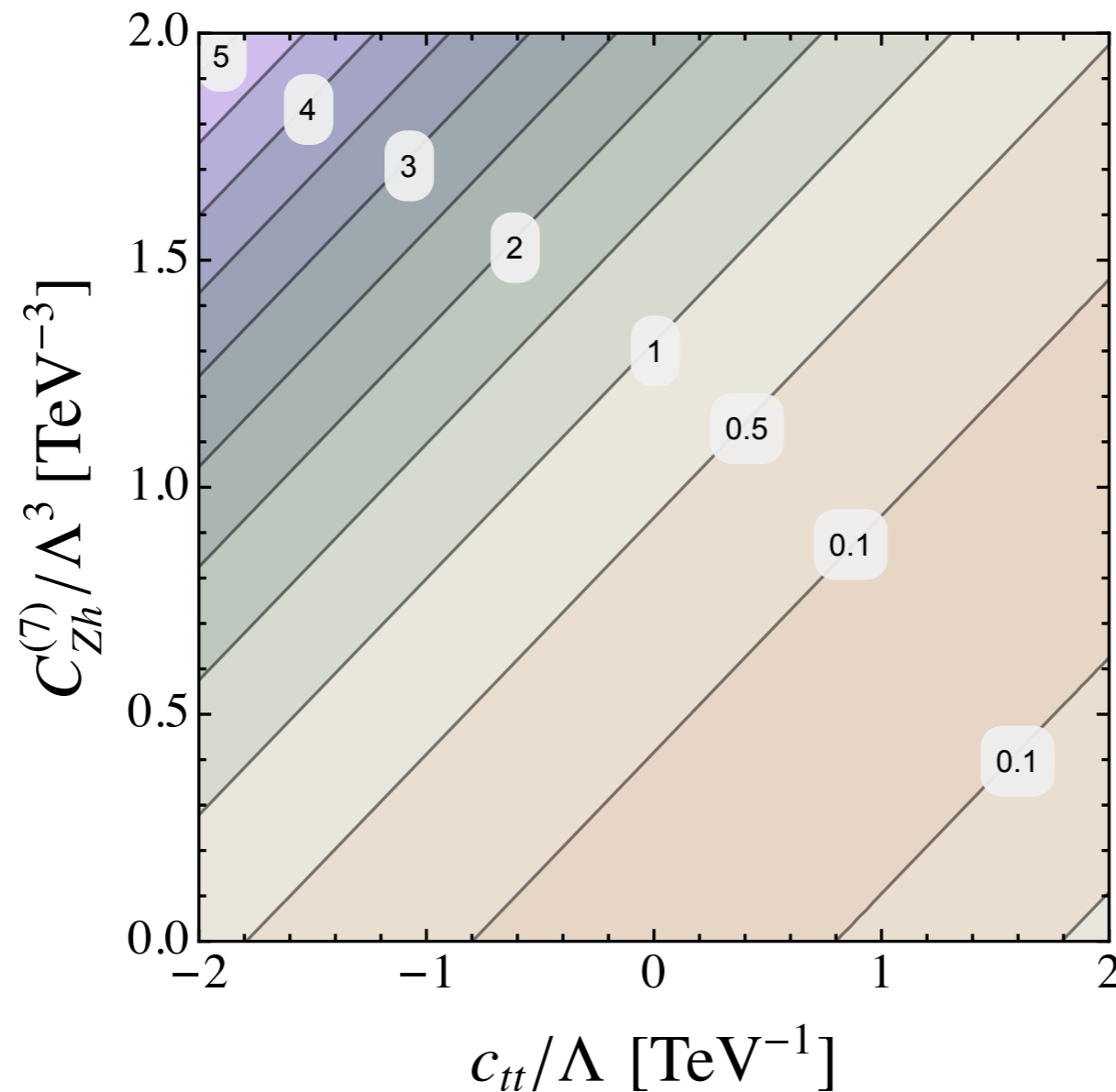
[Bauer, Neubert, Thamm:1610.00009]

# Exotic Higgs Decays $h \rightarrow Z a$

- Decay rate normalised to SM  $\Gamma(h \rightarrow Z\gamma)_{\text{SM}} = 6.32 \cdot 10^{-6} \text{ GeV}$



# Exotic Higgs Decays $h \rightarrow Za$



- This channel is a realistic target for discovery at LHC

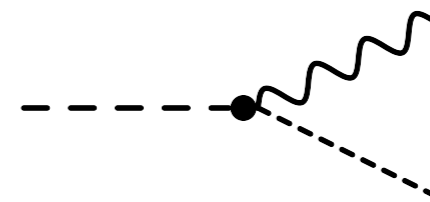
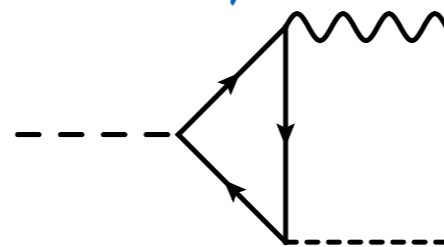


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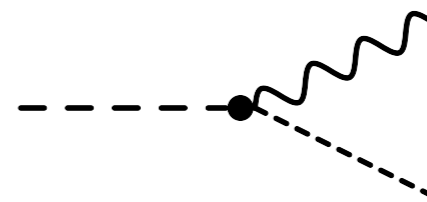
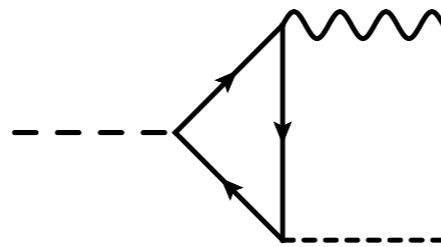
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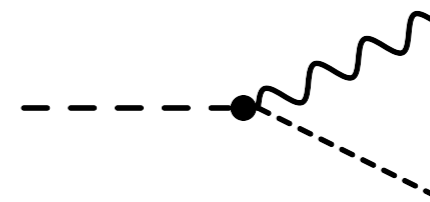
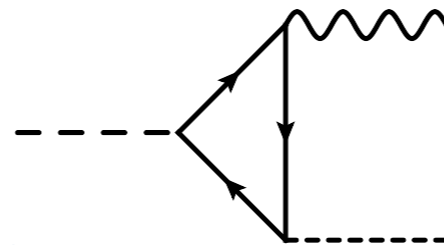


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Non-polynomial operator for models with new heavy particles whose mass arises from EWSB

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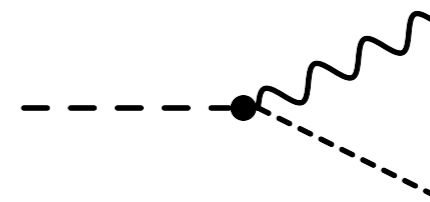
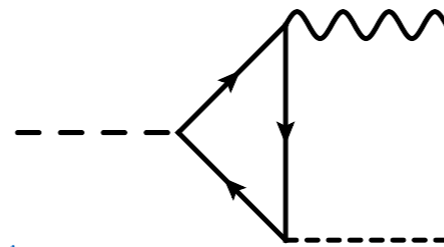
[\[Pierce, Thaler, Wang: 0609049\]](#)  
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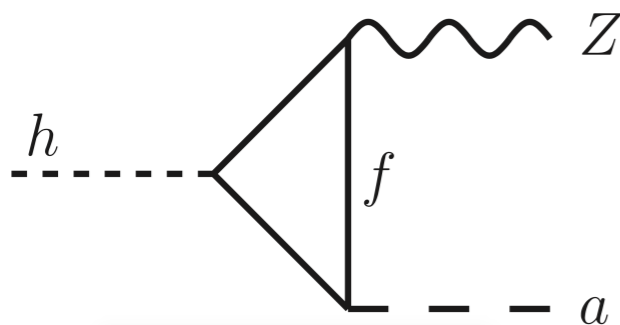
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- Non-polynomial operator

[Pierce, Thaler, Wang: 0609049]  
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$$\begin{aligned} \mathcal{L}_{\text{eff}}^{\text{non-pol}} &\ni \frac{C_{Zh}^{(5)}}{\Lambda} (\partial^\mu a) (\phi^\dagger iD_\mu \phi + \text{h.c.}) \ln \frac{\phi^\dagger \phi}{\mu^2} + \dots, \\ &= -\frac{C_{Zh}^{(5)}}{\Lambda} a (\phi^\dagger iD_\mu \phi + \text{h.c.}) \frac{\partial^\mu (\phi^\dagger \phi)}{\phi^\dagger \phi} + \dots \\ &\rightarrow -\frac{C_{Zh}^{(5)}}{\Lambda} \frac{g}{c_w} a Z_\mu (v + h) \partial^\mu h \end{aligned}$$

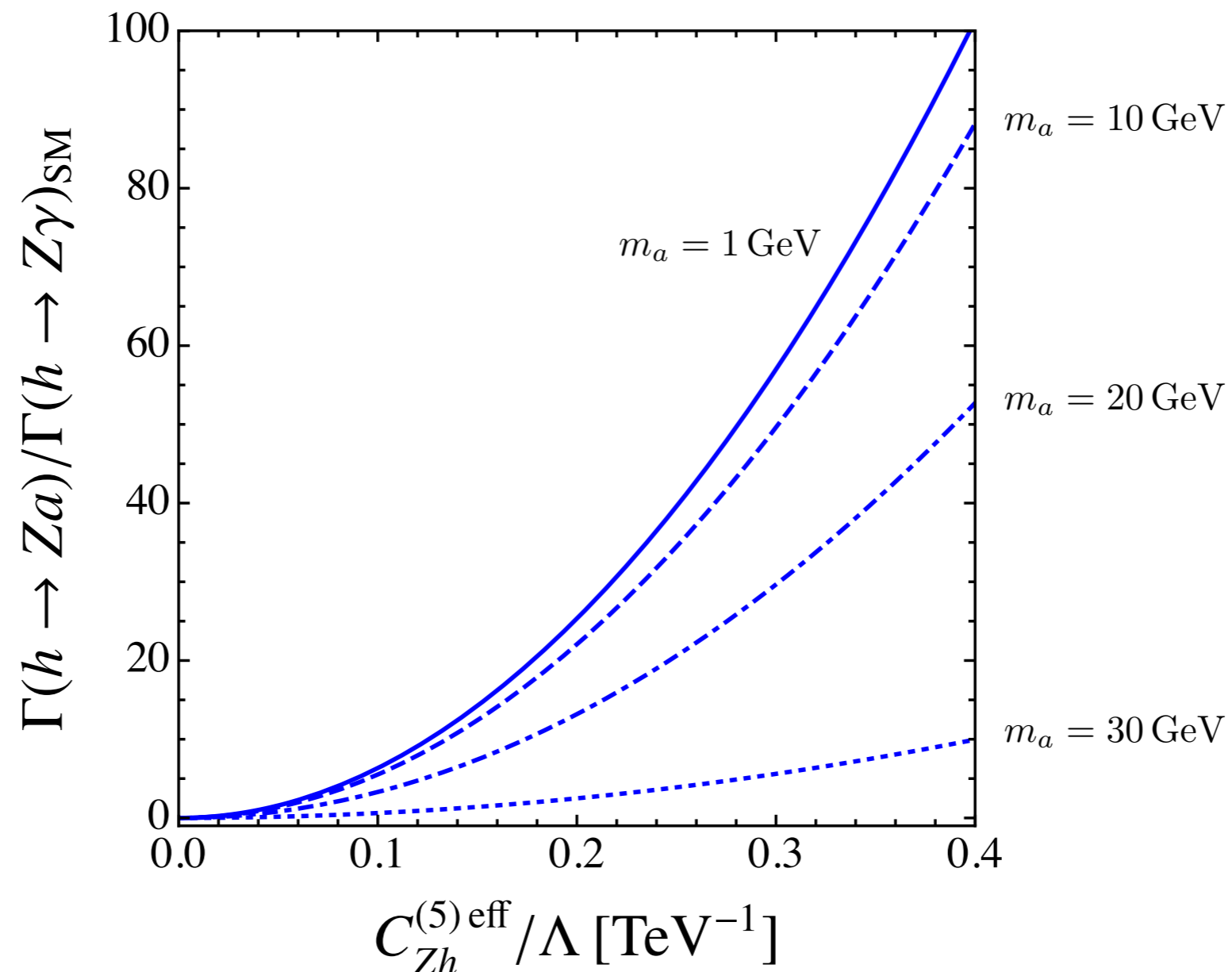


$$F = \int_0^1 d[xyz] \frac{2m_t^2 - xm_h^2 - zm_Z^2}{m_t^2 - xym_h^2 - yzm_Z^2 - xzm_a^2}$$

$$C_{Zh}^{(5)} = -\frac{N_c y_t^2}{8\pi^2} T_3^t \tilde{c}_{tt} F$$

# Exotic Higgs Decays $h \rightarrow Za$

- Enhanced rates for this process



# Exotic Higgs Decays $h \rightarrow Za$

- Current upper limit  $\text{Br}(h \rightarrow \text{BSM}) < 0.34$

*[ATLAS and CMS:1606.02266]*

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- Interesting final states

- ◆  $h \rightarrow Za \rightarrow Z\gamma\gamma$

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- All these modes can be reconstructed at run II

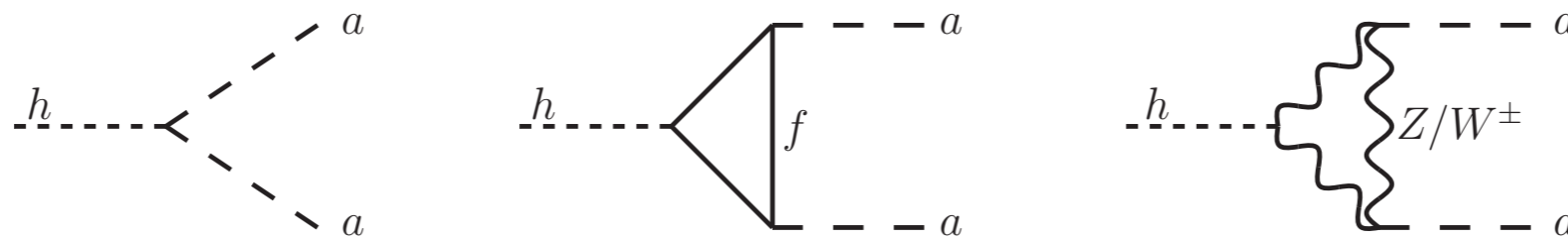
# Exotic Higgs Decays $h \rightarrow aa$

- Dim-6 Higgs portal and loop diagrams

[Dobrescu, Landsberg, Matchev: 0005308]

[Dobrescu, Matchev: 0008192]

[Chang, Fox, Weiner: 0608310]



$$C_{ah}^{\text{eff}} = C_{ah}(\mu) + \frac{N_c y_t^2}{4\pi^2} c_{tt}^2 \left[ \ln \frac{\mu^2}{m_t^2} - g_1(\tau_{t/h}) \right] - \frac{3\alpha}{2\pi s_w^2} (g^2 C_{WW})^2 \left[ \ln \frac{\mu^2}{m_W^2} + \delta_1 - g_2(\tau_{W/h}) \right]$$

$$- \frac{3\alpha}{4\pi s_w^2 c_w^2} \left( \frac{g^2}{c_w^2} C_{ZZ} \right)^2 \left[ \ln \frac{\mu^2}{m_Z^2} + \delta_1 - g_2(\tau_{Z/h}) \right]$$

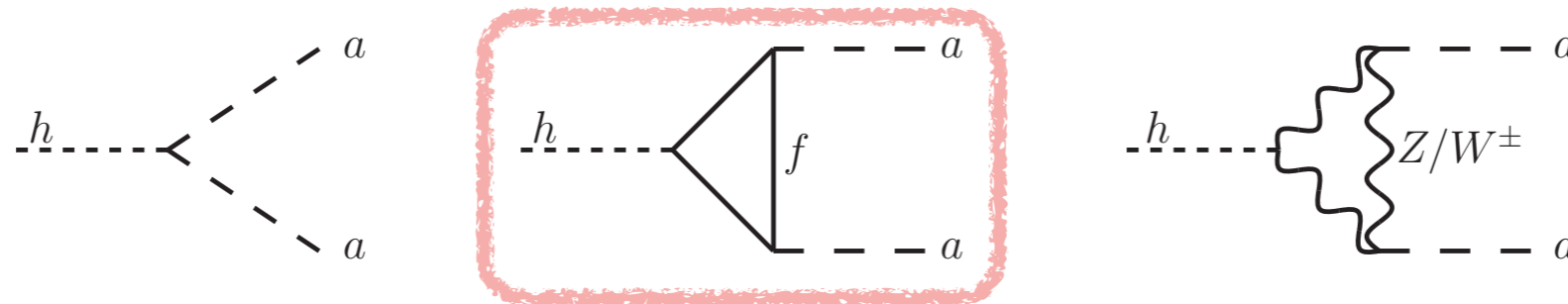
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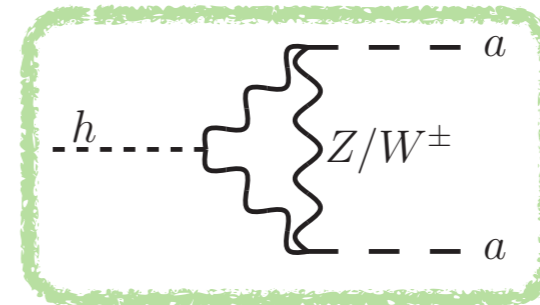
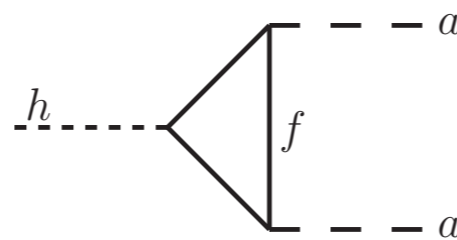
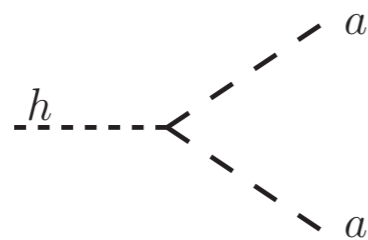
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[Dobrescu, Landsberg, Matchev: 0005308]

[Dobrescu, Matchev: 0008192]

[Chang, Fox, Weiner: 0608310]



$$C_{ah}^{\text{eff}} = C_{ah}(\mu) + \frac{N_c y_t^2}{4\pi^2} c_{tt}^2 \left[ \ln \frac{\mu^2}{m_t^2} - g_1(\tau_{t/h}) \right] - \frac{3\alpha}{2\pi s_w^2} (g^2 C_{WW})^2 \left[ \ln \frac{\mu^2}{m_W^2} + \delta_1 - g_2(\tau_{W/h}) \right]$$

$$- \frac{3\alpha}{4\pi s_w^2 c_w^2} \left( \frac{g^2}{c_w^2} C_{ZZ} \right)^2 \left[ \ln \frac{\mu^2}{m_Z^2} + \delta_1 - g_2(\tau_{Z/h}) \right]$$

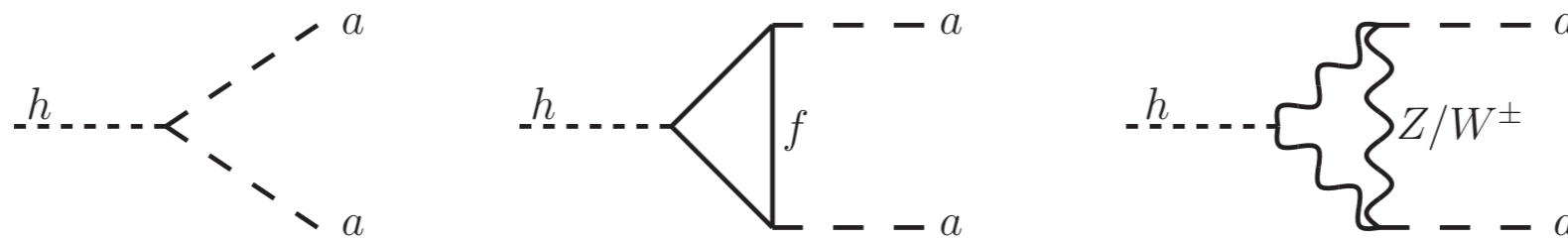
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$$C_{ah}^{\text{eff}} \approx C_{ah}(\Lambda) + 0.173 c_{tt}^2 - 0.0025 (C_{WW}^2 + C_{ZZ}^2)$$



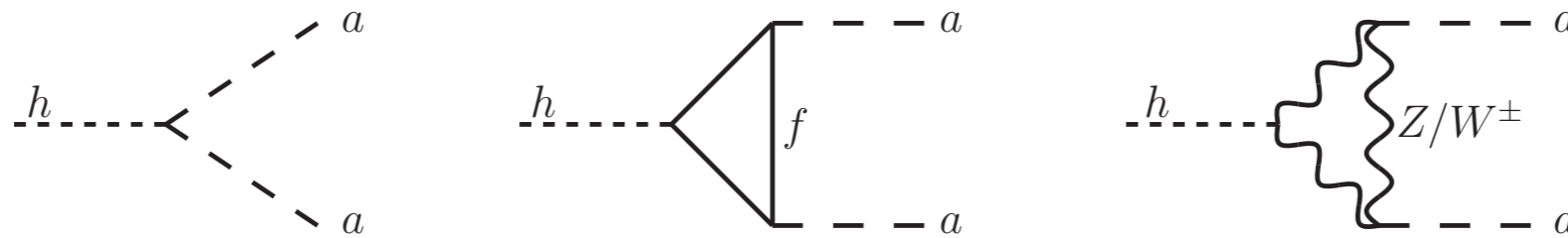
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$$\Gamma(h \rightarrow aa) = \frac{v^2 m_h^3}{32\pi \Lambda^4} |C_{ah}^{\text{eff}}|^2 \left( 1 - \frac{2m_a^2}{m_h^2} \right)^2 \sqrt{1 - \frac{4m_a^2}{m_h^2}}$$

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- Current upper limit  $\text{Br}(h \rightarrow \text{BSM}) < 0.34$

[ATLAS and CMS:1606.02266]

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- For  $\text{Br}(h \rightarrow aa) = 0.1$  need  $|C_{ah}|/\Lambda^2 \approx 0.62 \text{ TeV}^{-2}$
- From top-loop only:  $\text{Br}(h \rightarrow aa) = 0.01$  for  $|c_{tt}|/\Lambda \approx 1.04 \text{ TeV}^{-1}$

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- Interesting final states

◆  $h \rightarrow aa \rightarrow \gamma\gamma\gamma\gamma$

◆  $h \rightarrow aa \rightarrow 4\text{jets}$

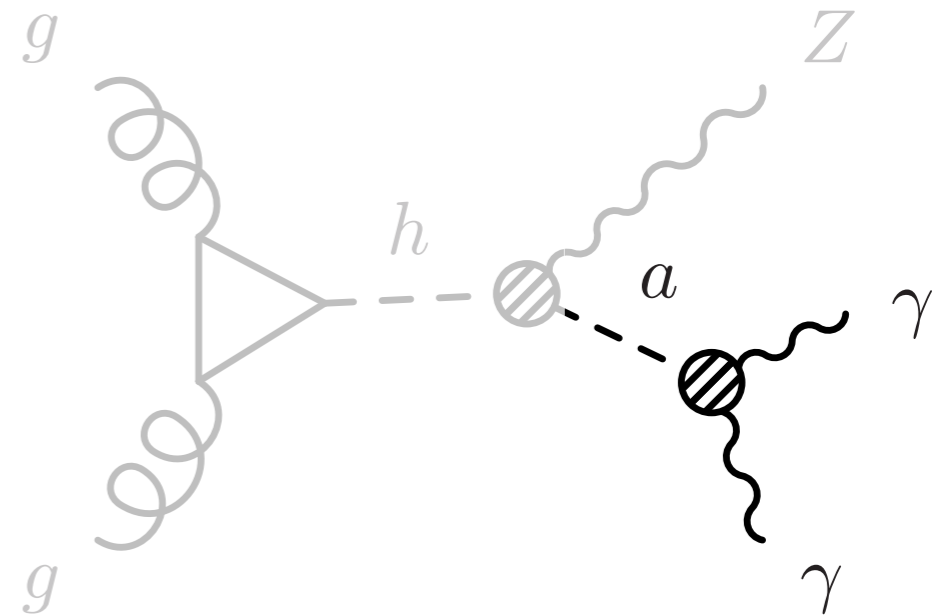
◆  $h \rightarrow aa \rightarrow l^+l^-l^+l^-$

◆  $h \rightarrow aa \rightarrow \text{invisible}$

- All these modes can be reconstructed at run II

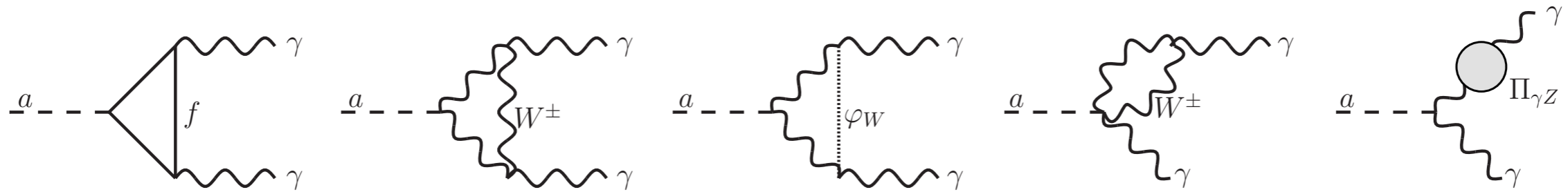
# Outline

- Motivation
- ALPs and collider probes
  - ♦ Effective Lagrangian
  - ♦ Exotic Higgs decays
  - ♦ ALP Decays
  - ♦ Probing the ALP parameter space
  - ♦ Muon  $(g - 2)_\mu$
  - ♦ Future Colliders
- Conclusions and Outlook



# ALP decays into photons

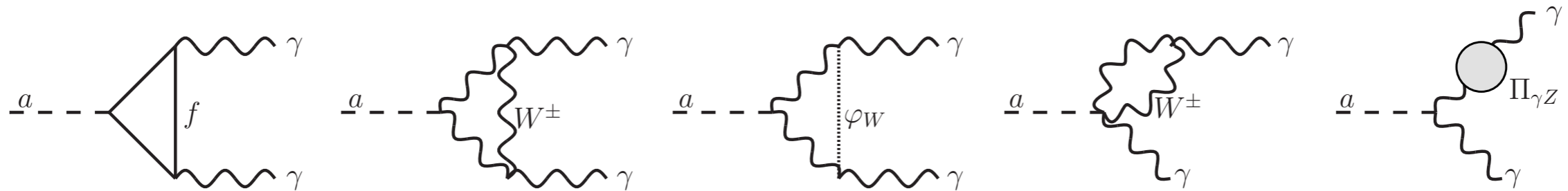
- Often considered as the dominant decay mode





# ALP decays into photons

- Often considered as the dominant decay mode



$$\Gamma(a \rightarrow \gamma\gamma) = \frac{4\pi\alpha^2 m_a^3}{\Lambda^2} \left| C_{\gamma\gamma} + \sum_f \frac{N_c^f Q_f^2}{16\pi^2} c_{ff} B_1(\tau_f) + \frac{2\alpha}{\pi} \frac{C_{WW}}{s_w^2} B_2(\tau_W) \right|^2$$

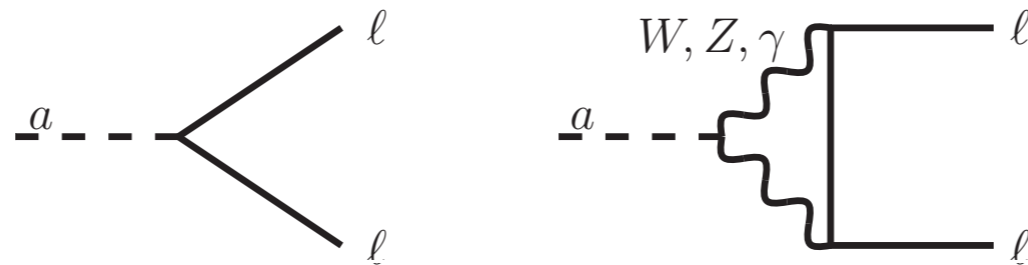
$$\equiv \frac{4\pi\alpha^2 m_a^3}{\Lambda^2} |C_{\gamma\gamma}^{\text{eff}}|^2$$

$$\tau_i \equiv 4m_i^2/m_a^2$$

- Only mode for  $m_a < 2m_e$

# ALP decays into leptons

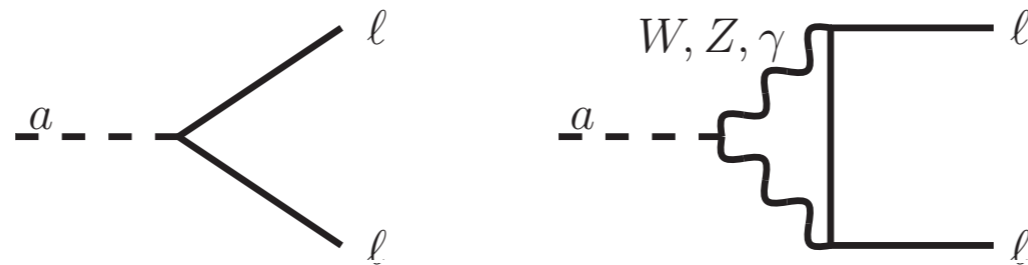
- For  $m_a > 2m_e$



$$\begin{aligned}
 c_{\ell\ell}^{\text{eff}} = & c_{\ell\ell}(\mu) [1 + \mathcal{O}(\alpha)] - 12Q_\ell^2 \alpha^2 C_{\gamma\gamma} \left[ \ln \frac{\mu^2}{m_\ell^2} + \delta_1 + g(\tau_\ell) \right] \\
 & - \frac{3\alpha^2}{s_w^4} C_{WW} \left( \ln \frac{\mu^2}{m_W^2} + \delta_1 + \frac{1}{2} \right) - \frac{12\alpha^2}{s_w^2 c_w^2} C_{\gamma Z} Q_\ell (T_3^\ell - 2Q_\ell s_w^2) \left( \ln \frac{\mu^2}{m_Z^2} + \delta_1 + \frac{3}{2} \right) \\
 & - \frac{12\alpha^2}{s_w^4 c_w^4} C_{ZZ} \left( Q_\ell^2 s_w^4 - T_3^\ell Q_\ell s_w^2 + \frac{1}{8} \right) \left( \ln \frac{\mu^2}{m_Z^2} + \delta_1 + \frac{1}{2} \right).
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 \end{aligned}$$

$$\Gamma(a \rightarrow \ell^+ \ell^-) = \frac{m_a m_\ell^2}{8\pi \Lambda^2} |c_{\ell\ell}^{\text{eff}}|^2 \sqrt{1 - \frac{4m_\ell^2}{m_a^2}}$$

# ALP decays into hadrons

- Decays into gluons and quarks
- For  $m_a > 2m_\pi$
- Can be computed only in perturbative regime for  $m_a \gg \Lambda_{\text{QCD}}$

$$\Gamma(a \rightarrow \text{hadrons}) = \frac{32\pi \alpha_s^2(m_a) m_a^3}{\Lambda^2} \left[ 1 + \left( \frac{97}{4} - \frac{7n_q}{6} \right) \frac{\alpha_s(m_a)}{\pi} \right] \left| C_{GG} + \sum_{q=1}^{n_q} \frac{c_{qq}}{32\pi^2} \right|^2$$

[Spira, Djouadi, Graudenz, Zerwas: 9504378]

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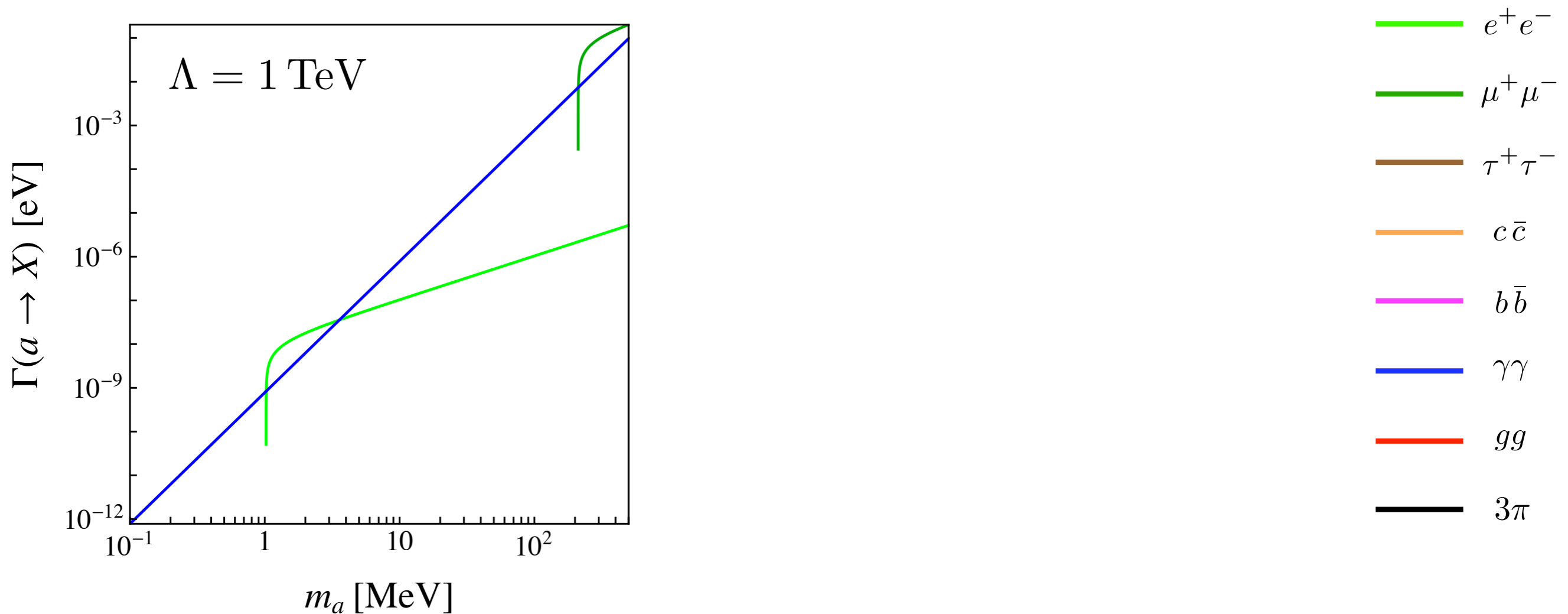
[Spira, Djouadi, Graudenz, Zerwas: 9504378]

- Decays into heavy quarks

$$\Gamma(a \rightarrow Q\bar{Q}) = \frac{3m_a \bar{m}_Q^2(m_a)}{8\pi\Lambda^2} |c_{QQ}^{\text{eff}}|^2 \sqrt{1 - \frac{4m_Q^2}{m_a^2}}$$

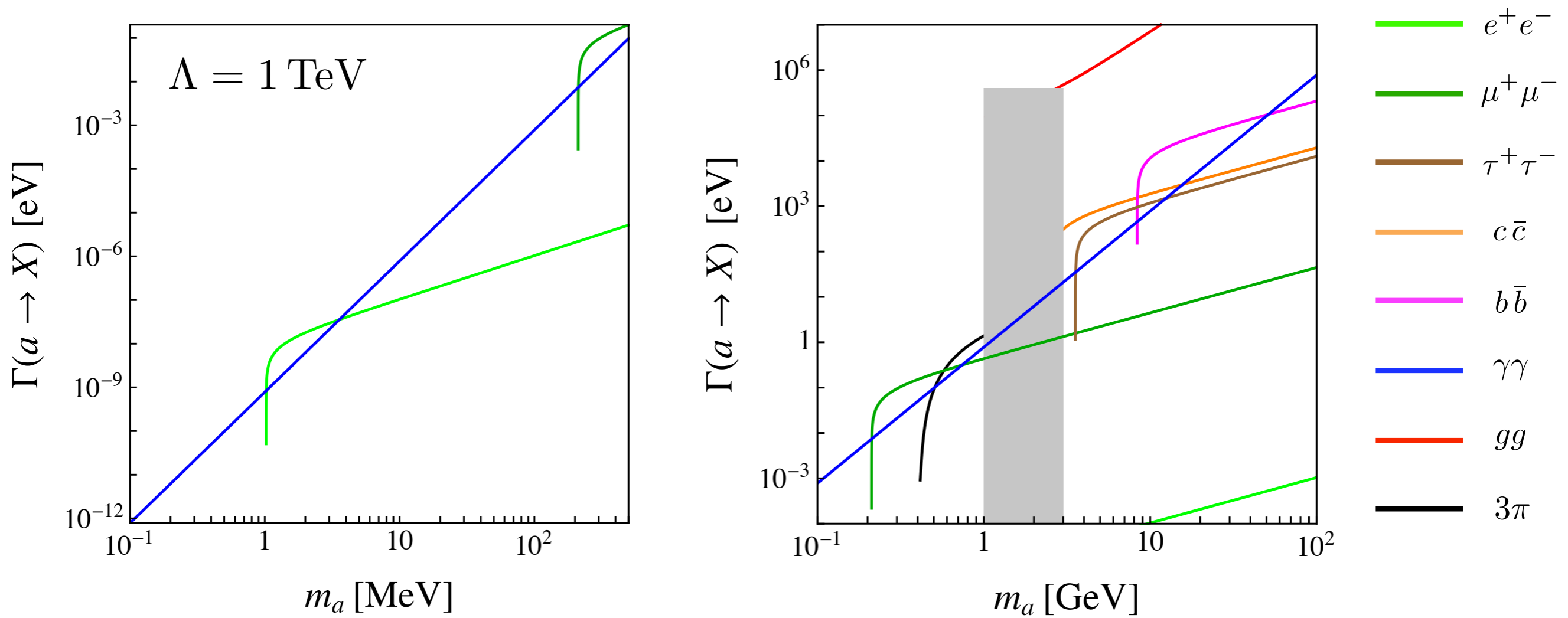
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- Assuming effective Wilson coefficients to be 1



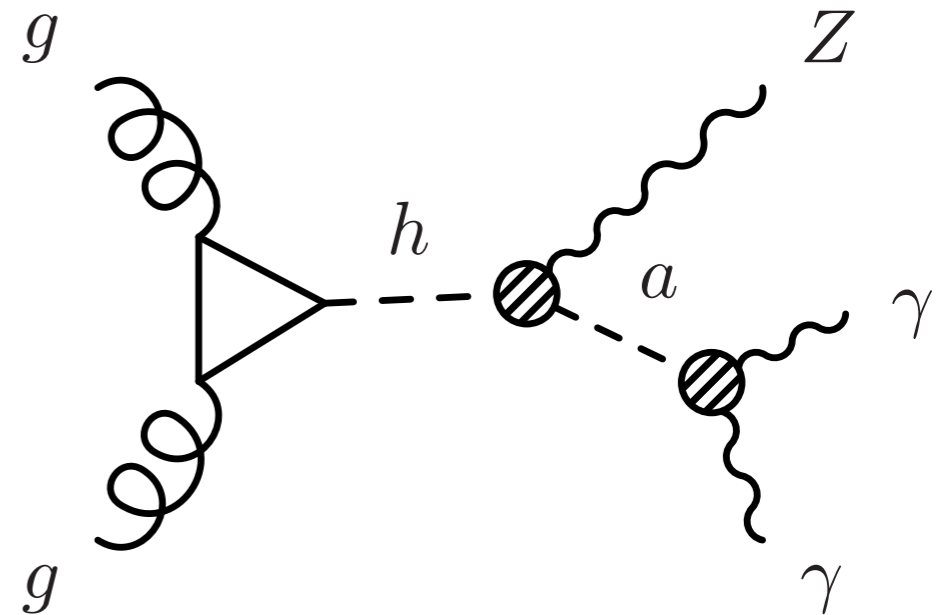
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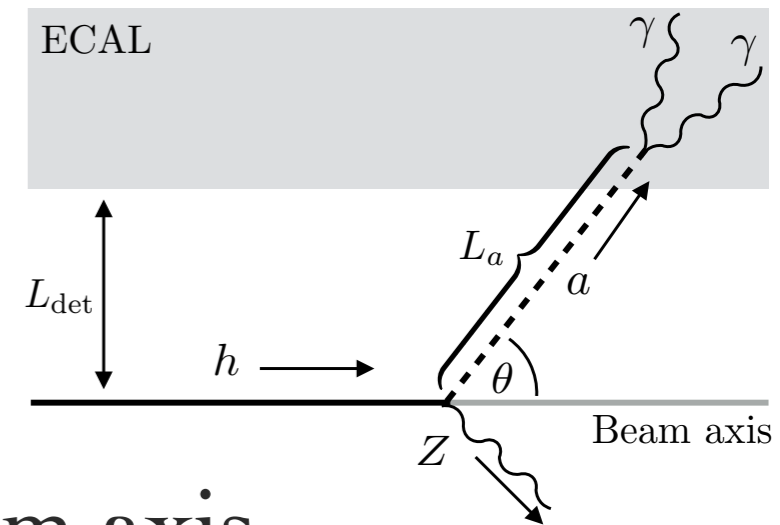
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  - ♦ Electroweak precision test
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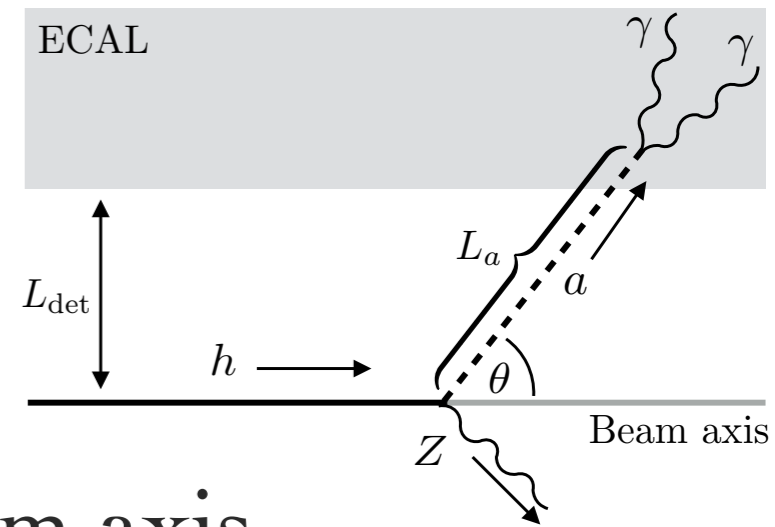
# Detecting ALPs



- Average decay length perpendicular to beam axis

$$L_a^\perp(\theta) = \sin \theta \frac{\beta_a \gamma_a}{\Gamma_a}$$

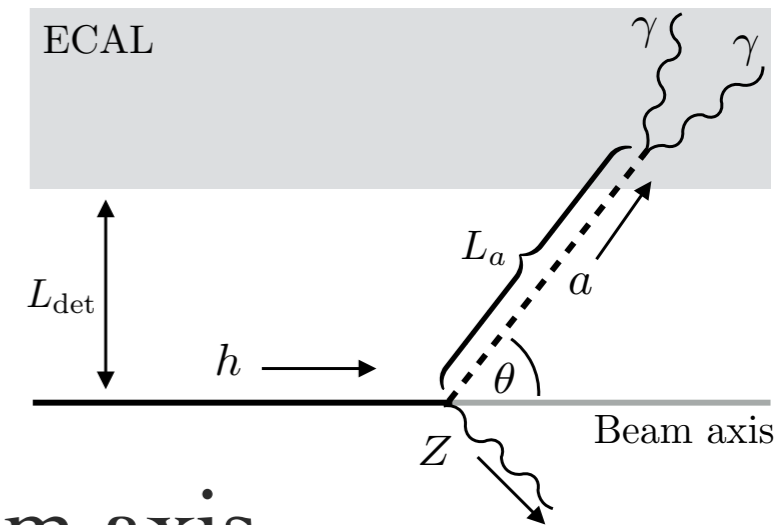
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$$\begin{aligned}
 L_a^\perp(\theta) &= \sin \theta \frac{\beta_a \gamma_a}{\Gamma_a} \\
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# Detecting ALPs



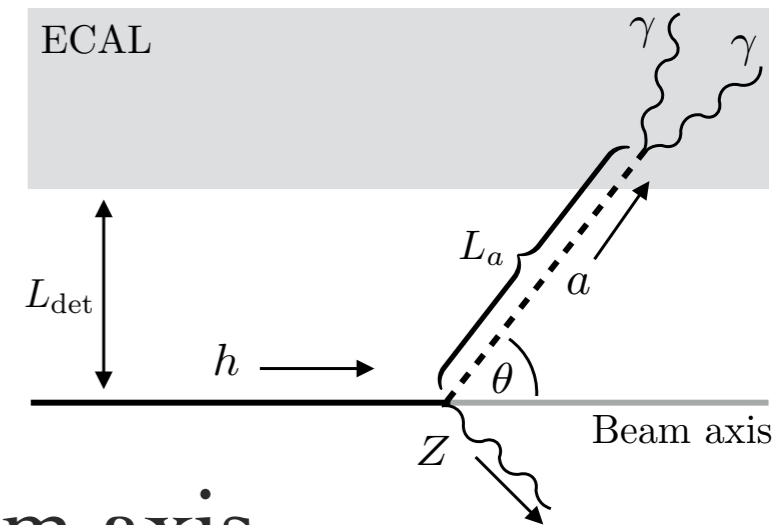
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- Fraction of ALPs decaying before travelling a certain distance

$$f_{\text{det}} = \int_0^{\pi/2} d\theta \sin \theta \left( 1 - e^{-L_{\text{det}}/L_a^\perp(\theta)} \right)$$

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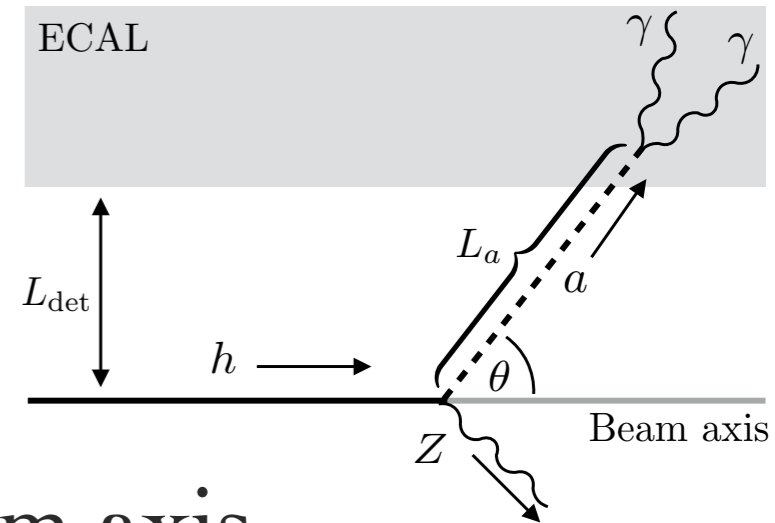
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Decay into photons  
before EM calorimeter

$$L_{\text{det}} = 1.5 \text{ m}$$

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Decay into photons  
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$$L_{\text{det}} = 1.5 \text{ m}$$

Decay into electrons  
before inner tracker

$$L_{\text{det}} = 2 \text{ cm}$$

# Detecting ALPs

- Effective branching ratios

$$\text{Br}(h \rightarrow Za \rightarrow \ell^+ \ell^- X \bar{X})|_{\text{eff}} = \text{Br}(h \rightarrow Za) \times \text{Br}(a \rightarrow X \bar{X}) f_{\text{dec}} \text{Br}(Z \rightarrow \ell^+ \ell^-)$$

$$\text{Br}(h \rightarrow aa \rightarrow 4X)|_{\text{eff}} = \text{Br}(h \rightarrow aa) \text{Br}(a \rightarrow X \bar{X})^2 f_{\text{dec}}^2$$

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- Requiring 100 events at  $\sqrt{s} = 13 \text{ TeV}$  with  $300 \text{ fb}^{-1}$  in

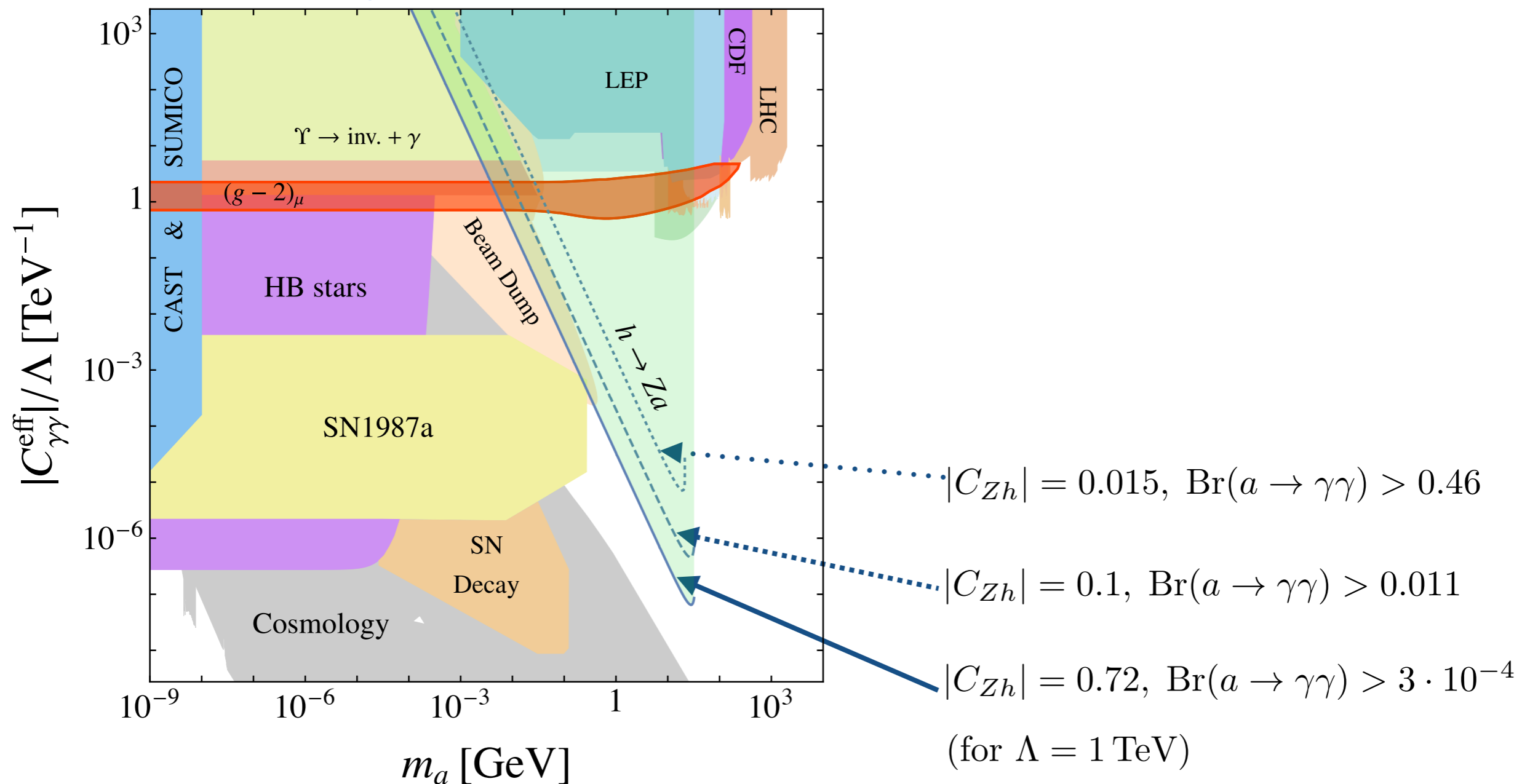
$$h \rightarrow Za \rightarrow \ell^+ \ell^- \gamma \gamma$$

$$h \rightarrow aa \rightarrow 4\gamma$$

# Probing the parameter space

- Constraints on ALP mass and coupling to photons

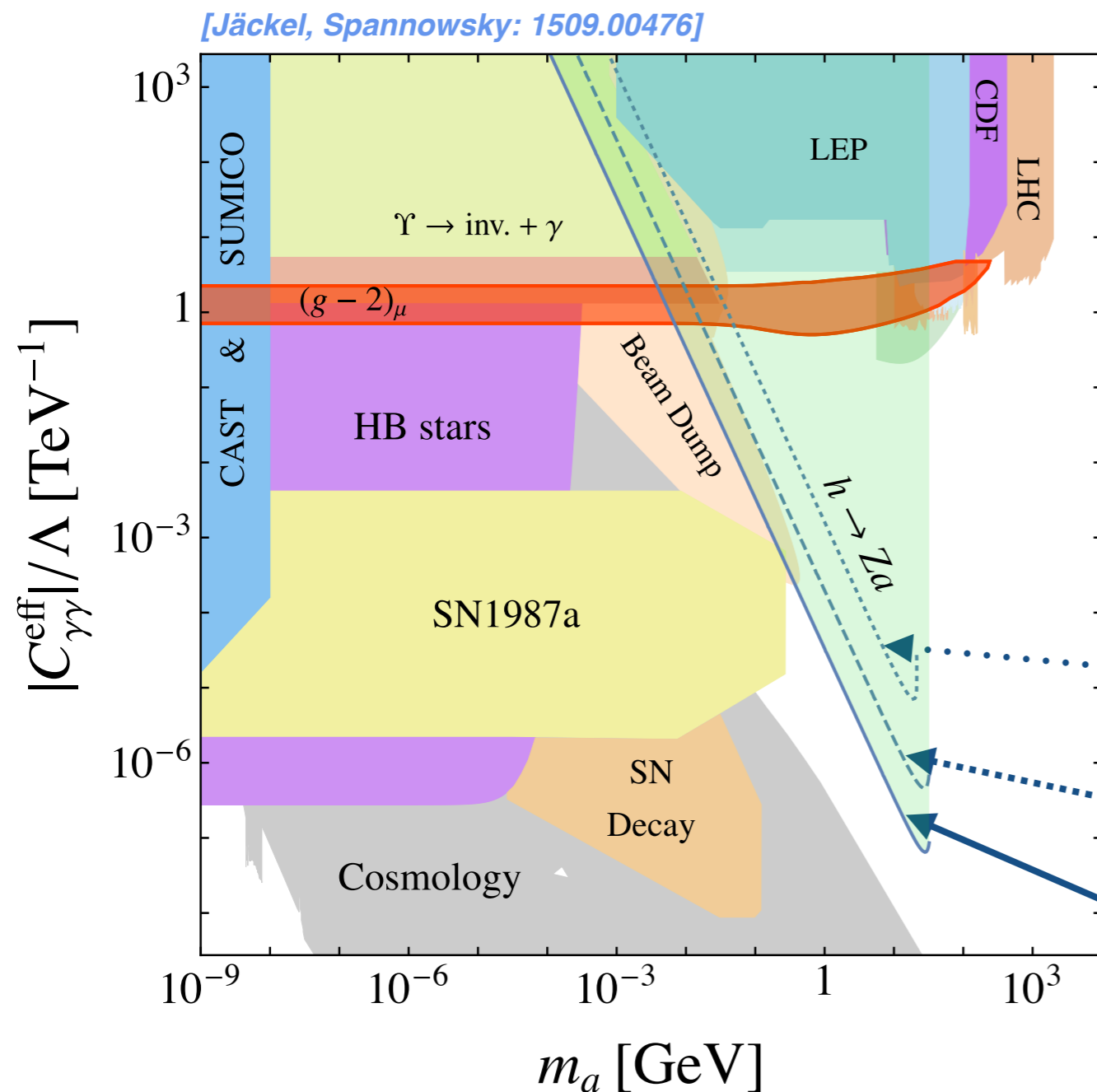
[Jäckel, Spannowsky: 1509.00476]





# Probing the parameter space

- Constraints on ALP mass and coupling to photons



- ALP-photon coupling can be probed if ALP decays predominantly into other particles
- Region preferred by  $(g-2)_\mu$  almost completely covered

$|C_{Zh}| = 0.015, \text{Br}(a \rightarrow \gamma\gamma) > 0.46$

$|C_{Zh}| = 0.1, \text{Br}(a \rightarrow \gamma\gamma) > 0.011$

$|C_{Zh}| = 0.72, \text{Br}(a \rightarrow \gamma\gamma) > 3 \cdot 10^{-4}$

(for  $\Lambda = 1 \text{ TeV}$ )

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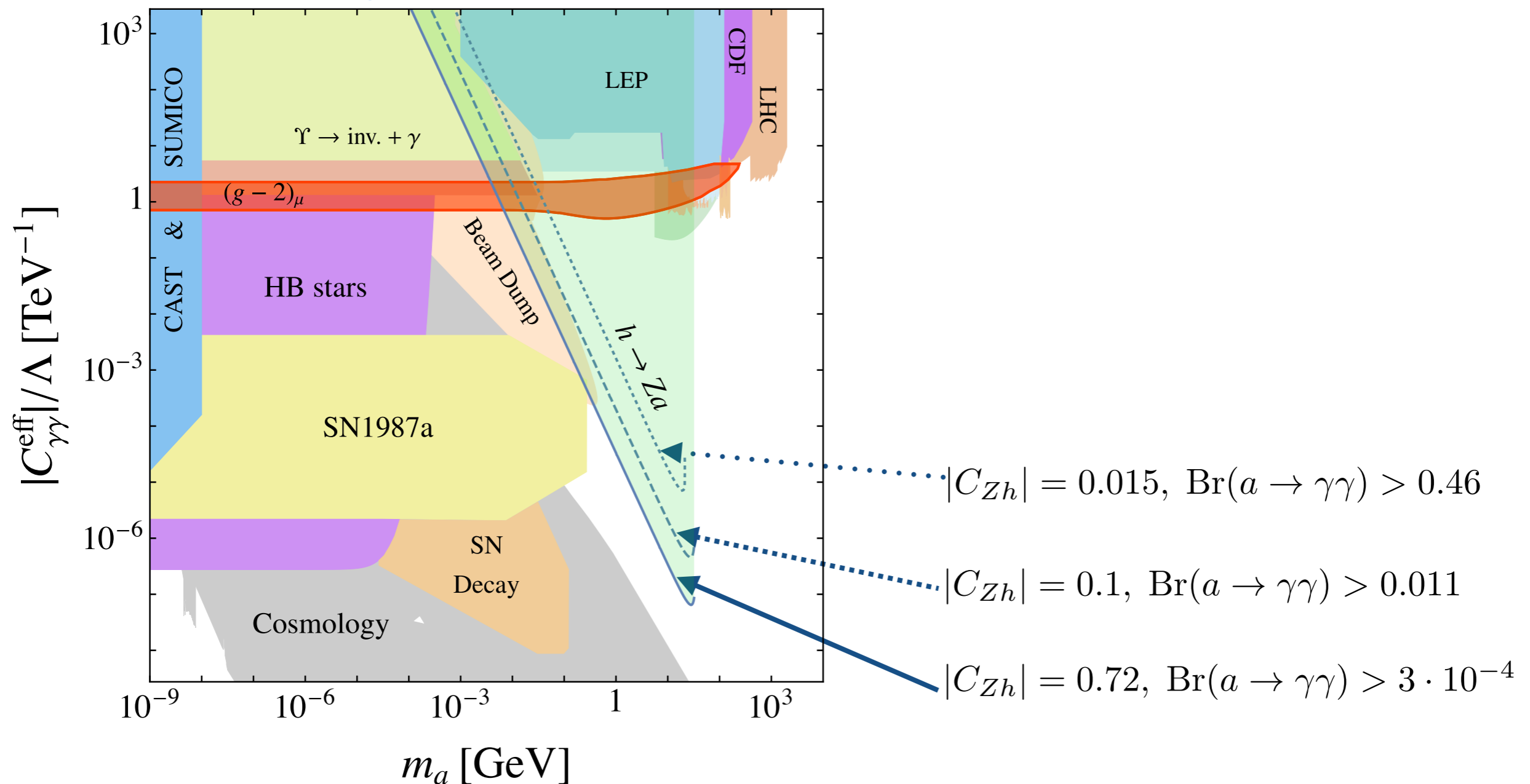
- For  $L_a \gg L_{\text{det}}$ , effective BR independent of  $\text{Br}(a \rightarrow X \bar{X})$

$$f_{\text{dec}} \approx (\pi/2) \frac{L_{\text{det}}}{L_a} \propto \frac{\Gamma(a \rightarrow X \bar{X})}{\text{Br}(a \rightarrow X \bar{X})}$$

# Probing the parameter space

- Large hierarchy in couplings can be plausible

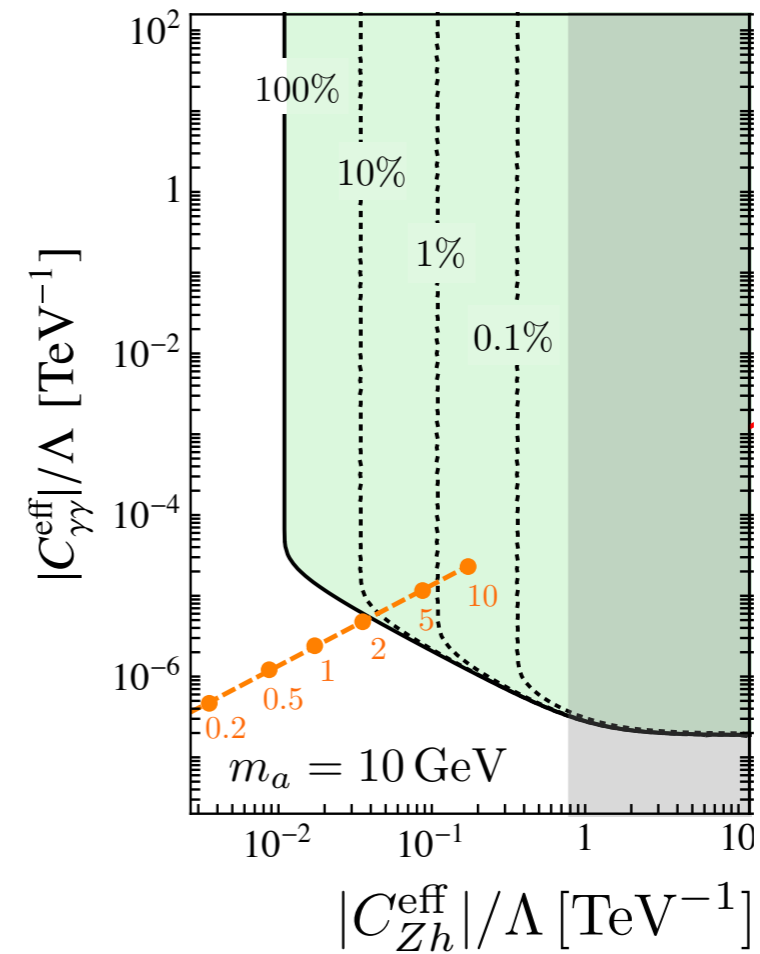
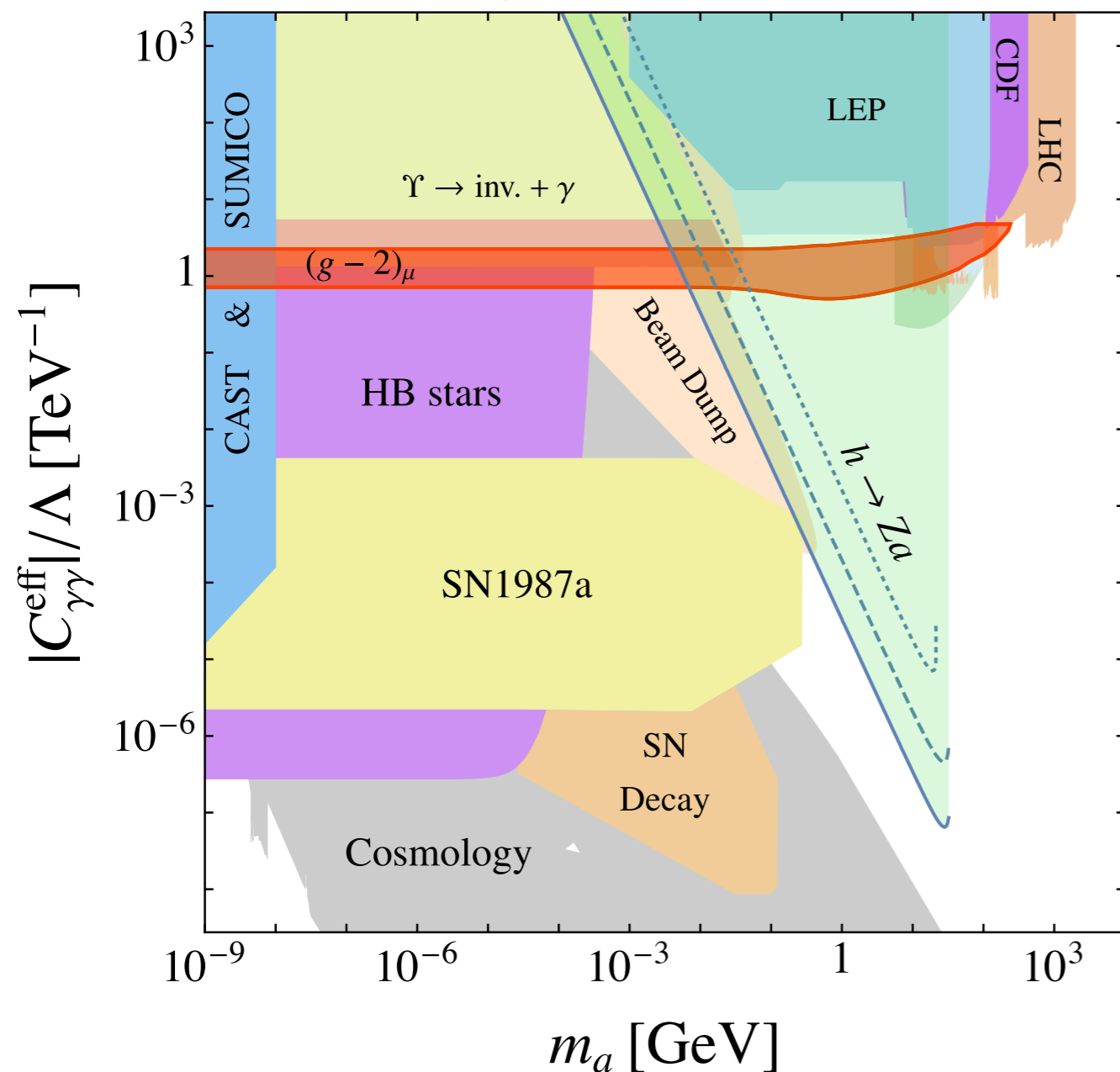
[Jäckel, Spannowsky: 1509.00476]



# Probing the parameter space

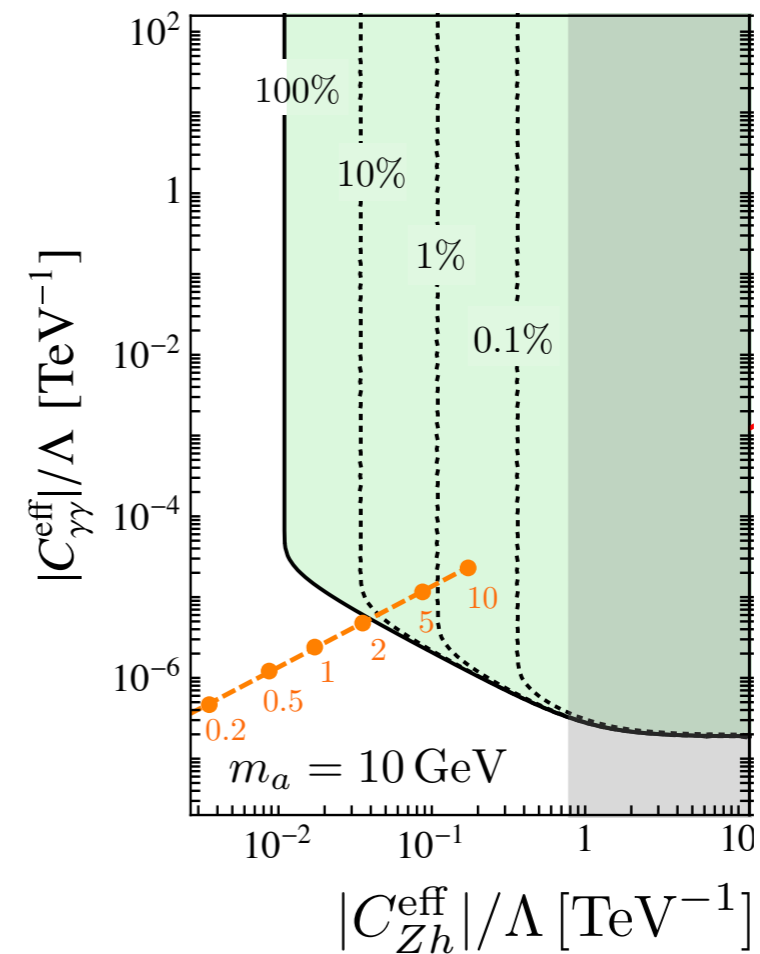
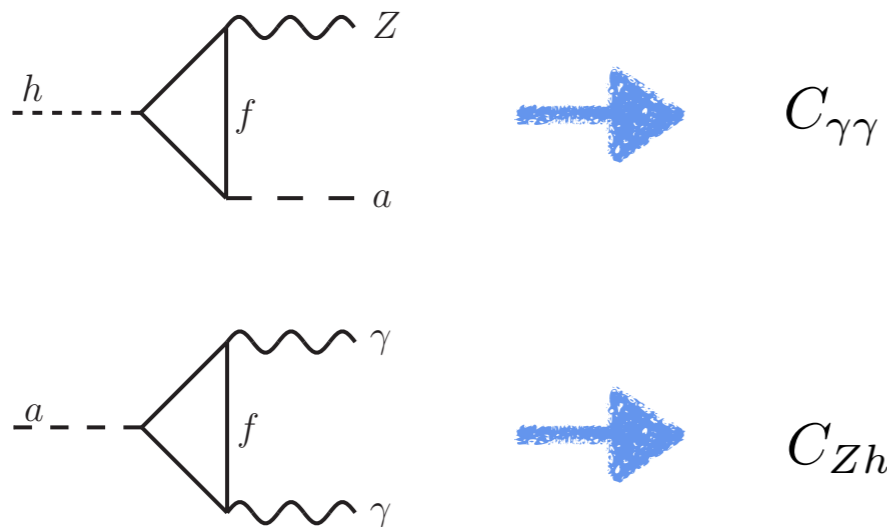
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[Jäckel, Spannowsky: 1509.00476]



# Probing the parameter space

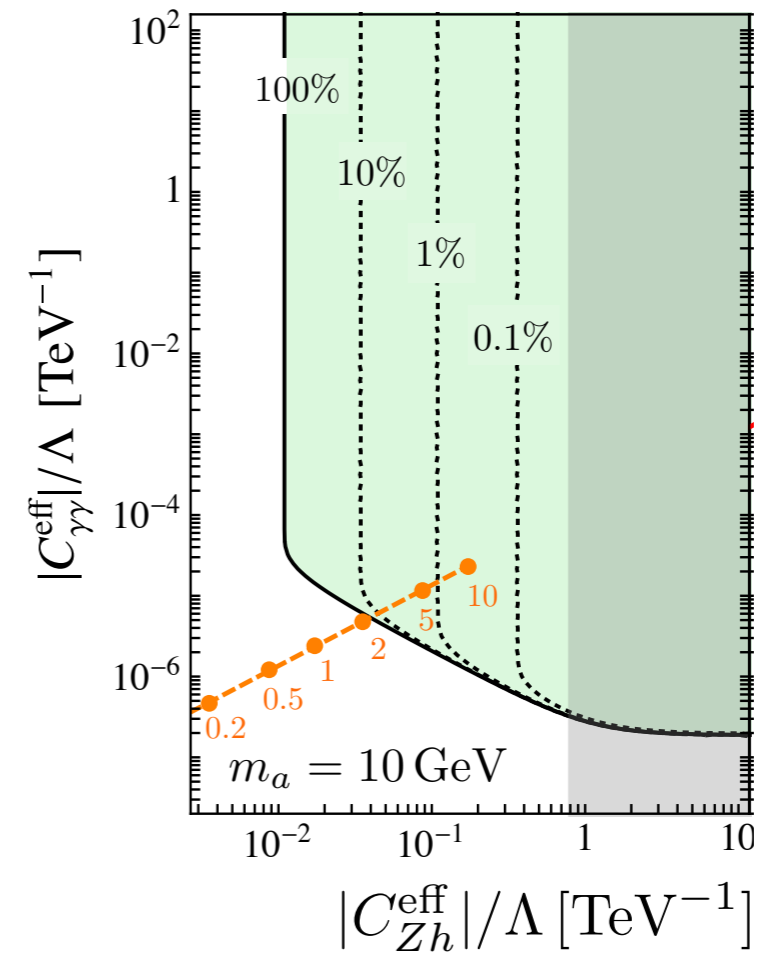
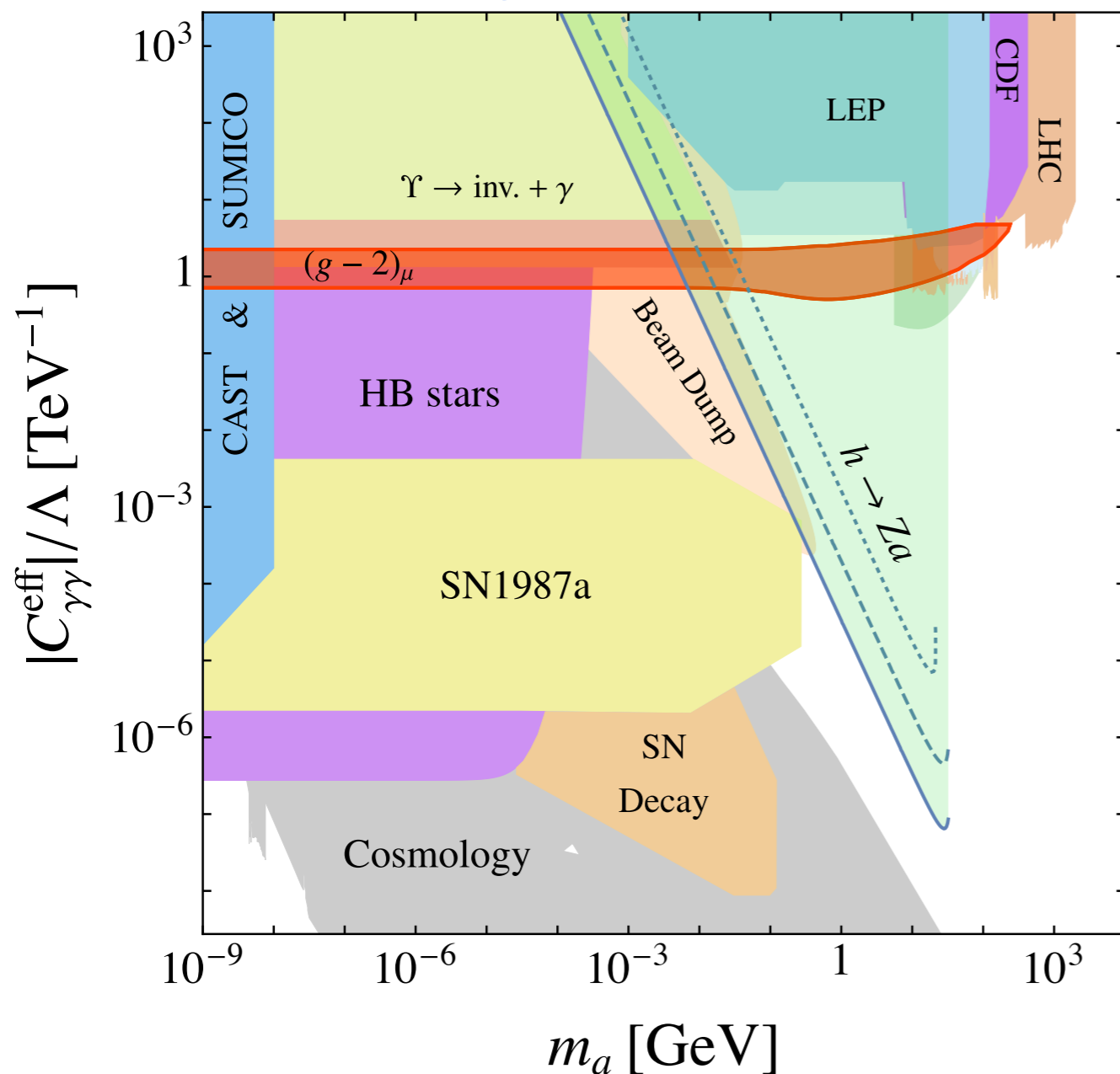
- Large hierarchy in couplings can be plausible
- Integrating out the top



# Probing the parameter space

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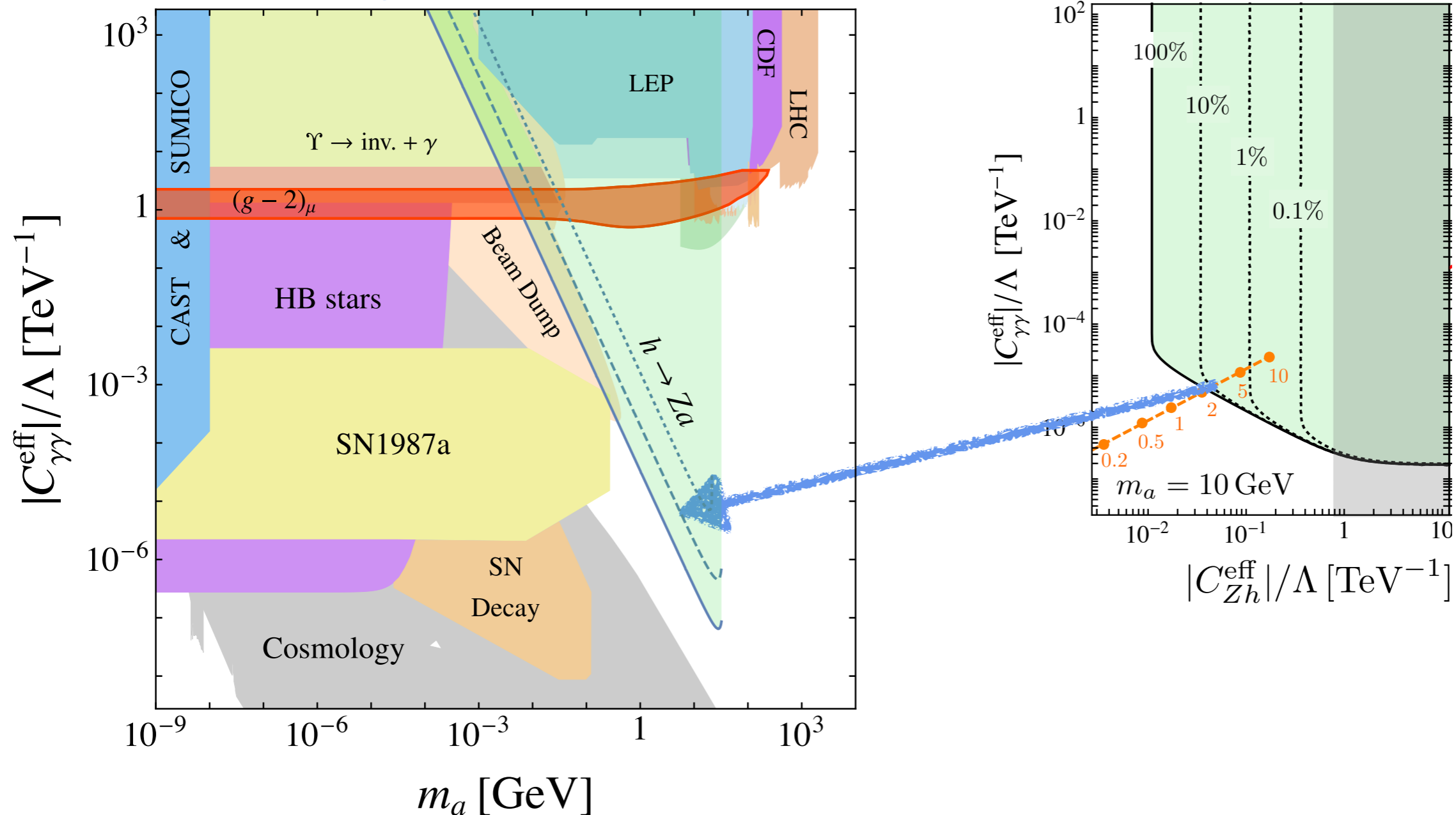
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[Jäckel, Spannowsky: 1509.00476]

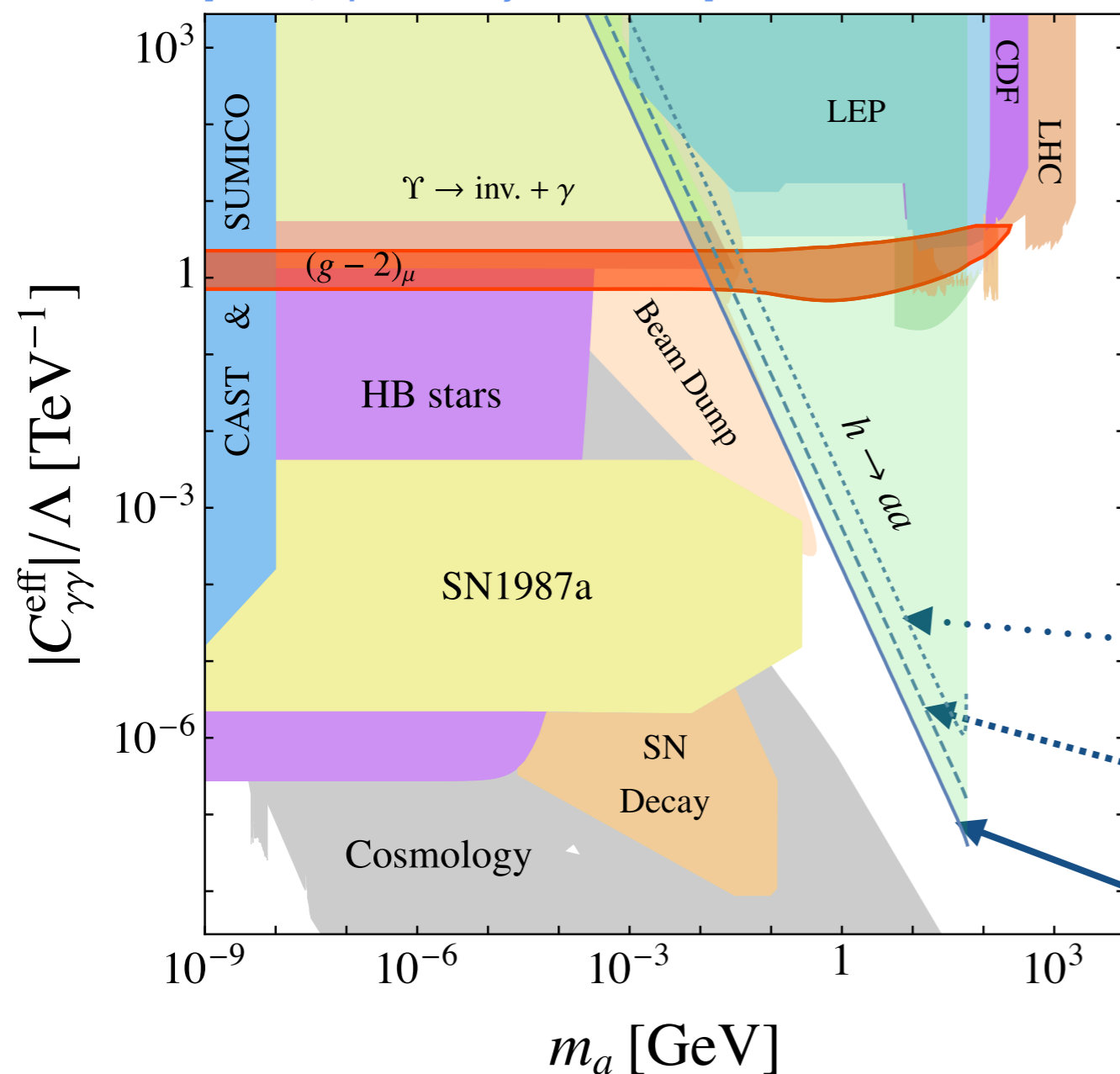




# Probing the parameter space

- Constraints on ALP mass and coupling to photons

[Jäckel, Spannowsky: 1509.00476]



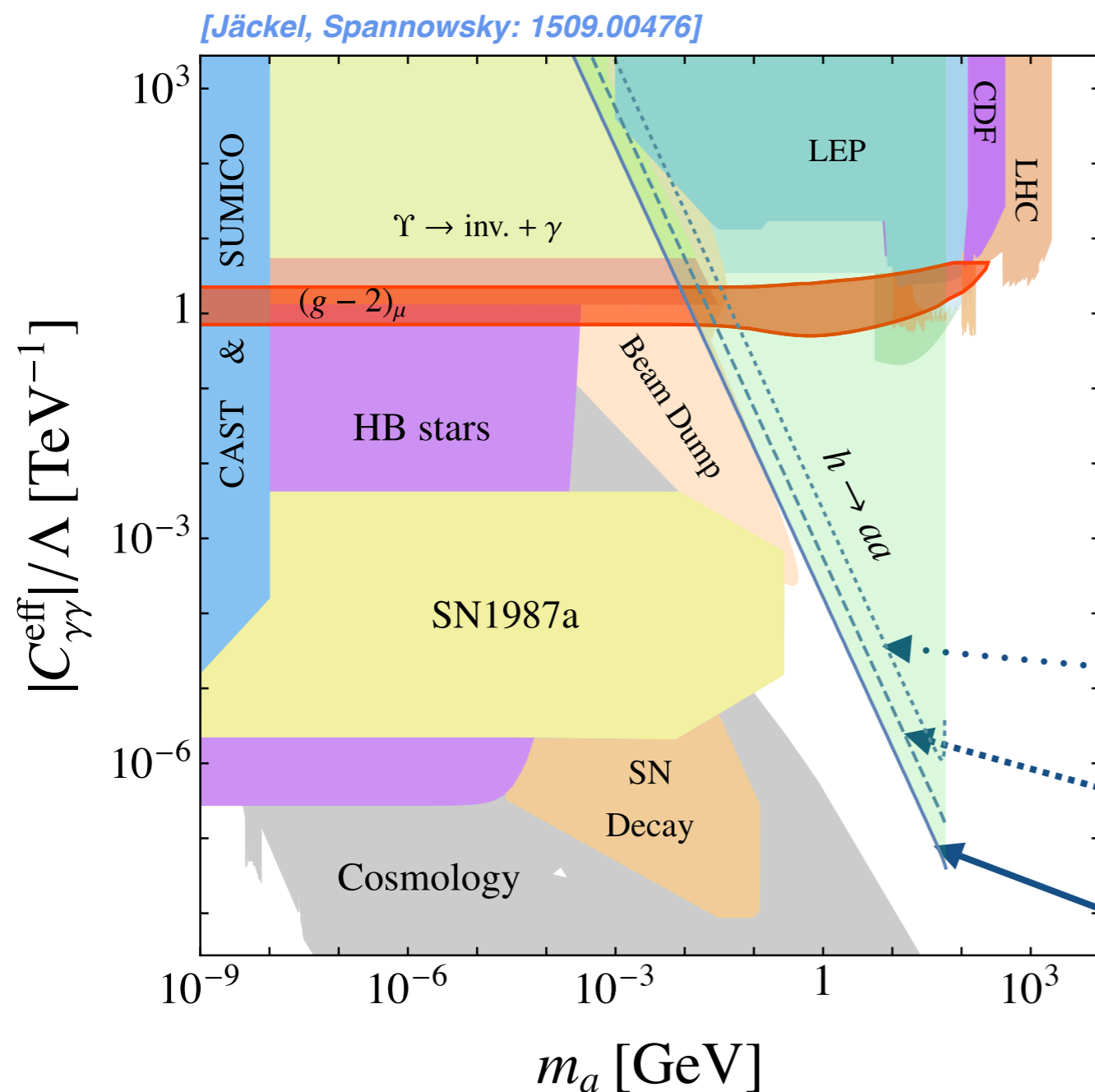
$|C_{ah}| = 0.01, \text{Br}(a \rightarrow \gamma\gamma) > 0.49$

$|C_{ah}| = 0.1, \text{Br}(a \rightarrow \gamma\gamma) > 0.049$

$|C_{ah}| = 1, \text{Br}(a \rightarrow \gamma\gamma) > 0.006$   
(for  $\Lambda = 1 \text{ TeV}$ )

# Probing the parameter space

- Constraints on ALP mass and coupling to photons



- ALP-photon coupling can be probed if ALP decays predominantly into other particles
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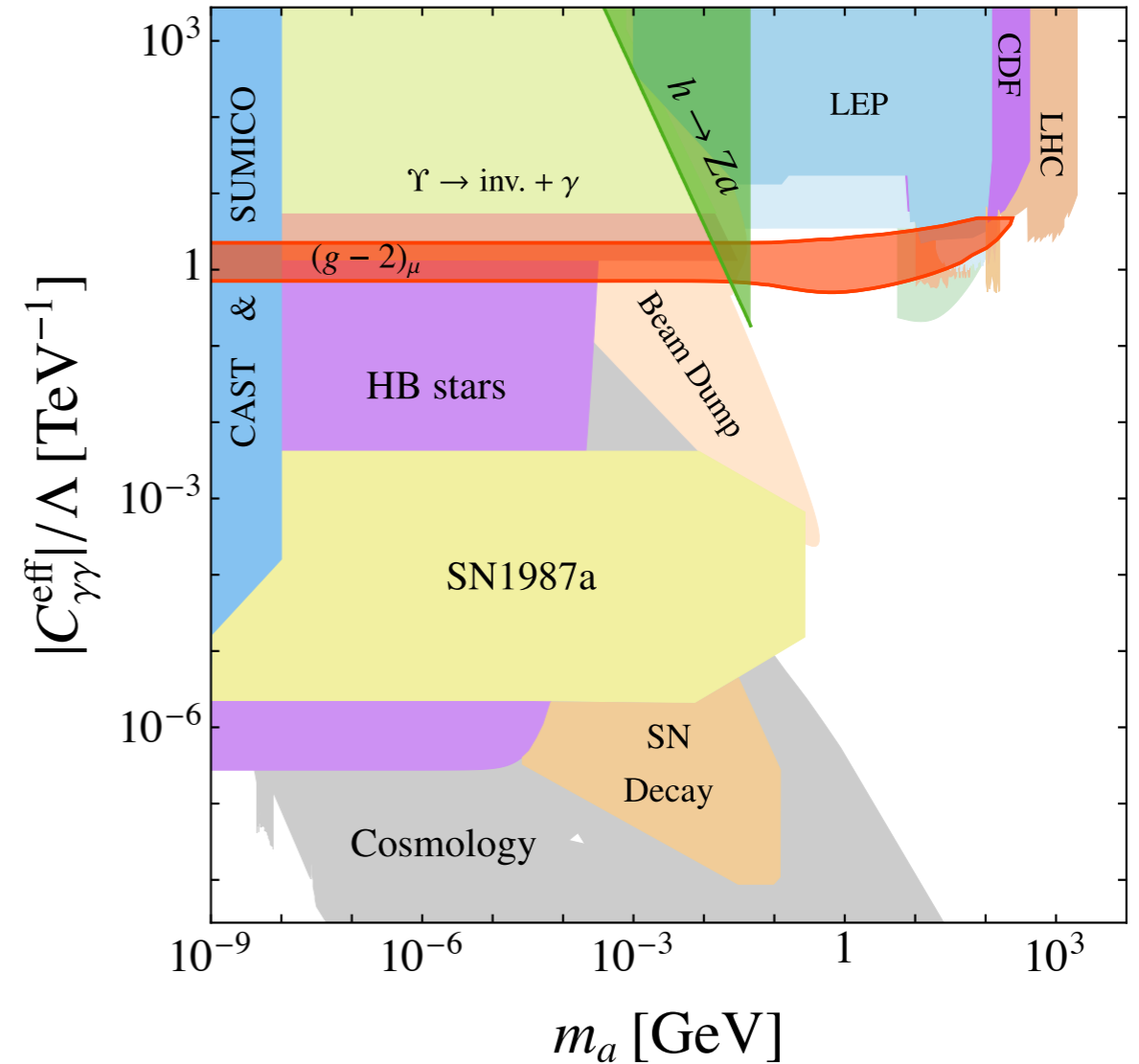
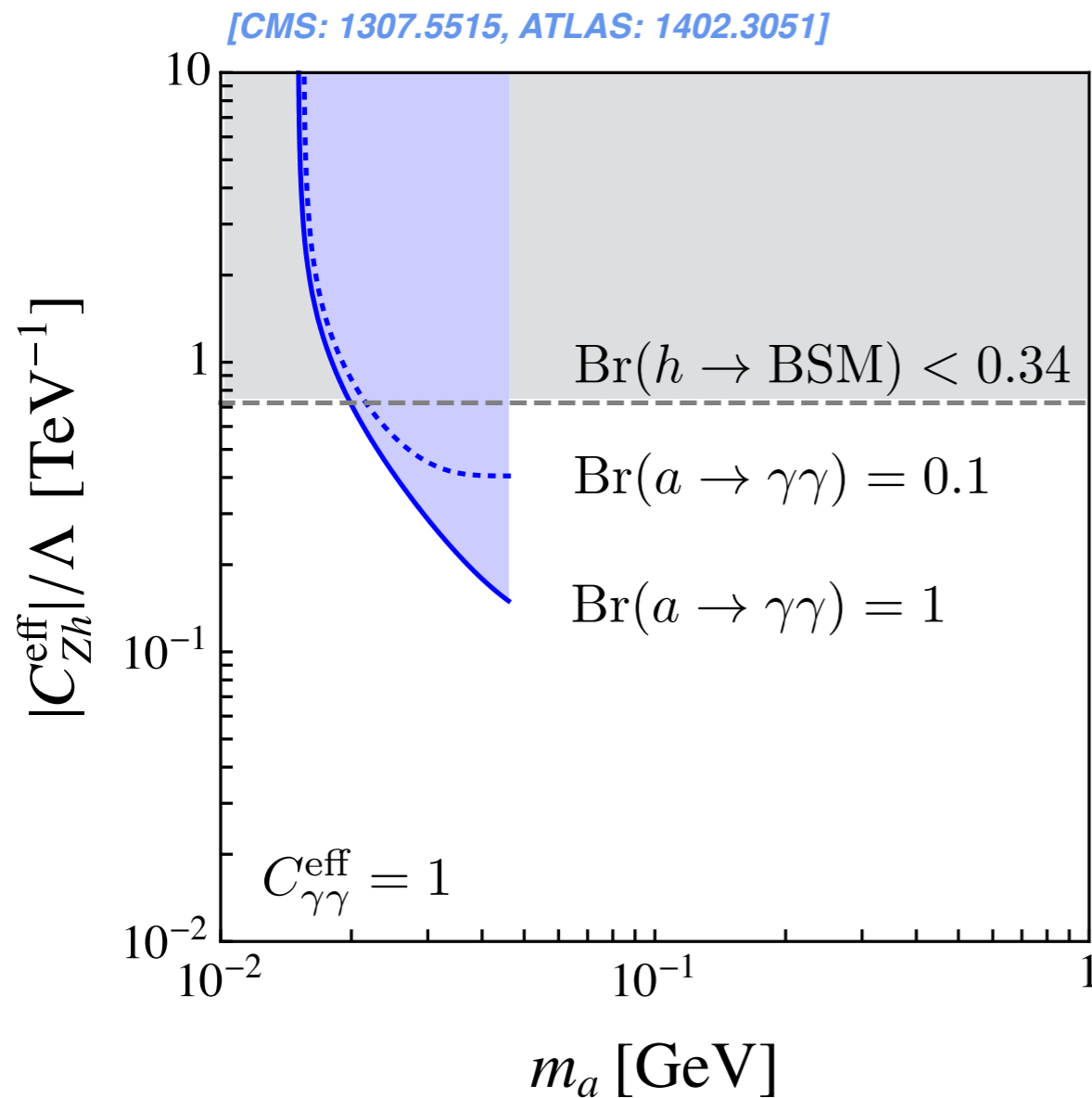
$|C_{ah}| = 0.01, \text{Br}(a \rightarrow \gamma\gamma) > 0.49$

$|C_{ah}| = 0.1, \text{Br}(a \rightarrow \gamma\gamma) > 0.049$

$|C_{ah}| = 1, \text{Br}(a \rightarrow \gamma\gamma) > 0.006$   
(for  $\Lambda = 1 \text{ TeV}$ )

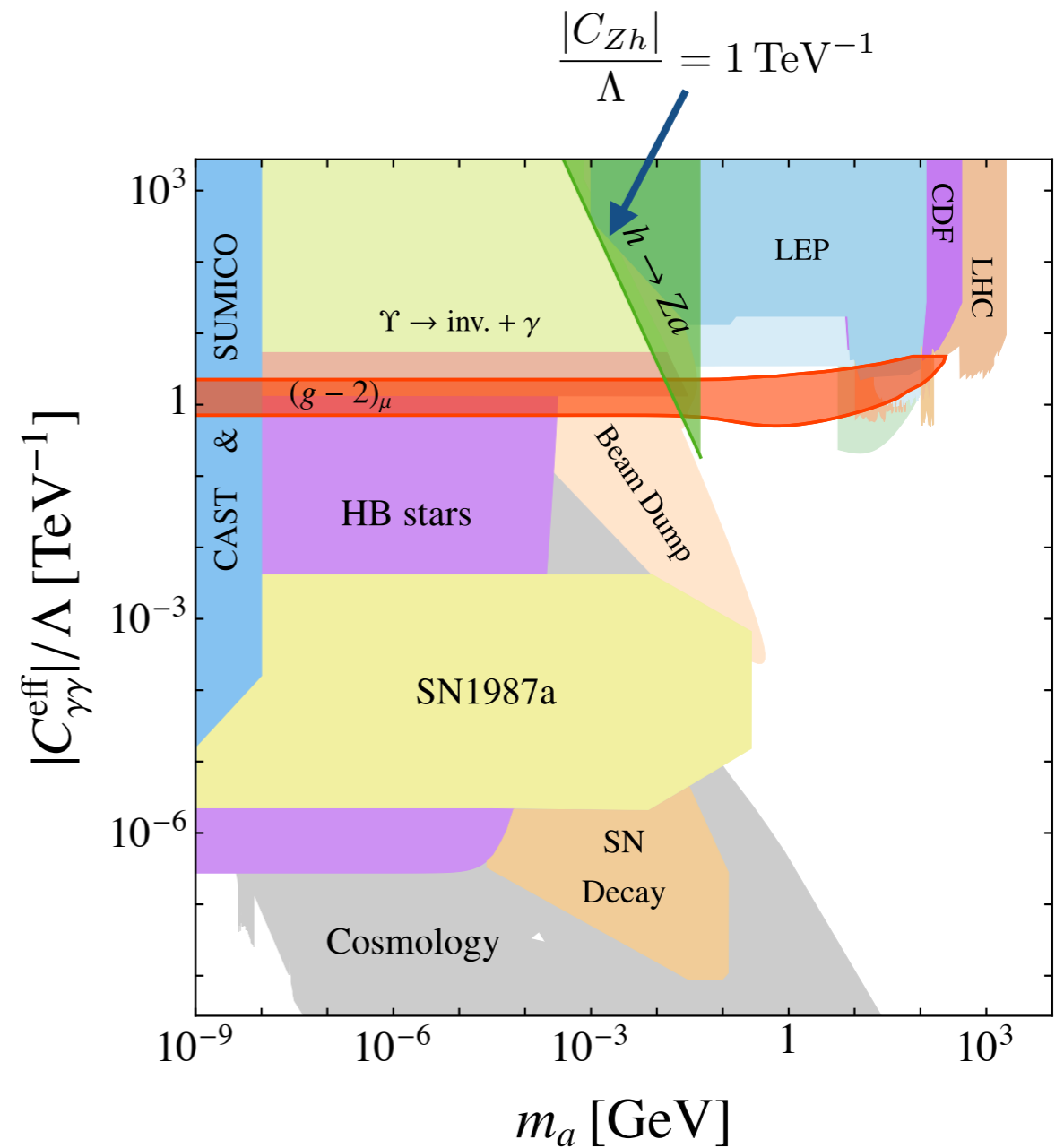
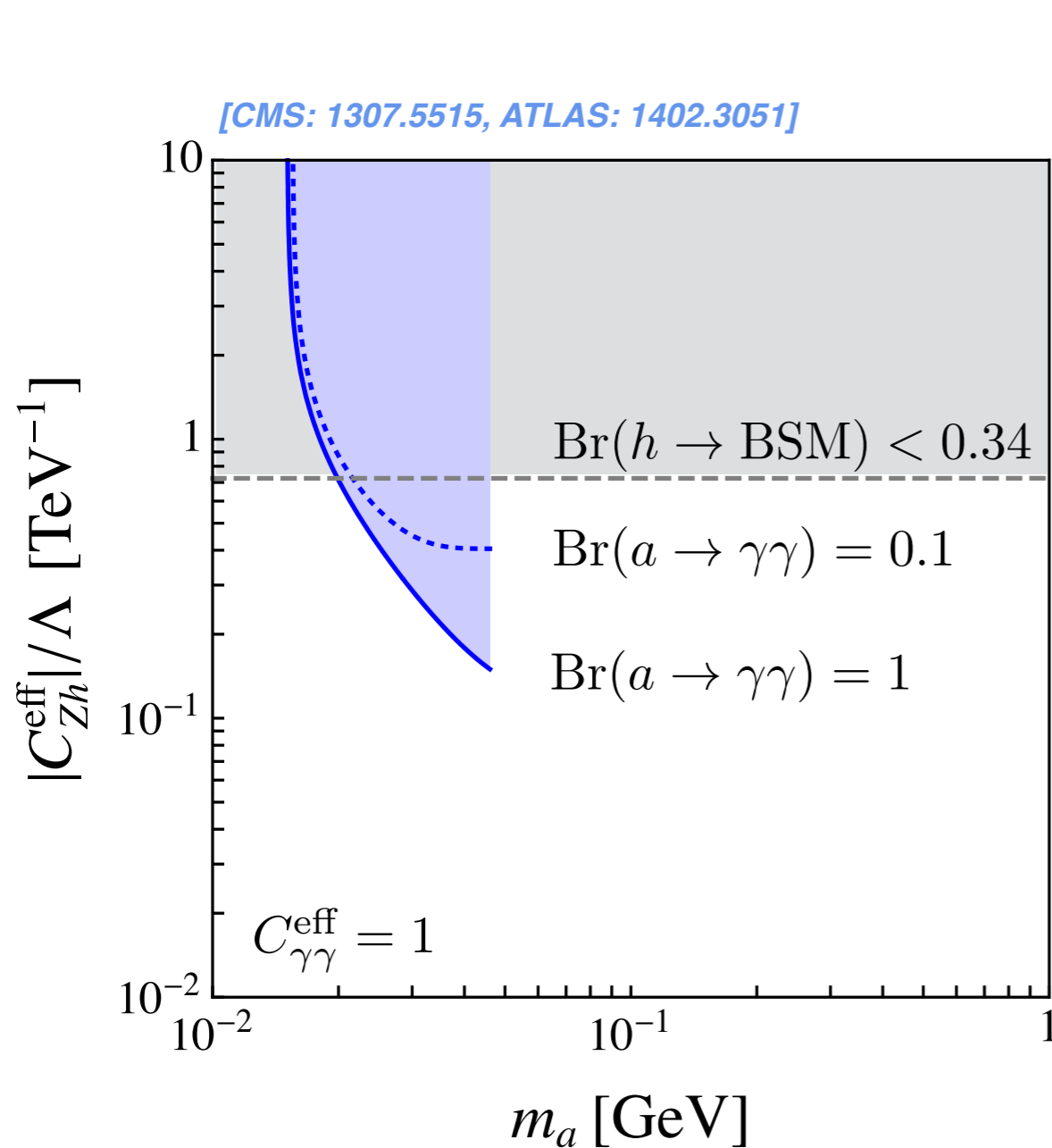
# Current exclusion bounds

- Current bounds on  $h \rightarrow Za$



# Current exclusion bounds

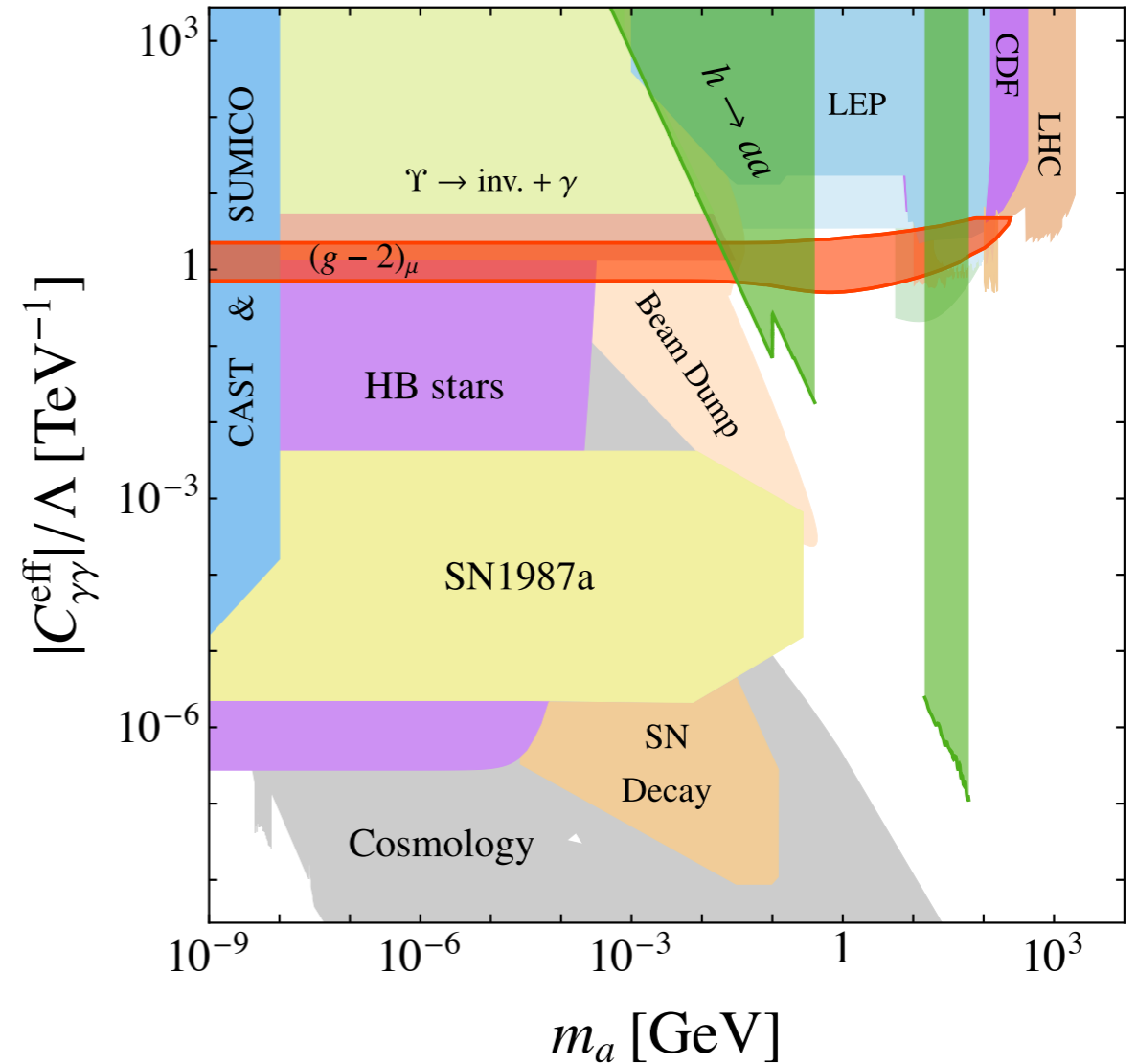
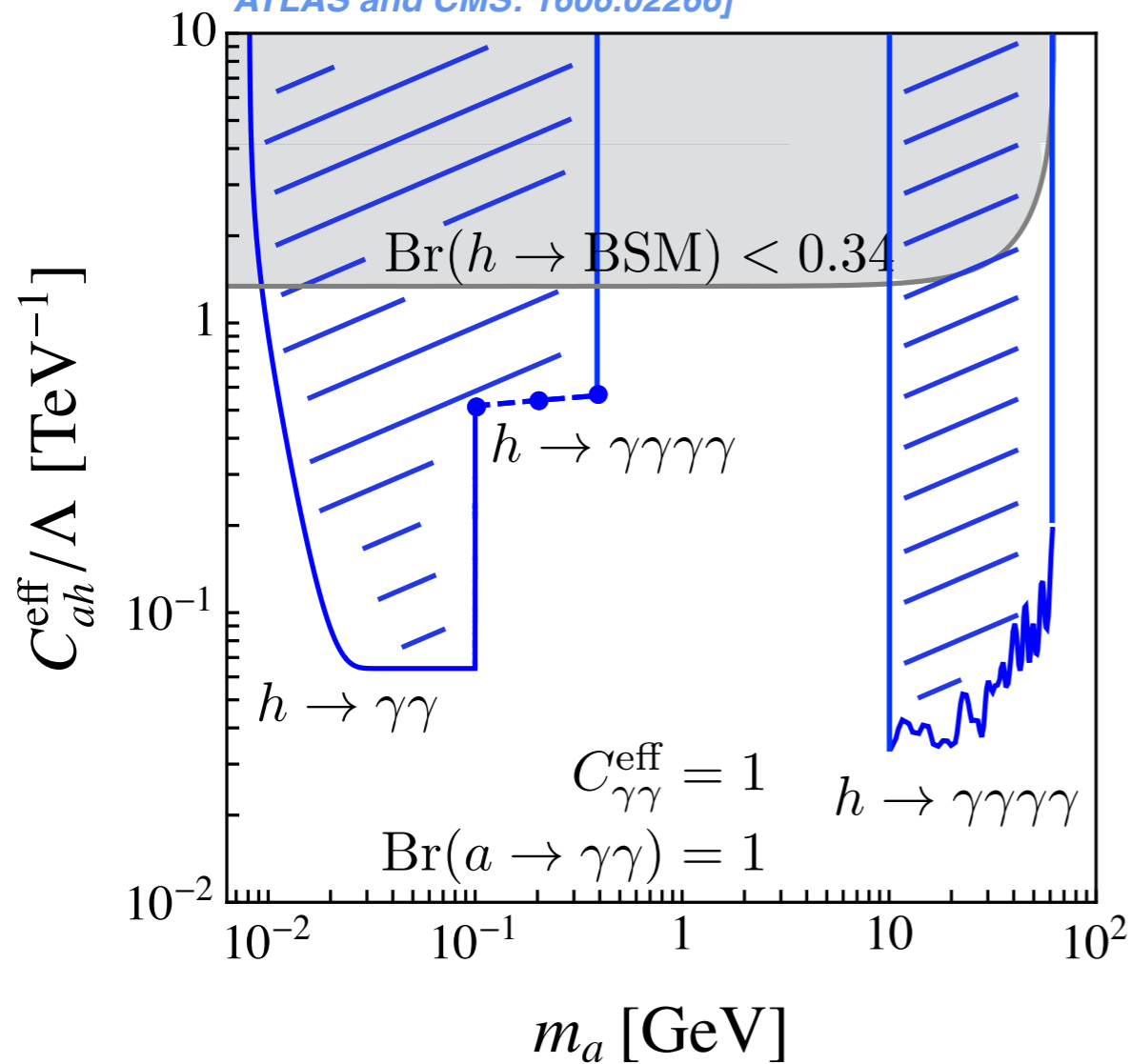
- Current bounds on  $h \rightarrow Za$



# Current exclusion bounds

- Current bounds on  $h \rightarrow aa$

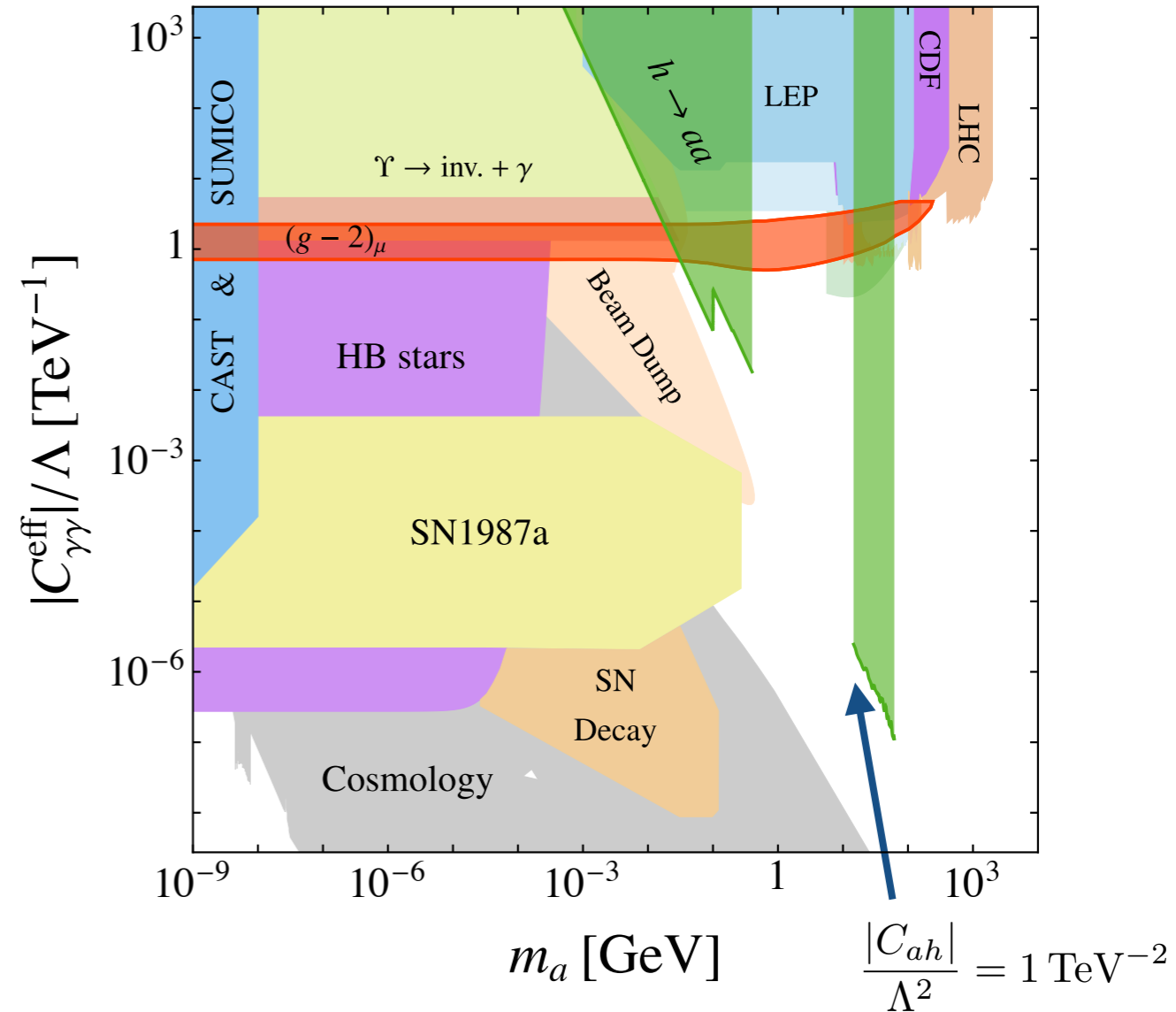
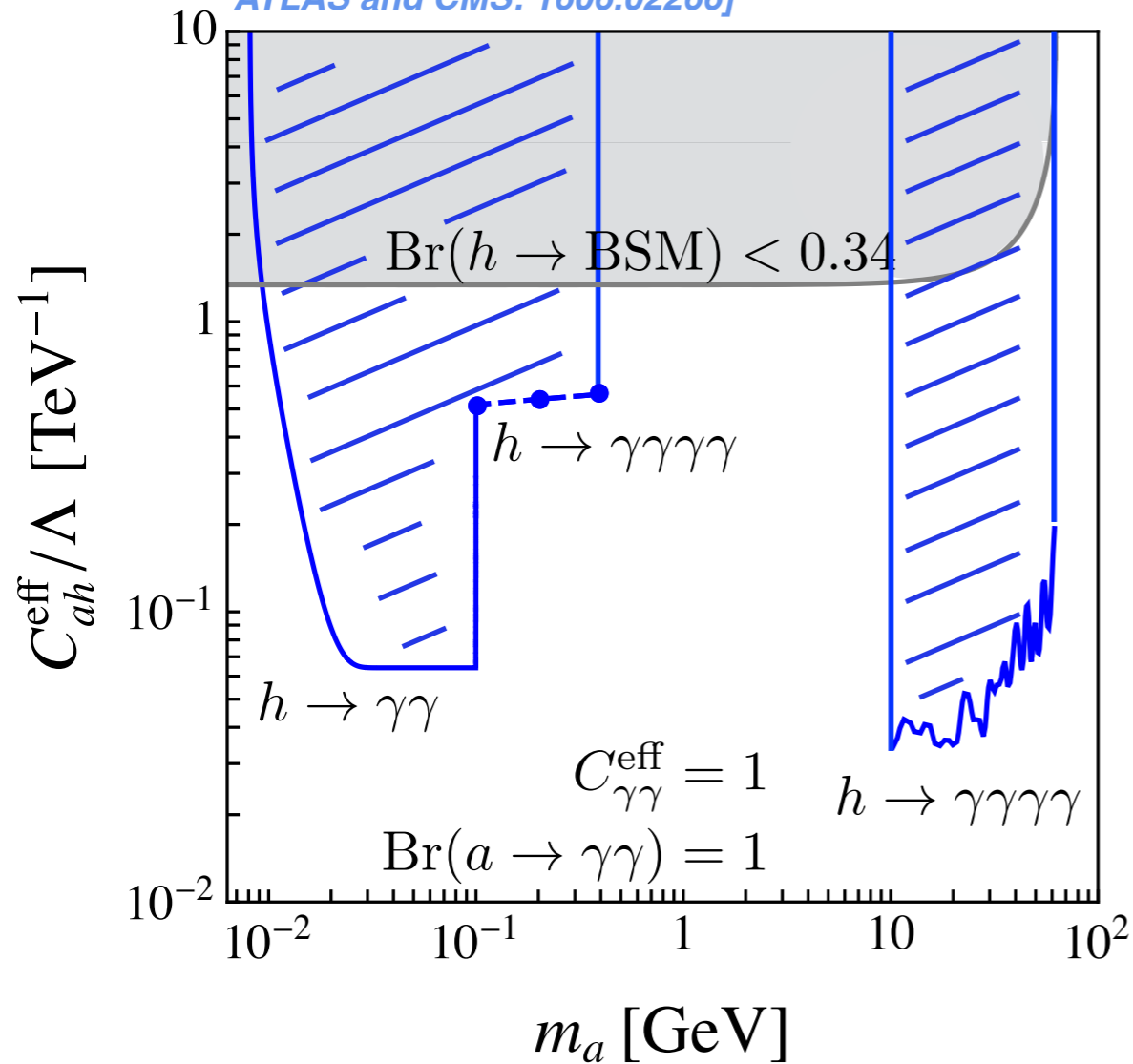
[ATLAS-CONF-2012-079, ATLAS: 1509.05051,  
ATLAS and CMS: 1606.02266]



# Current exclusion bounds

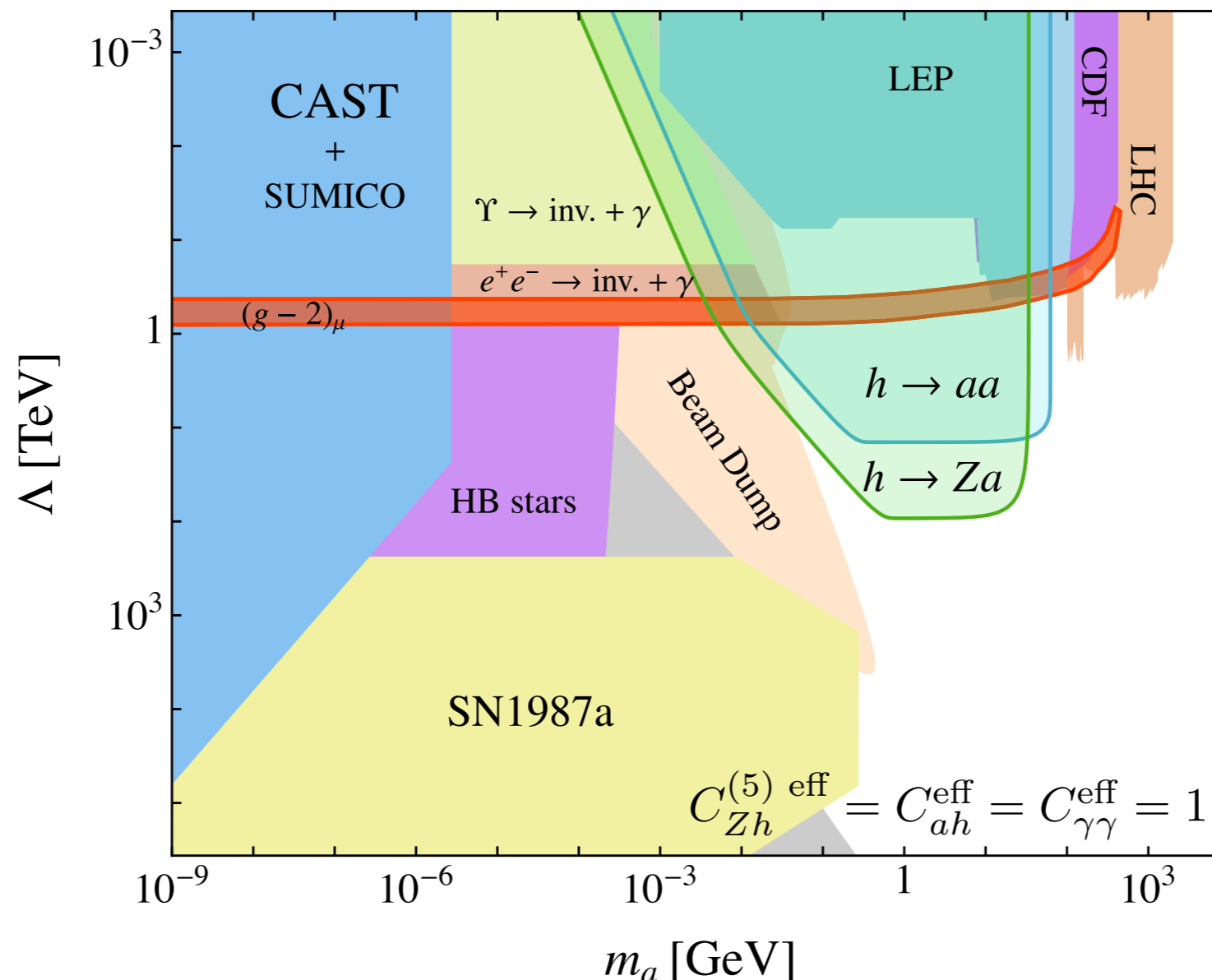
- Current bounds on  $h \rightarrow aa$

[ATLAS-CONF-2012-079, ATLAS: 1509.05051,  
ATLAS and CMS: 1606.02266]



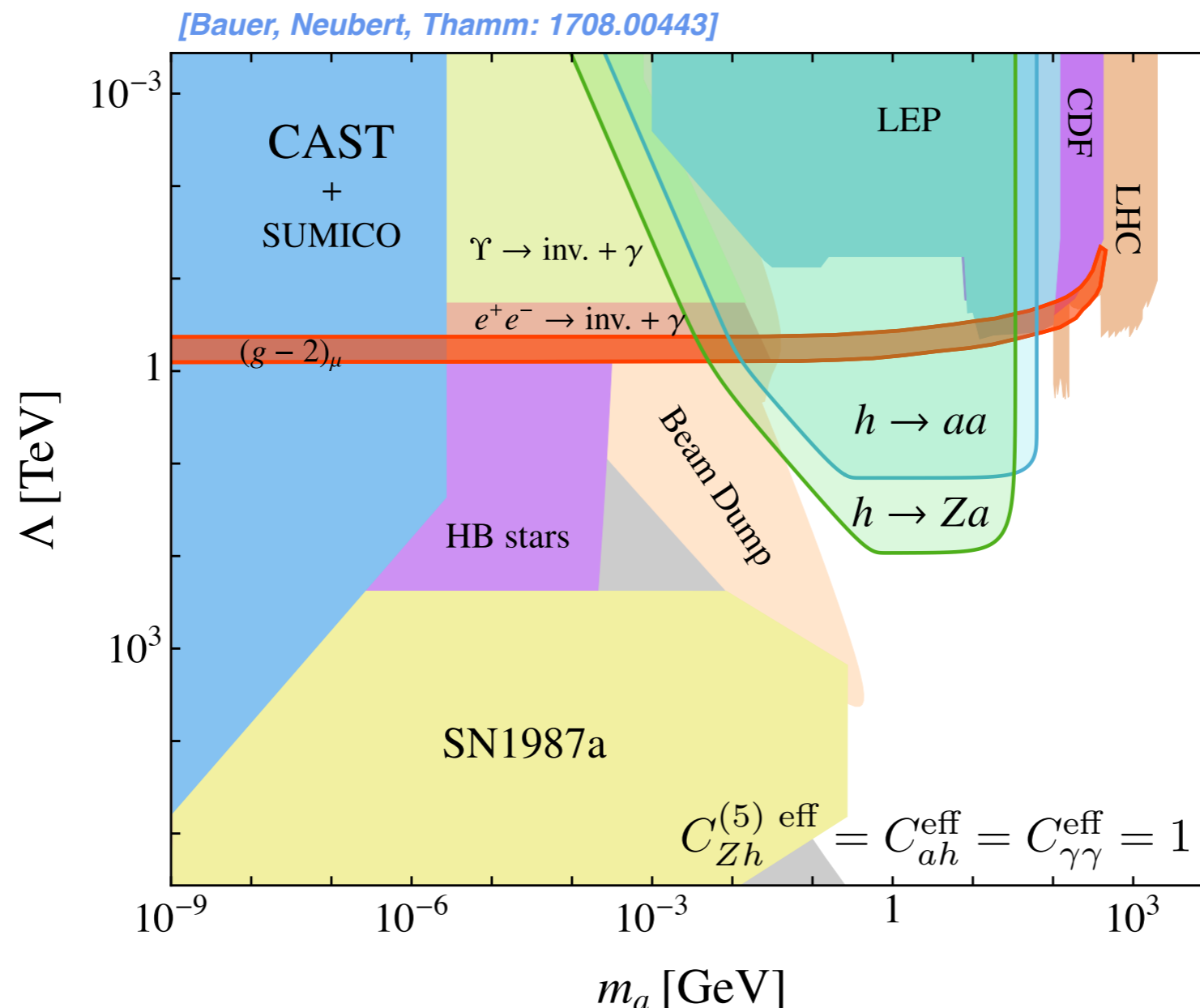
# Probing the parameter space

- Constraints on ALP mass and coupling to photons



# Probing the parameter space

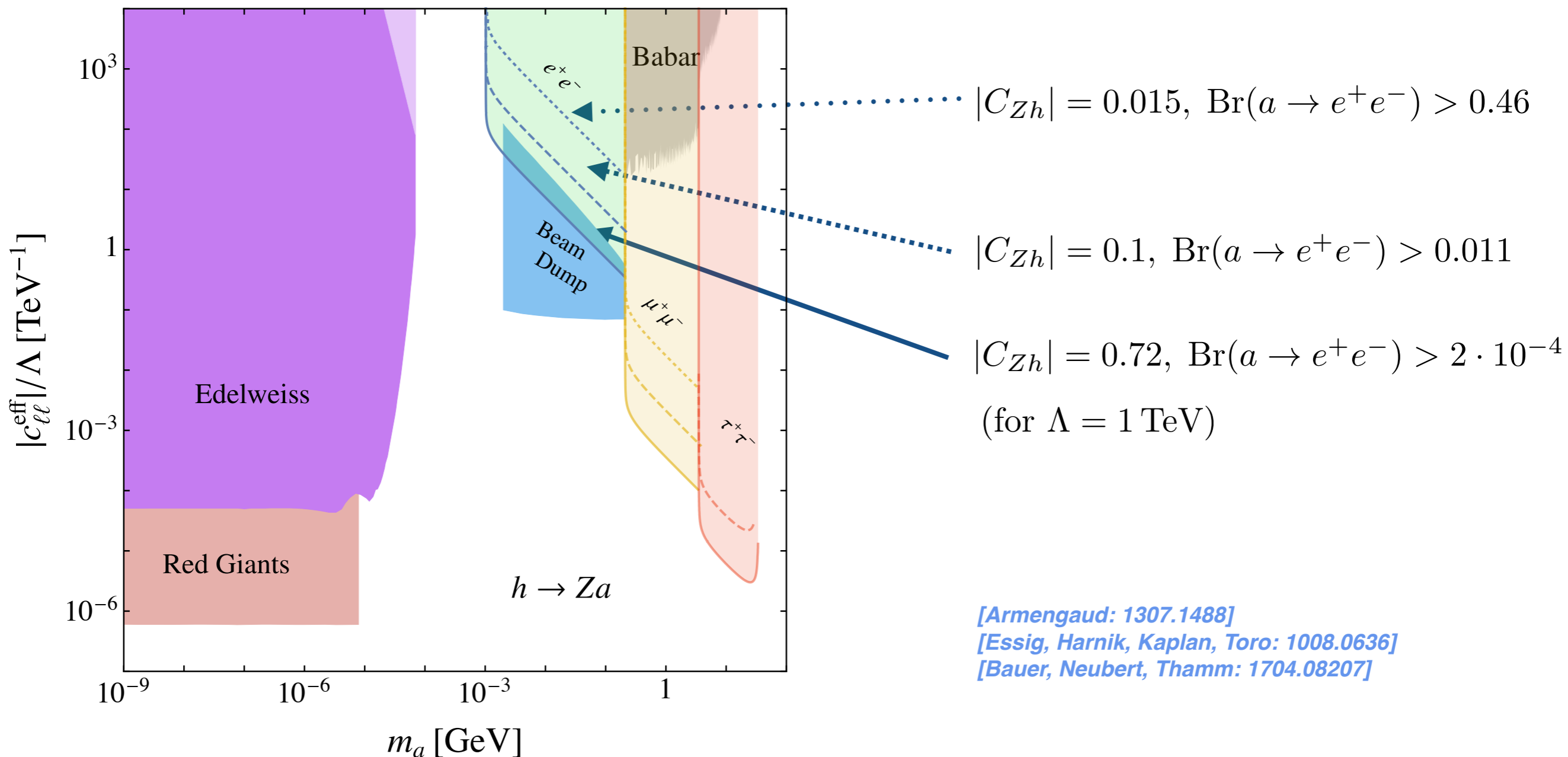
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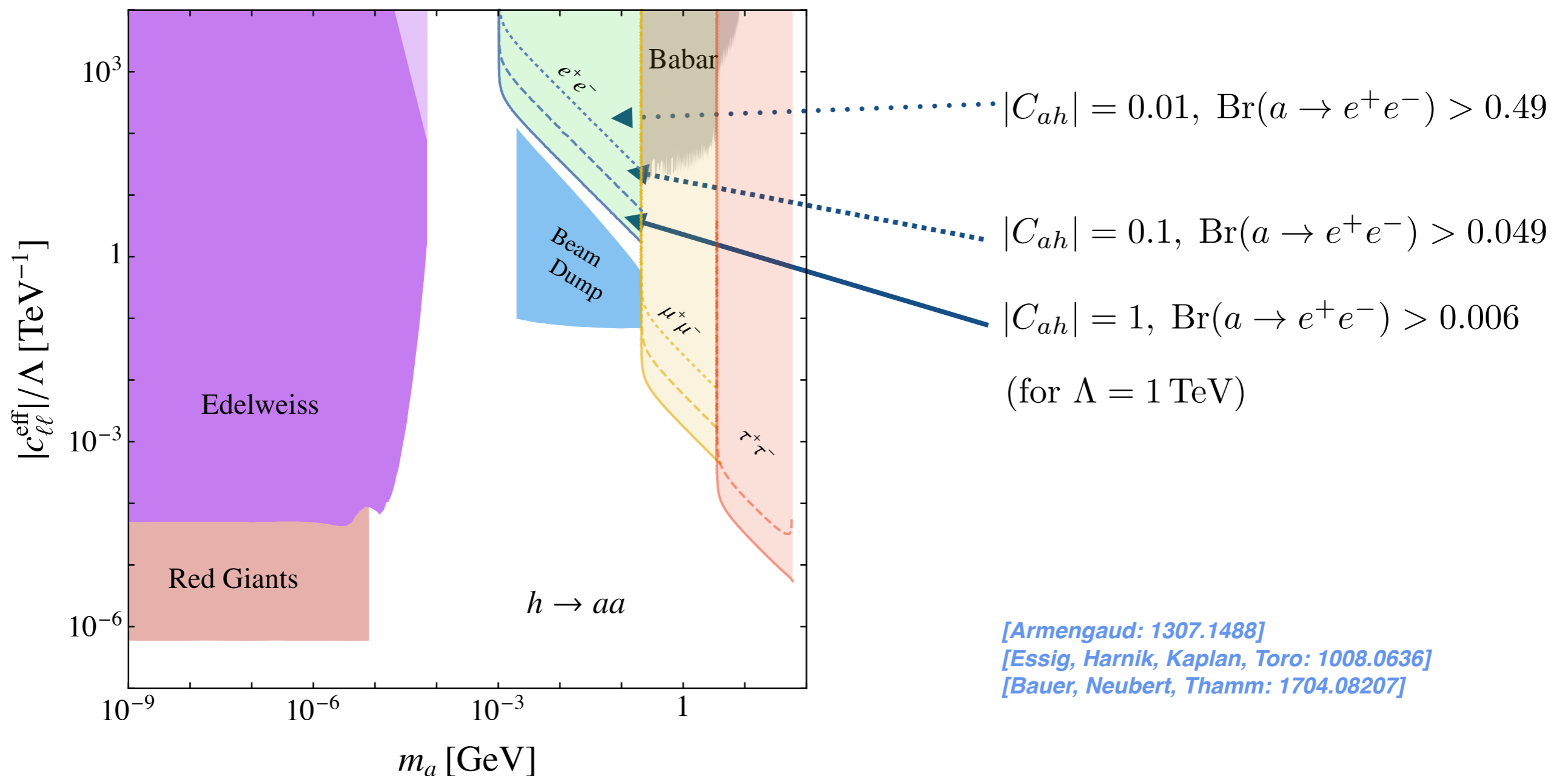
# Probing the parameter space

- Constraints on ALP mass and coupling to leptons



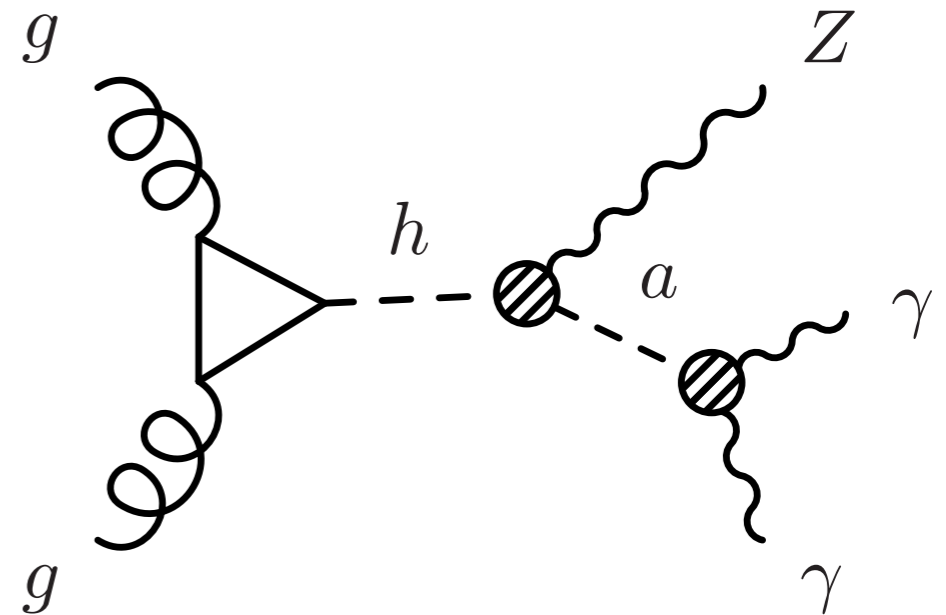
# Probing the parameter space

- Constraints on ALP mass and coupling to leptons



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- Motivation
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# Muon $(g - 2)_\mu$

- Persistent deviation

*[Particle Data Group 2016]*

$$a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = (288 \pm 63 \pm 49) \cdot 10^{-11}$$

# Muon $(g - 2)_\mu$

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- Differs from zero by more than  $3\sigma$

# Muon $(g - 2)_\mu$

consistent dev

[Particle Data Group 2016]

$$a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = (288 \pm 63 \pm 49) \cdot 10^{-11}$$

- Differs from zero by more than 5

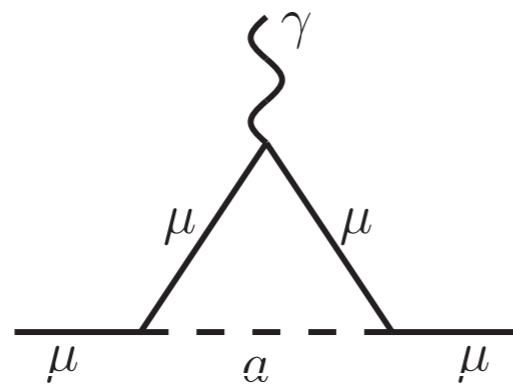
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- Persistent deviation

[Particle Data Group 2016]

$$a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = (288 \pm 63 \pm 49) \cdot 10^{-11}$$

- Differs from zero by more than  $3\sigma$
- ALP can account for discrepancy



[Haber, Kane, Sterling: Nucl. Phys. B 161 (1979)]

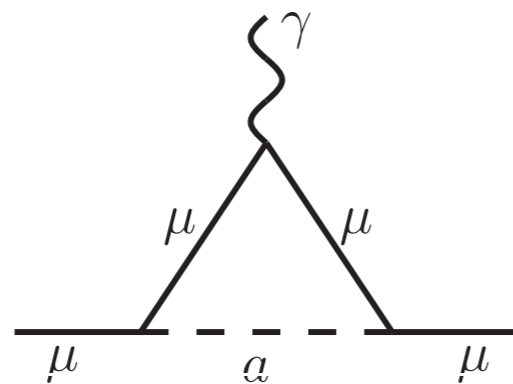
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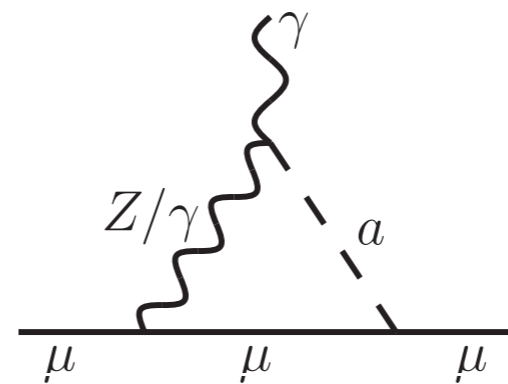
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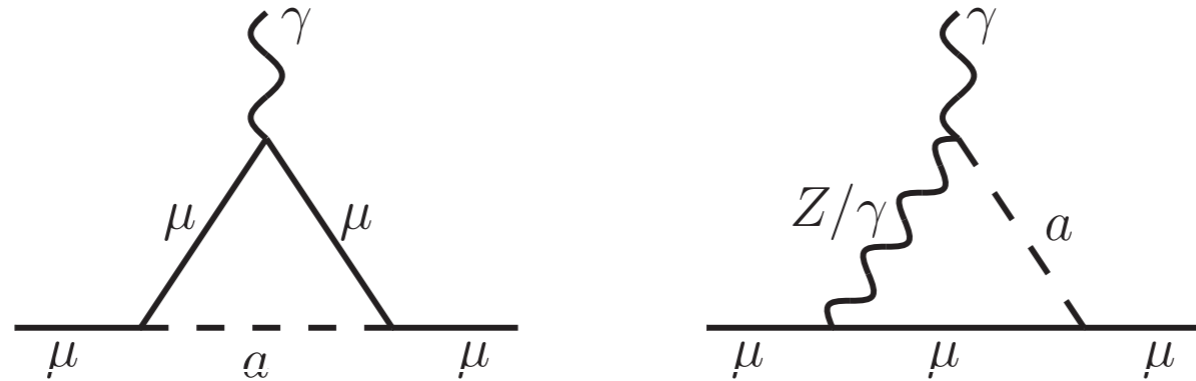
[Haber, Kane, Sterling: Nucl. Phys. B 161 (1979)]



[Chang, Chang, Chou, Keung: 0009292]  
[Marciano, Masiero, Paradisi, Passera: 1607.010122]

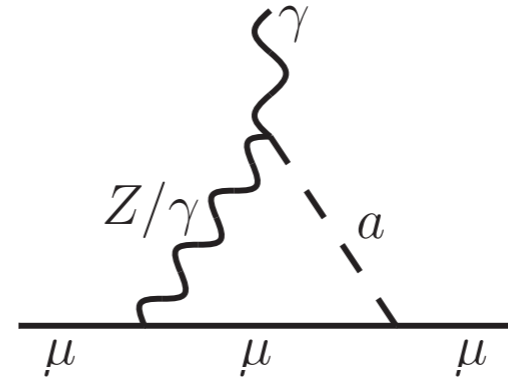
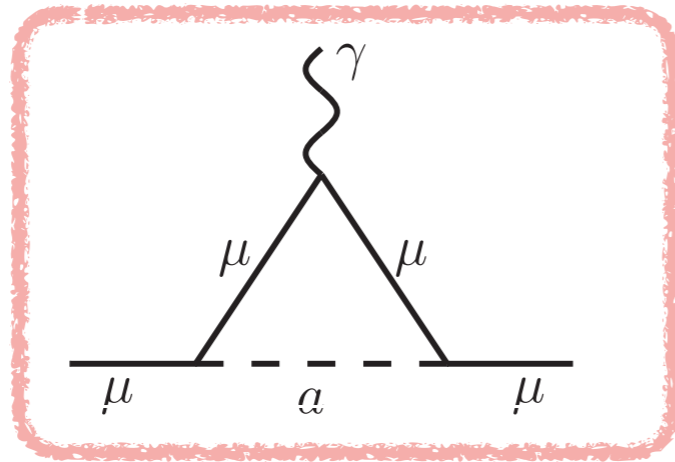


# Muon $(g - 2)_\mu$



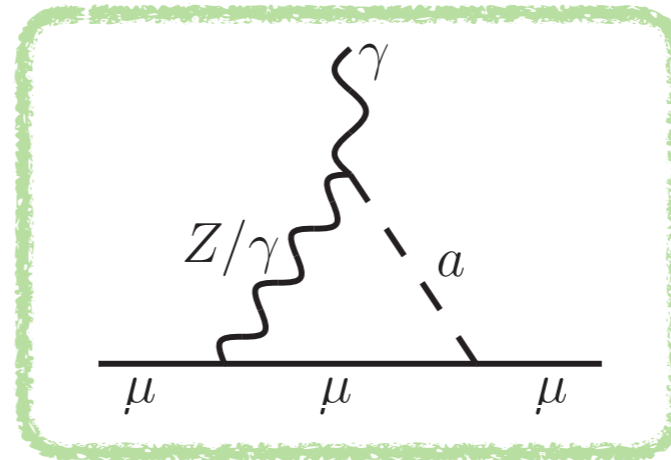
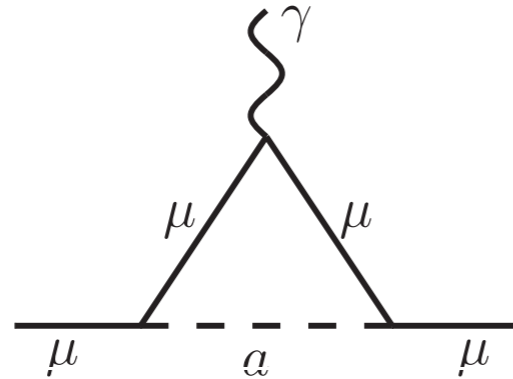
$$\delta a_\mu = \frac{m_\mu^2}{\Lambda^2} \left\{ K_{a\mu}(\mu) - \frac{(c_{\mu\mu})^2}{16\pi^2} h_1\left(\frac{m_a^2}{m_\mu^2}\right) - \frac{2\alpha}{\pi} c_{\mu\mu} C_{\gamma\gamma} \left[ \ln \frac{\mu^2}{m_\mu^2} + \delta_2 + 2 - h_2\left(\frac{m_a^2}{m_\mu^2}\right) \right] - \frac{\alpha}{2\pi} \frac{1 - 4s_w^2}{s_w c_w} c_{\mu\mu} C_{\gamma Z} \left( \ln \frac{\mu^2}{m_Z^2} + \delta_2 + \frac{3}{2} \right) \right\}$$

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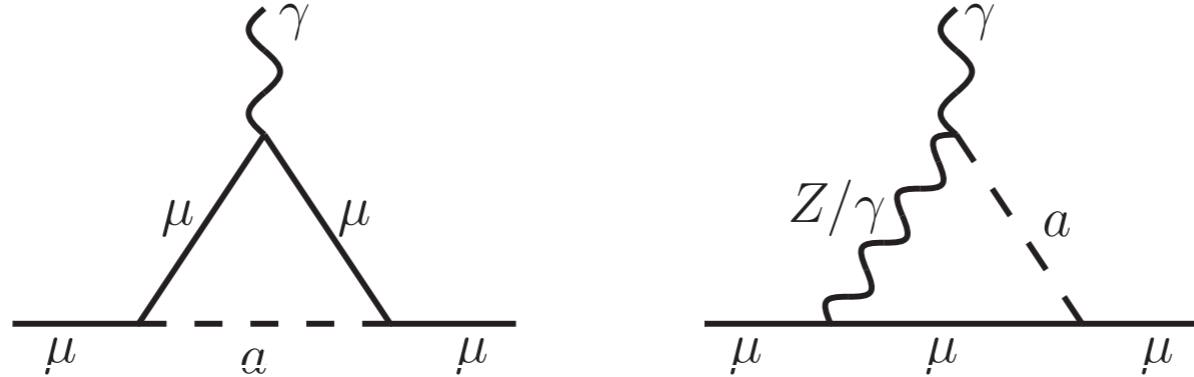
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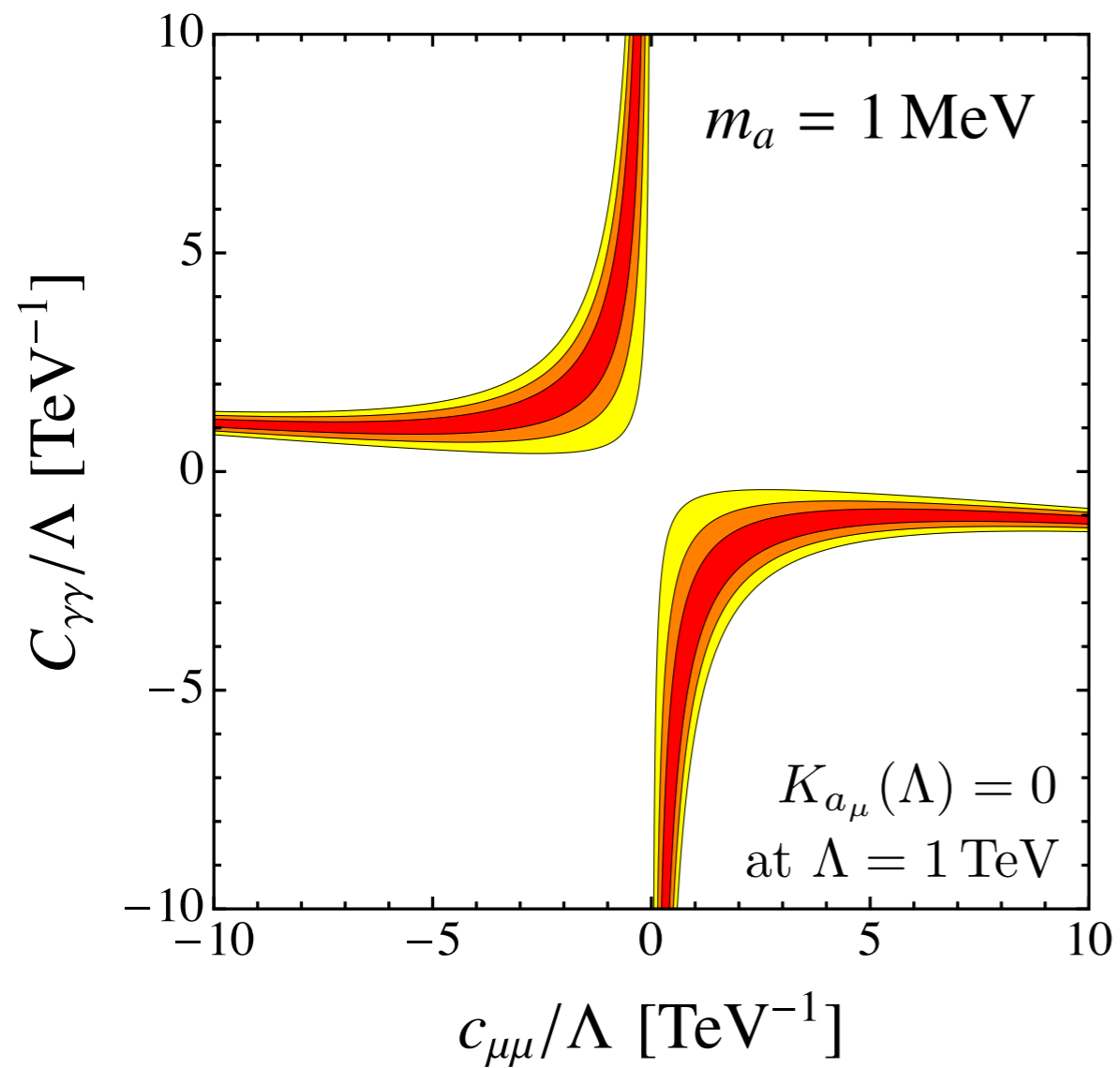
$$\mathcal{L}_{\text{eff}}^{D=6} \ni -K_{a\mu} \frac{em_\mu}{4\Lambda^2} \bar{\mu} \sigma_{\mu\nu} F^{\mu\nu} \mu$$

$$h_1(0) = 1 \quad h_1(x) \approx (2/x)(\ln x - \frac{11}{6}) \text{ for } x \gg 1$$

$$h_2(0) = 0 \quad h_2(x) \approx (\ln x + \frac{1}{2})$$

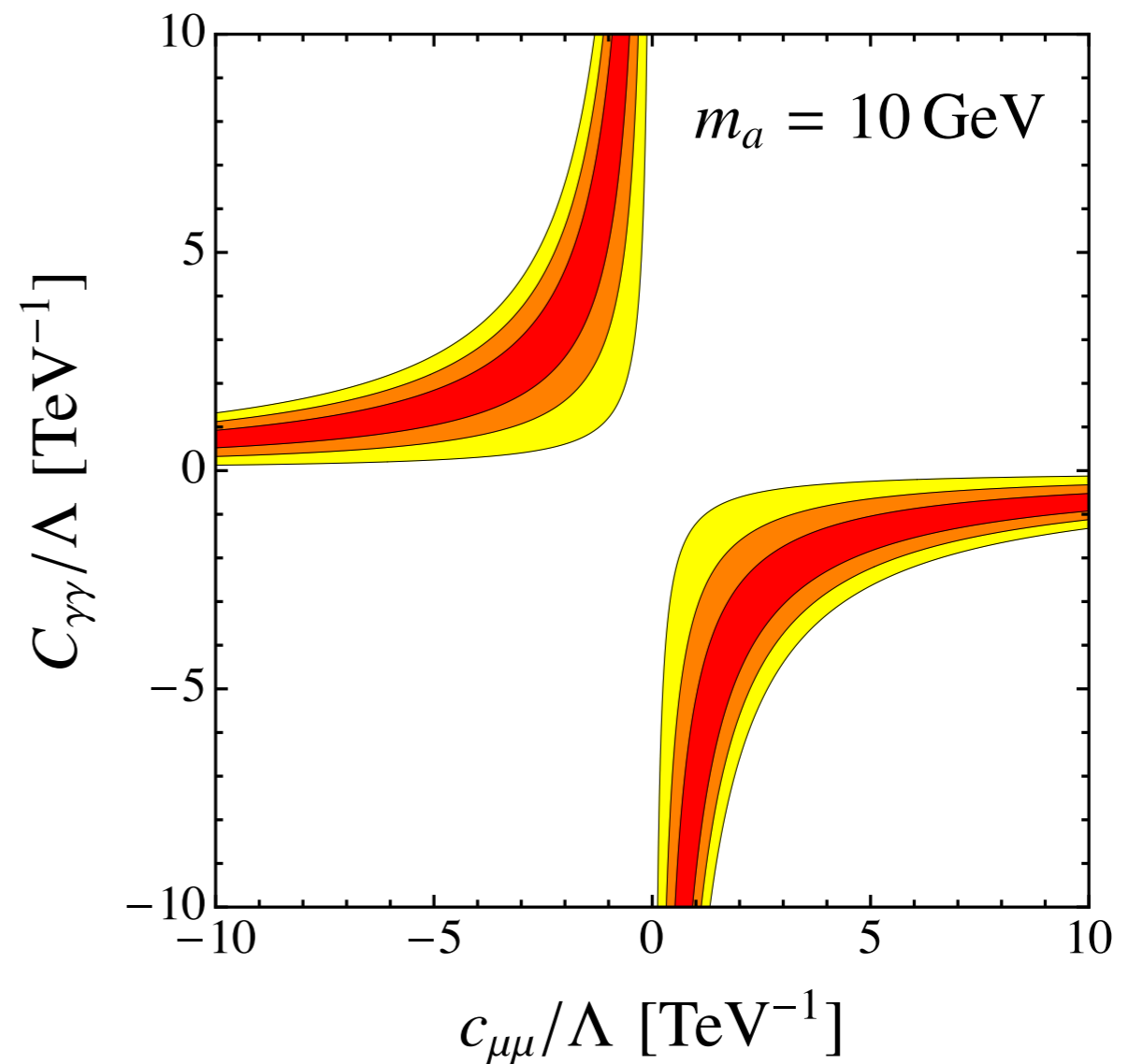
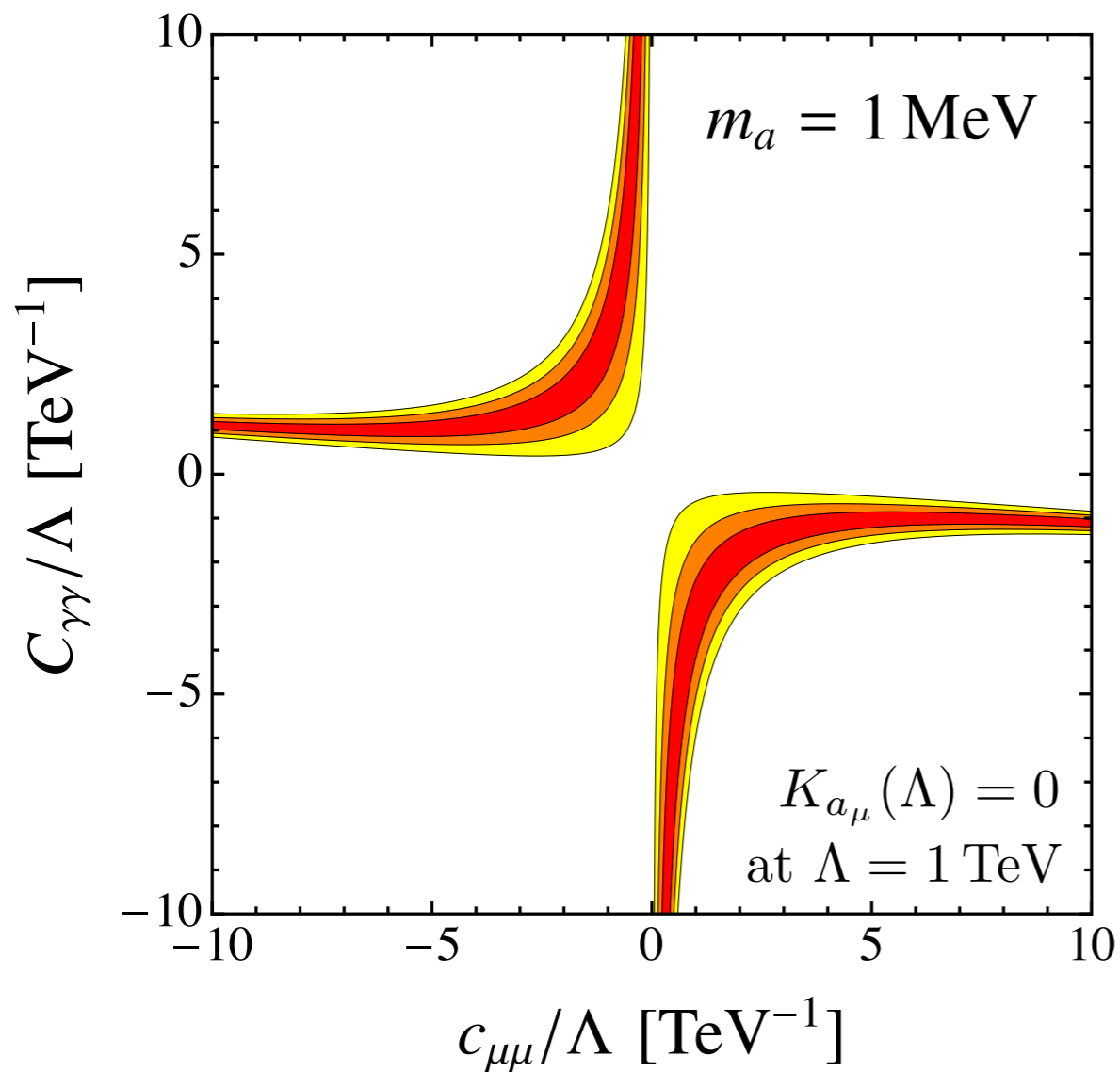
# Muon $(g - 2)_\mu$

- Allowed parameter space



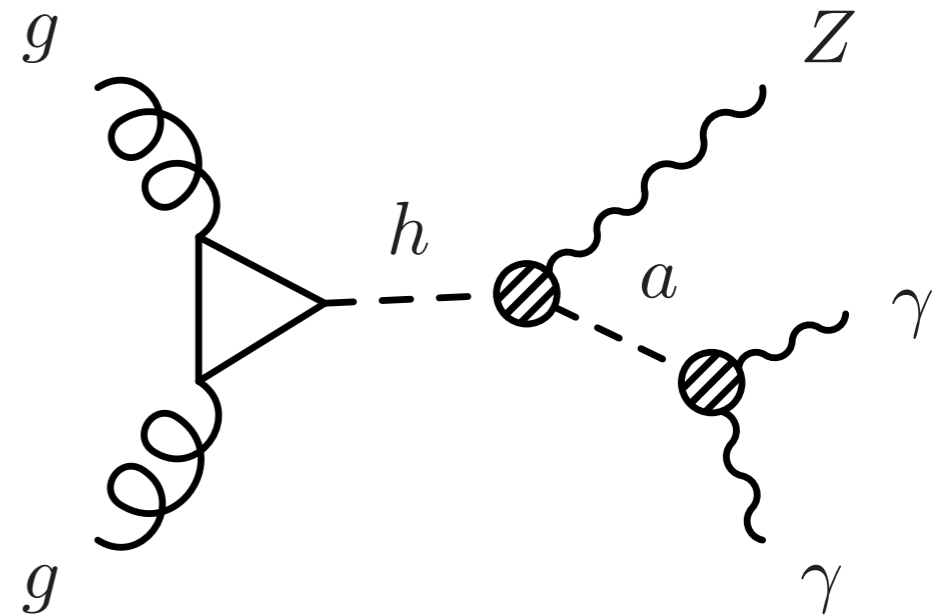
# Muon $(g - 2)_\mu$

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# Conclusions

- Rare Higgs decays provide a powerful way to probe the existence of ALPs with masses between 30 MeV and 60 GeV and couplings suppressed by the 1 - 100 TeV scale
- Connection to low-energy physics probes such as  $(g - 2)_\mu$



# Outlook

- Dedicated analyses with reconstruction efficiencies and exploiting displaced-vertex signatures
- Investigating the flavour sector
- Looking at various anomalies

*[Bauer, Neubert, Thamm: to appear]*

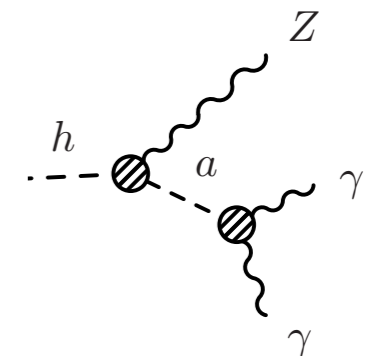
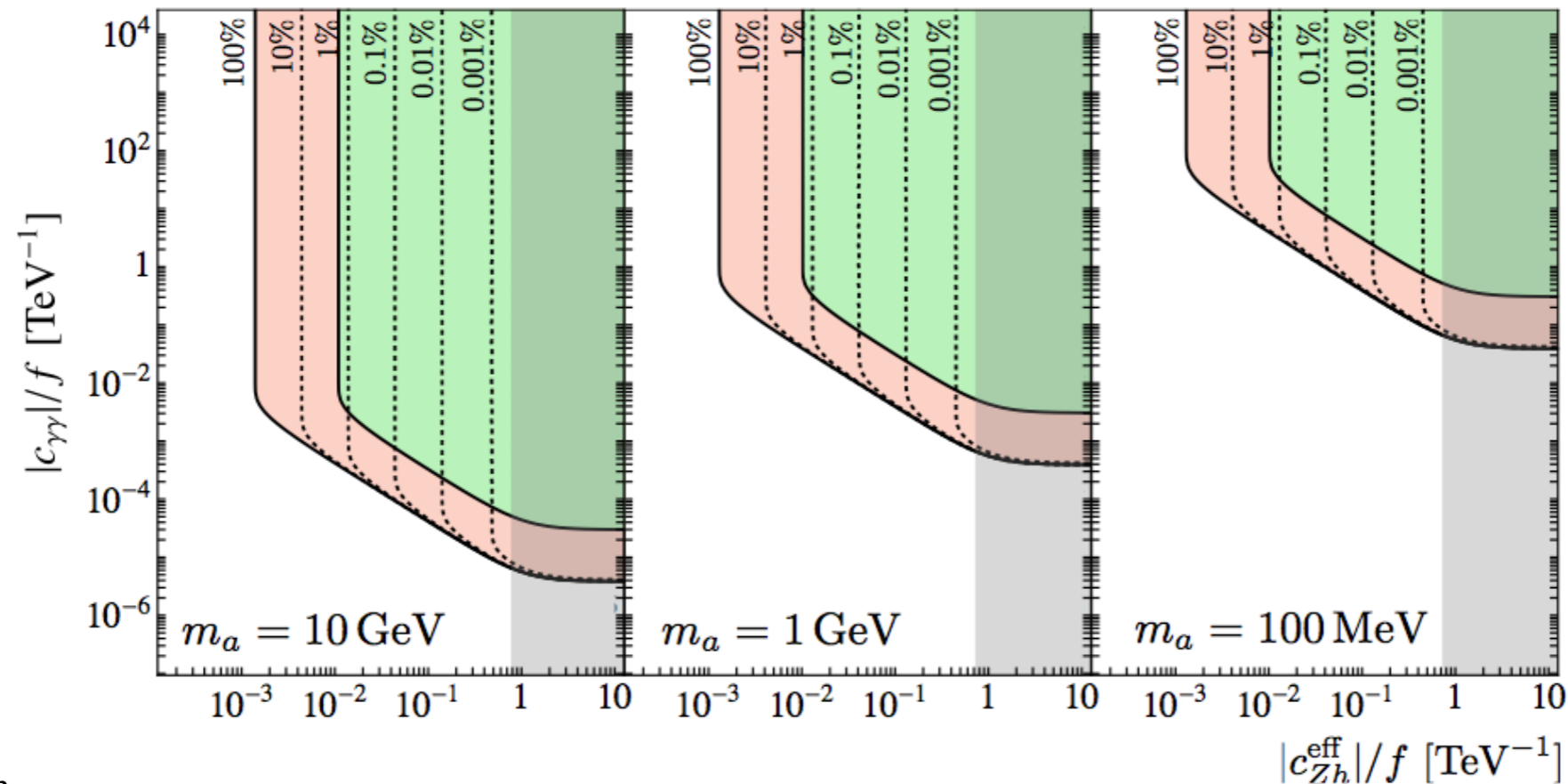
# Backup

# Parameter space at the FCC-ee

[arXiv:1308.6176]

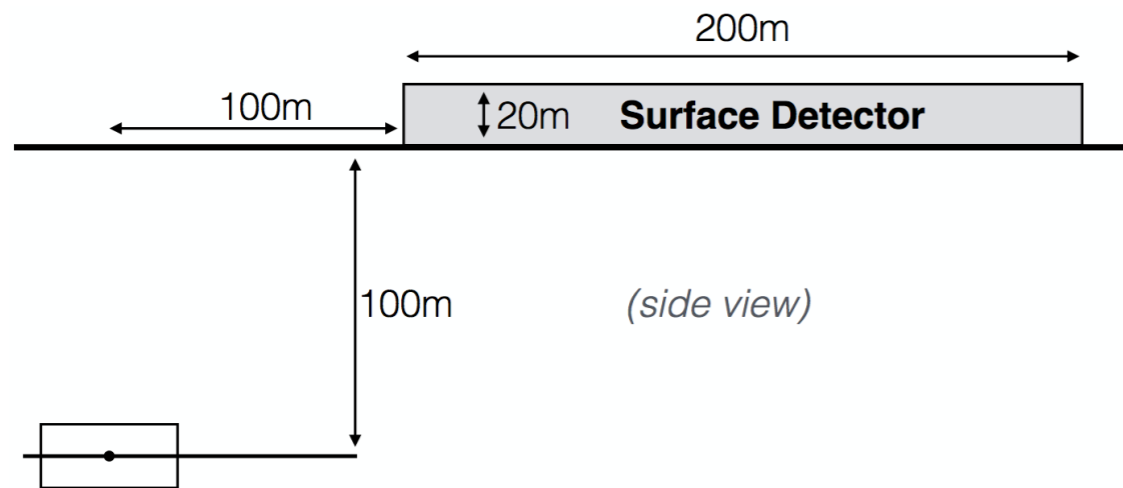


- $e^+e^-$  collider
- 240 and 350 GeV
- 3 million Higgses

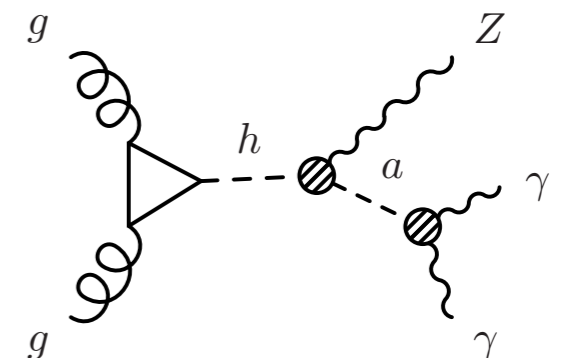
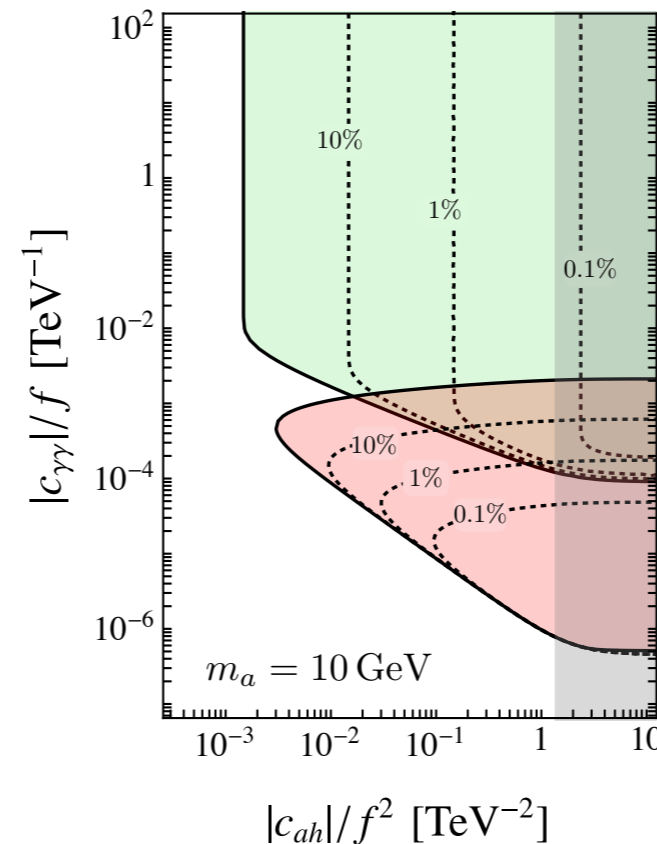
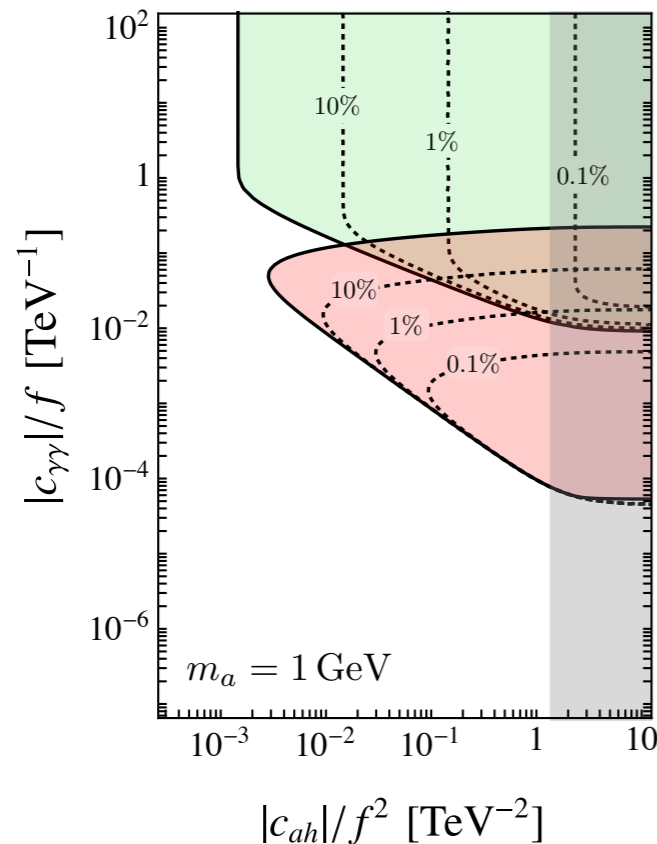


# Parameter space at MATHUSLA

[Chou, Curtin, Lubatti: 1606.06298]

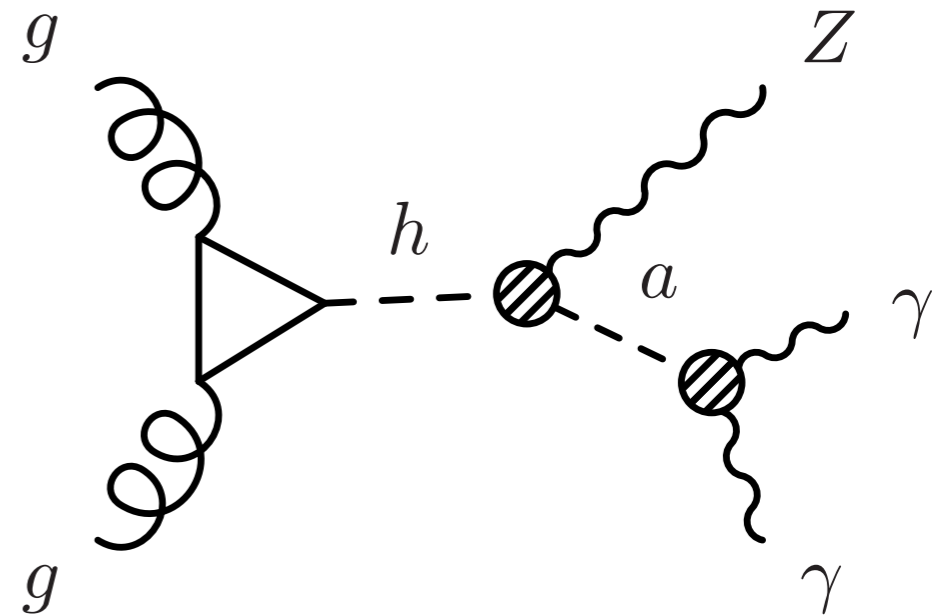


- Long-lived particles at LHC



# Outline

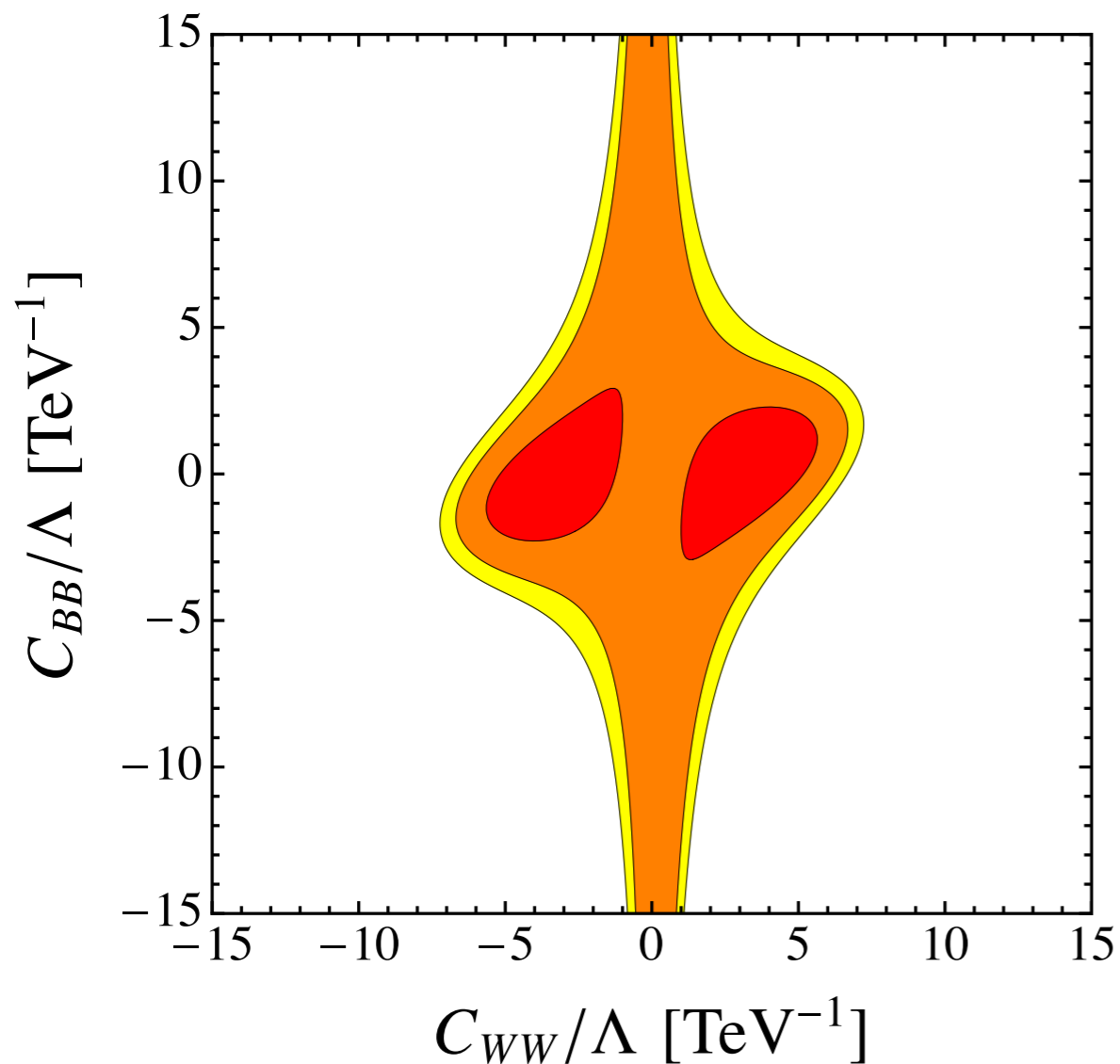
- Motivation
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  - ♦ Electroweak precision test
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# Electroweak precision tests

- Allowed parameter space for  $C_{Zh}^{(5)} = 0$

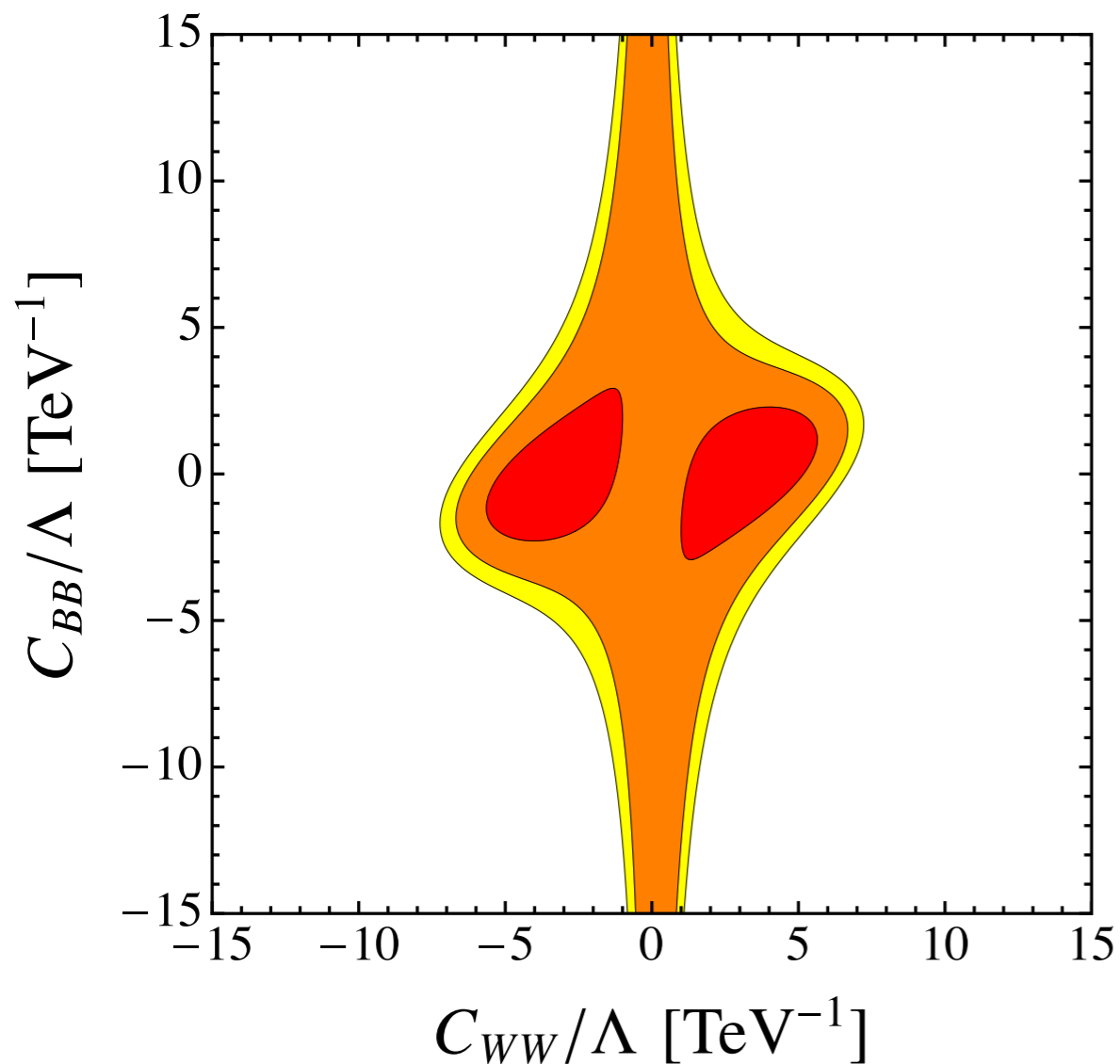
[Baak et al.: 1407.3792]



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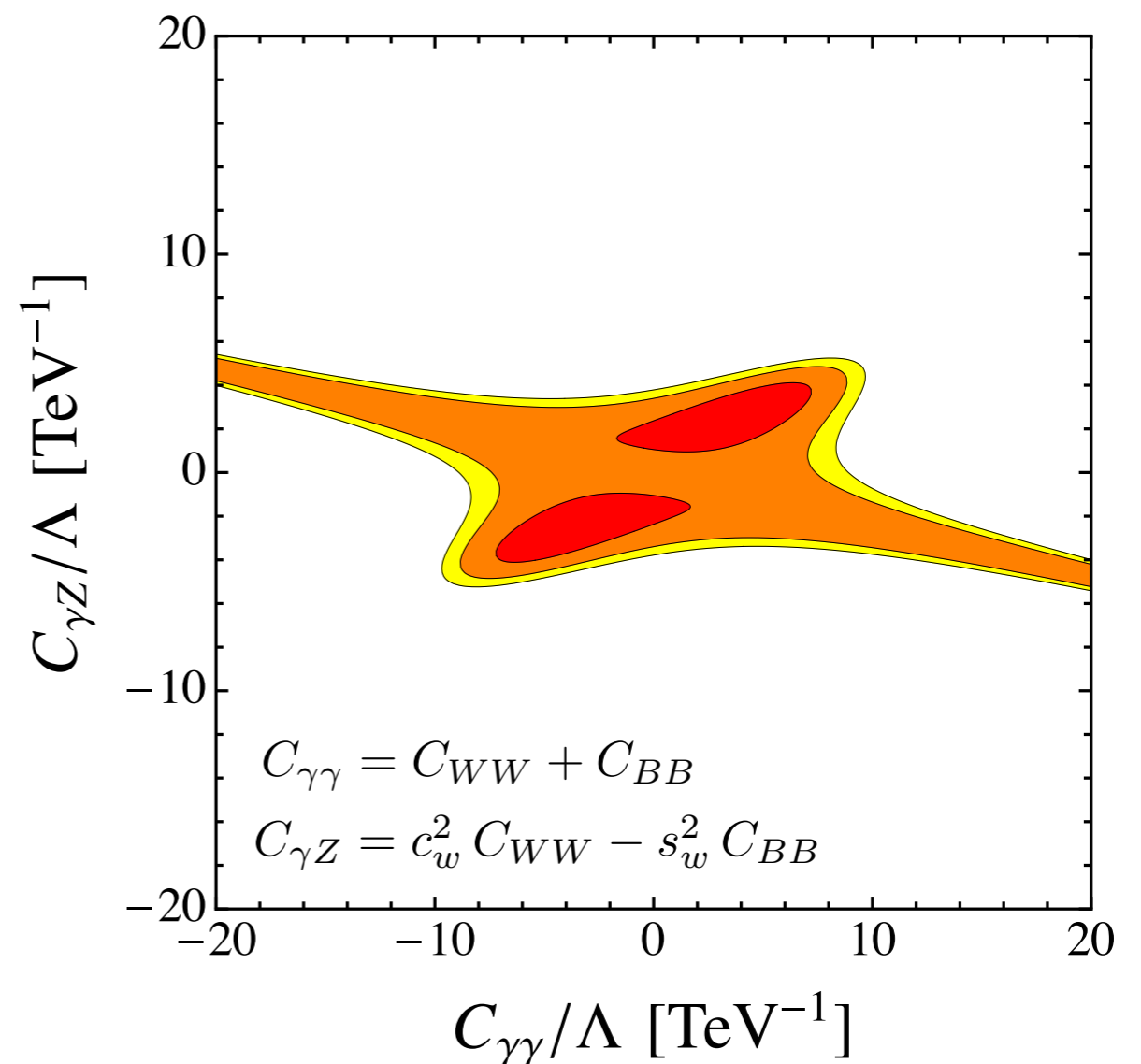
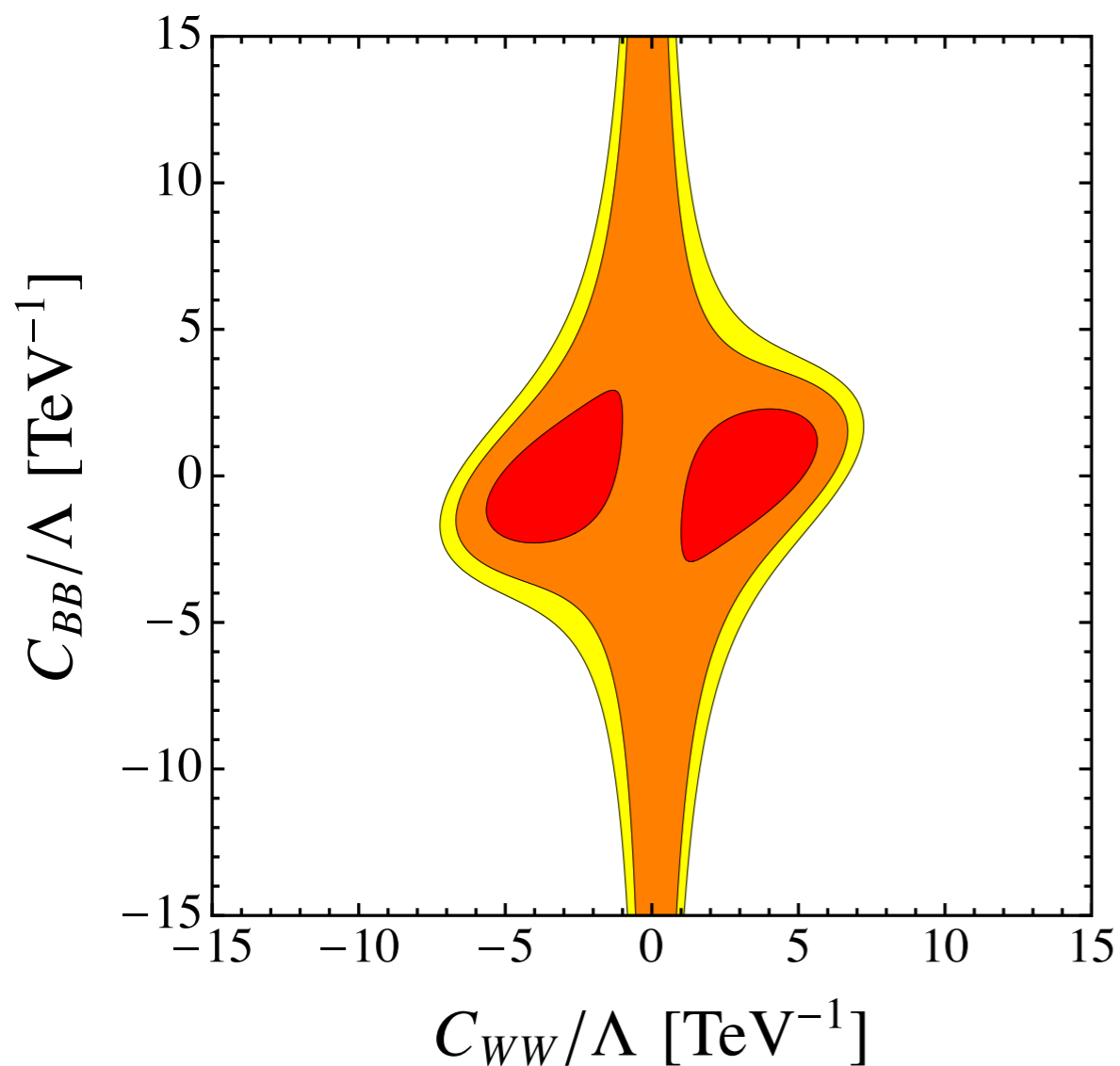
$$C_{\gamma\gamma} = C_{WW} + C_{BB}$$

$$C_{\gamma Z} = c_w^2 C_{WW} - s_w^2 C_{BB}$$

# Electroweak precision tests

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[Baak et al.: 1407.3792]

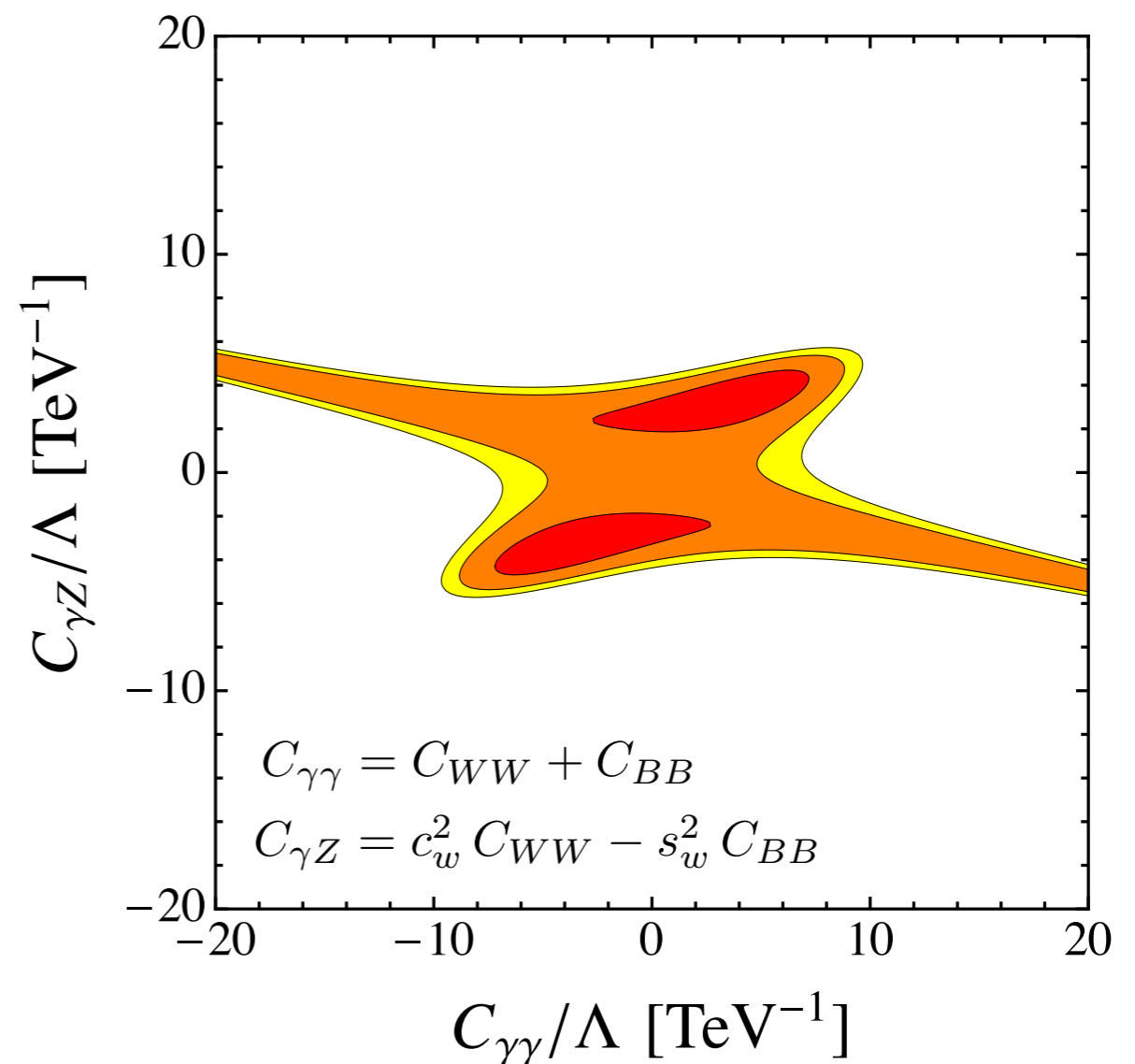
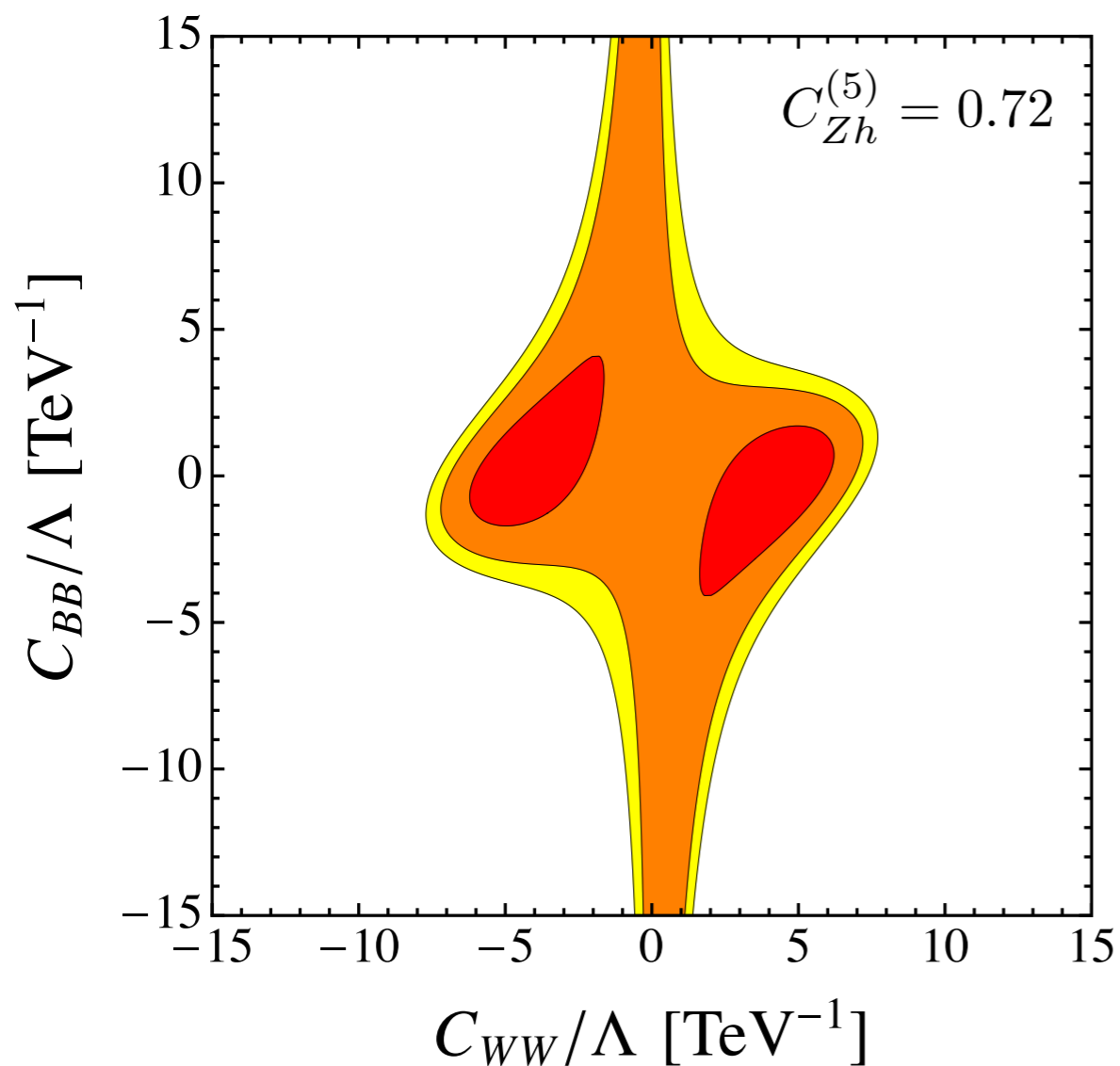




# Electroweak precision tests

- Allowed parameter space for  $C_{Zh}^{(5)} \neq 0$

[Baak et al.: 1407.3792]



# Electroweak precision tests

- Measurement of OPAL at per-cent level
- Compatible with  $C_{WW}$  and  $C_{BB}$  of order  $\sim 30$

*[Abbiendi et al.: 0309052]*

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- Measurement of OPAL at per-cent level
- Compatible with  $C_{WW}$  and  $C_{BB}$  of order  $\sim 30$  *[Abbiendi et al.: 0309052]*
- FCC-ee expectation of  $10^{-5}$  uncertainty *[Janot: 1512.05544]*  
*[Blas, Cuichini, Franco, Mishima, Pierini, Reina, Silvestrini: 1608.01509]*

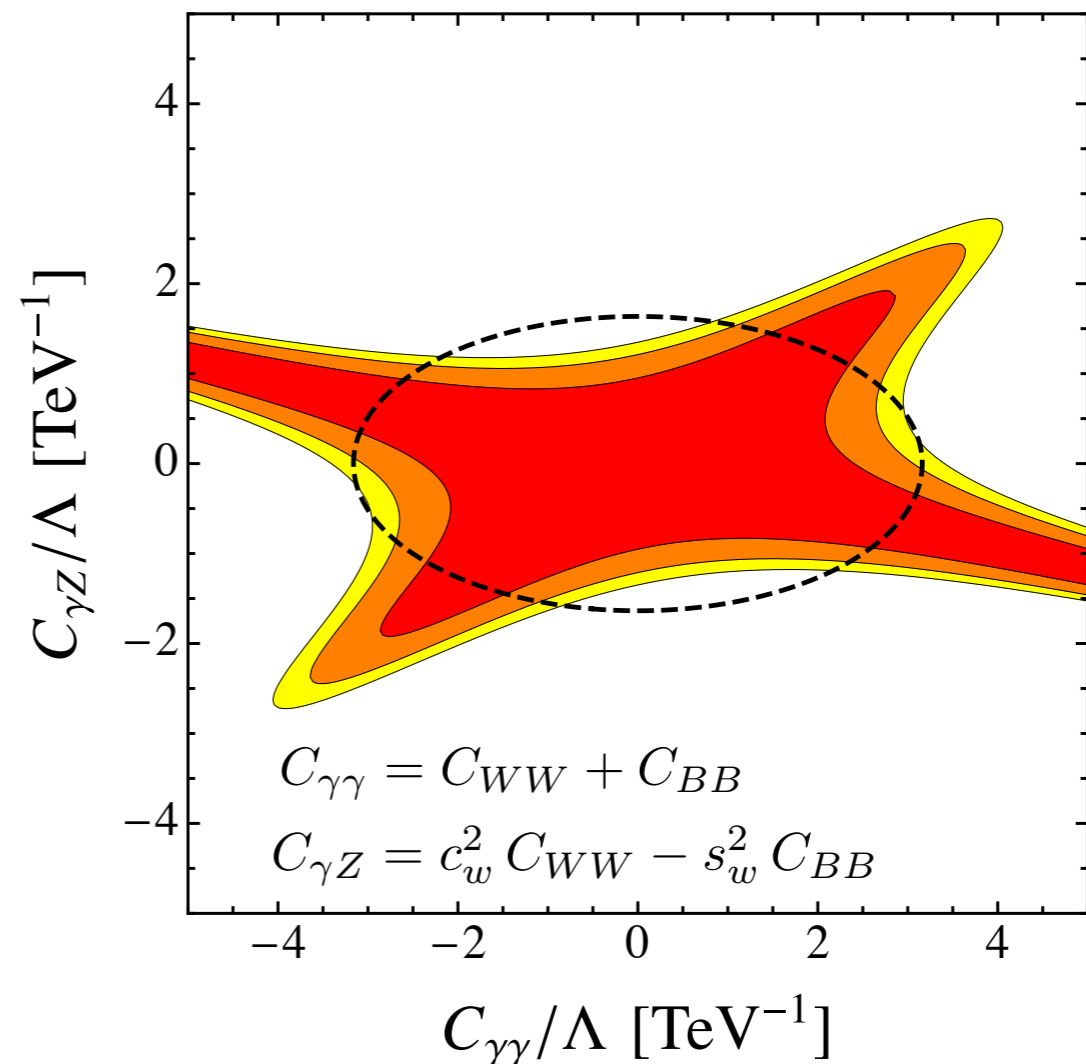
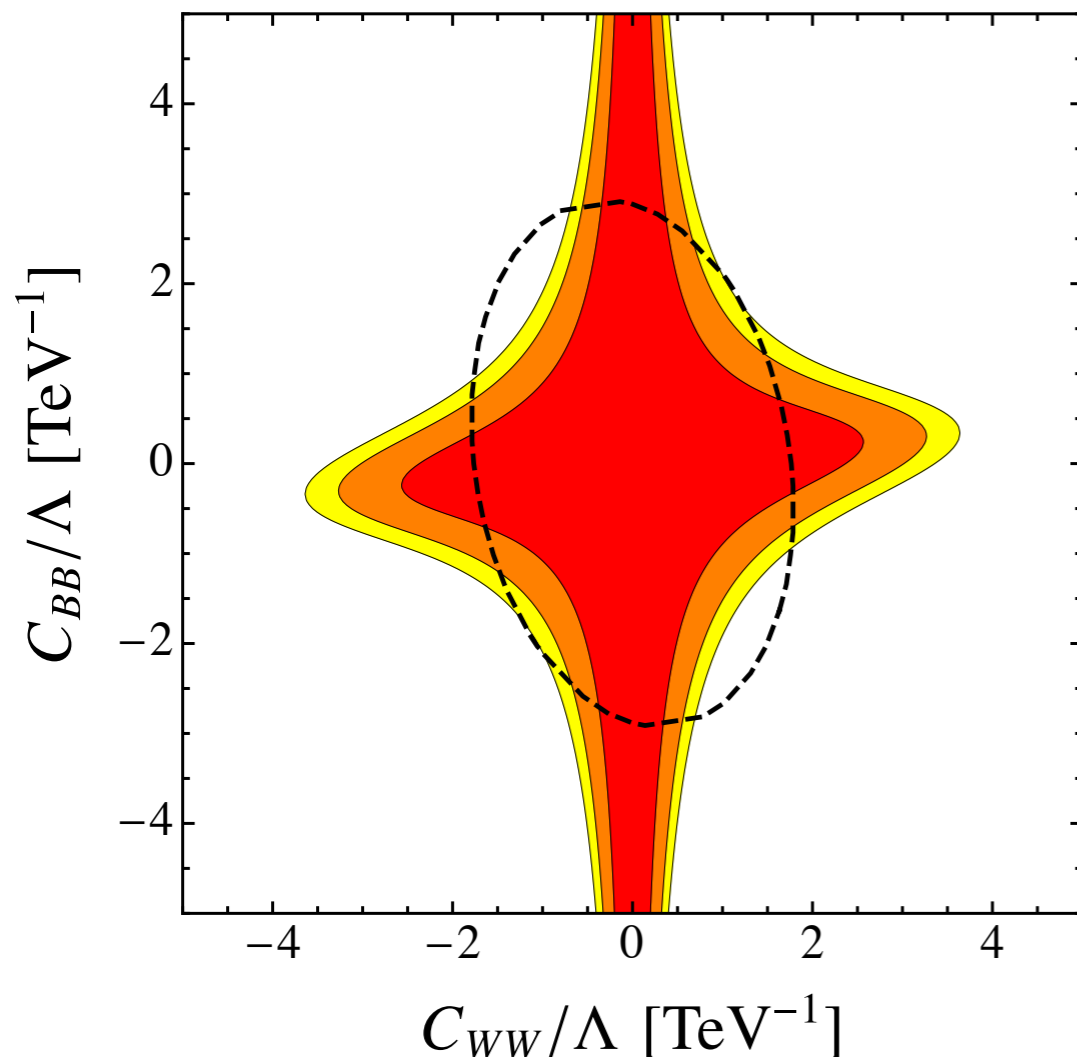
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[Abbiendi et al.: 0309052]

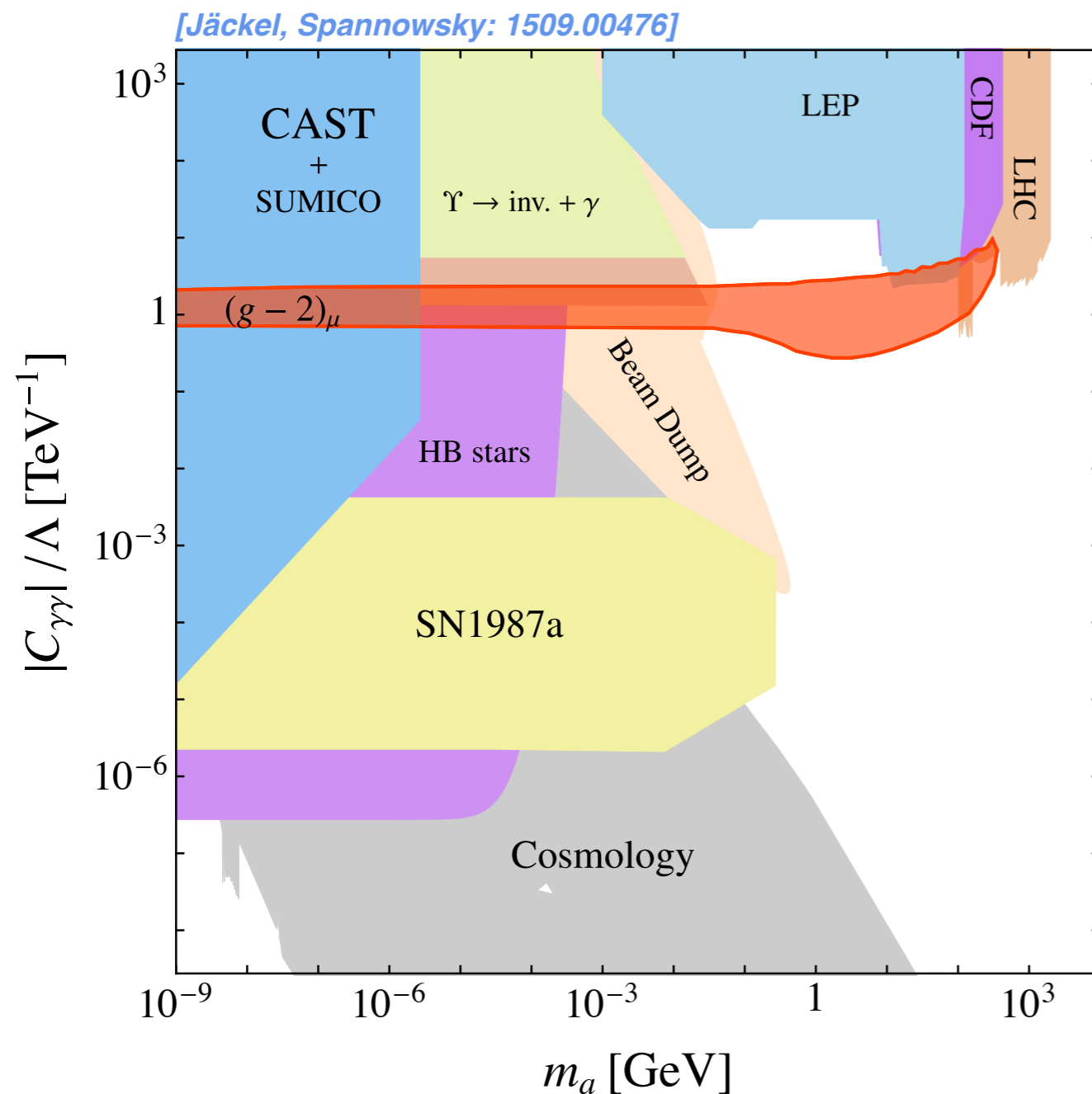
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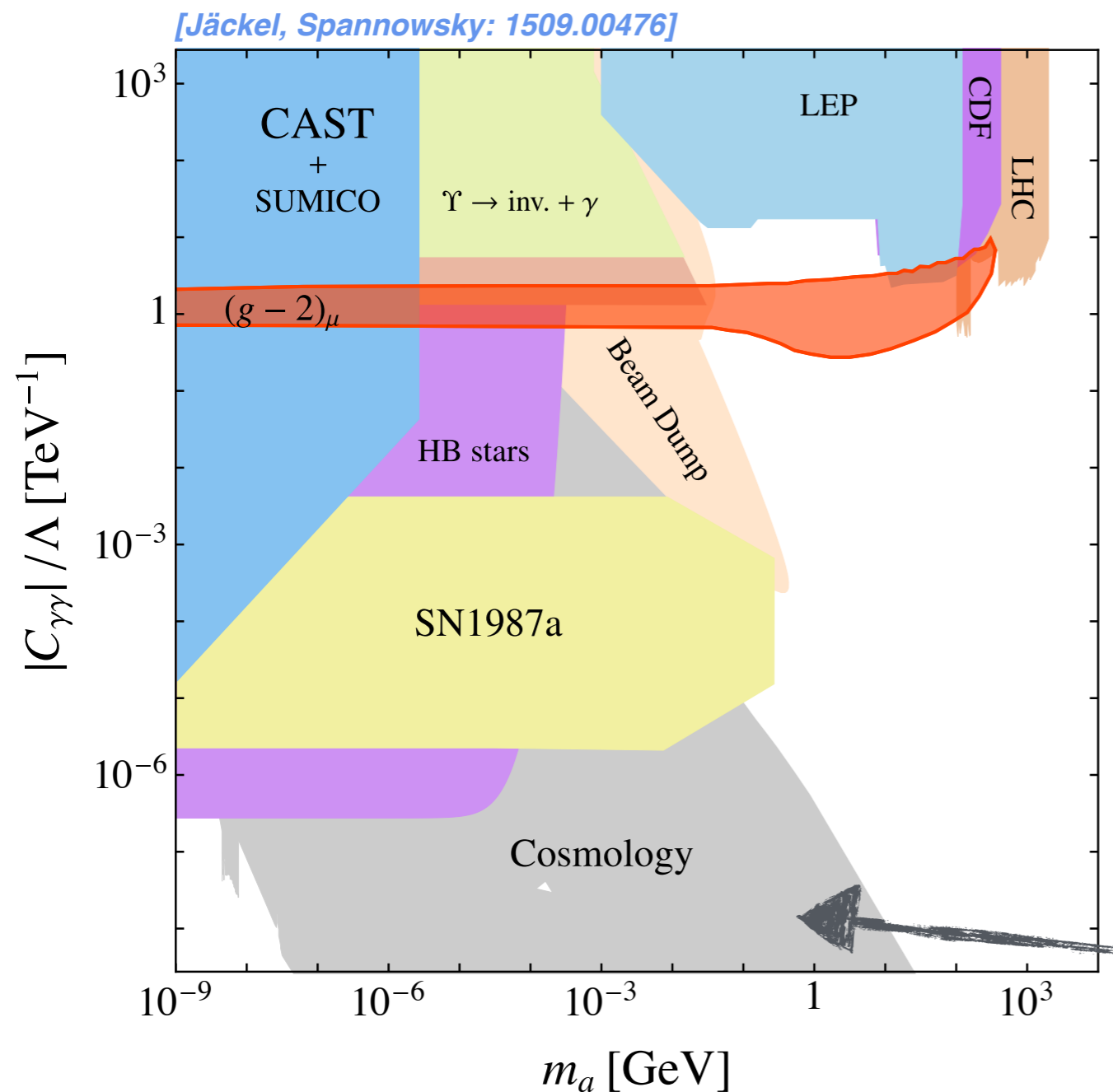
# Probing the parameter space

- Constraints on ALP mass and coupling to photons



# Probing the parameter space

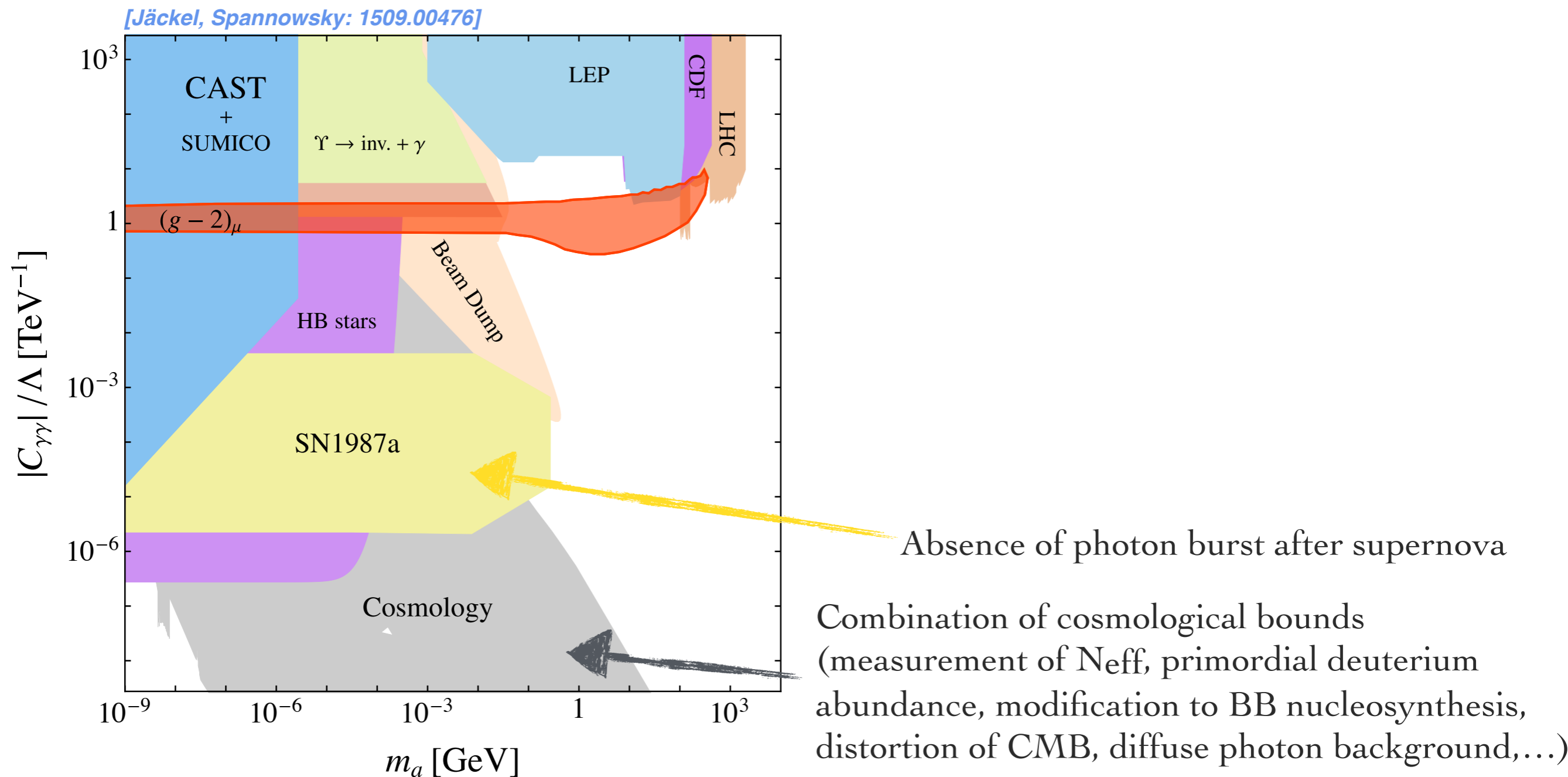
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Combination of cosmological bounds  
(measurement of  $N_{\text{eff}}$ , primordial deuterium abundance, modification to BB nucleosynthesis, distortion of CMB, diffuse photon background,...)

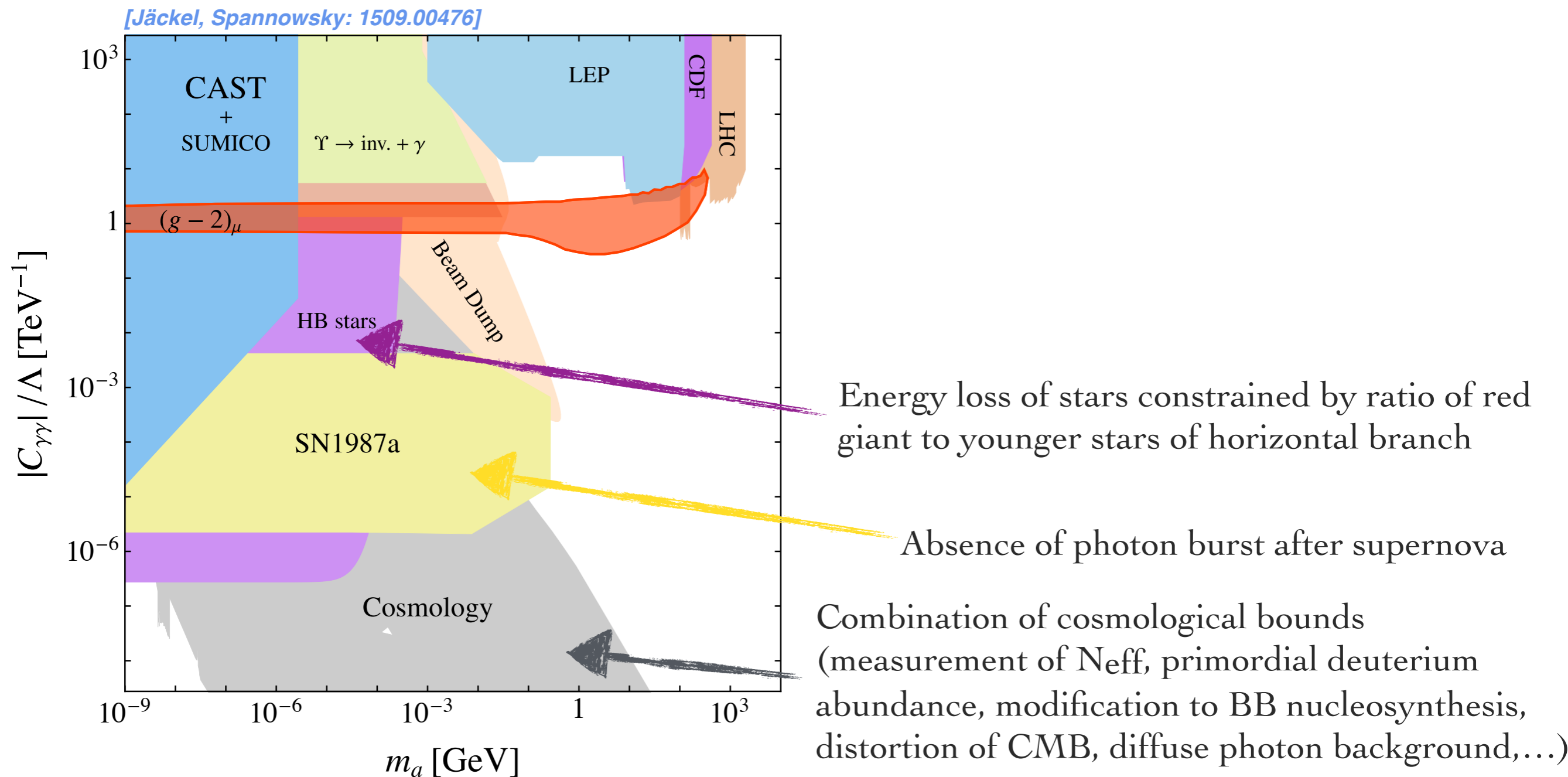
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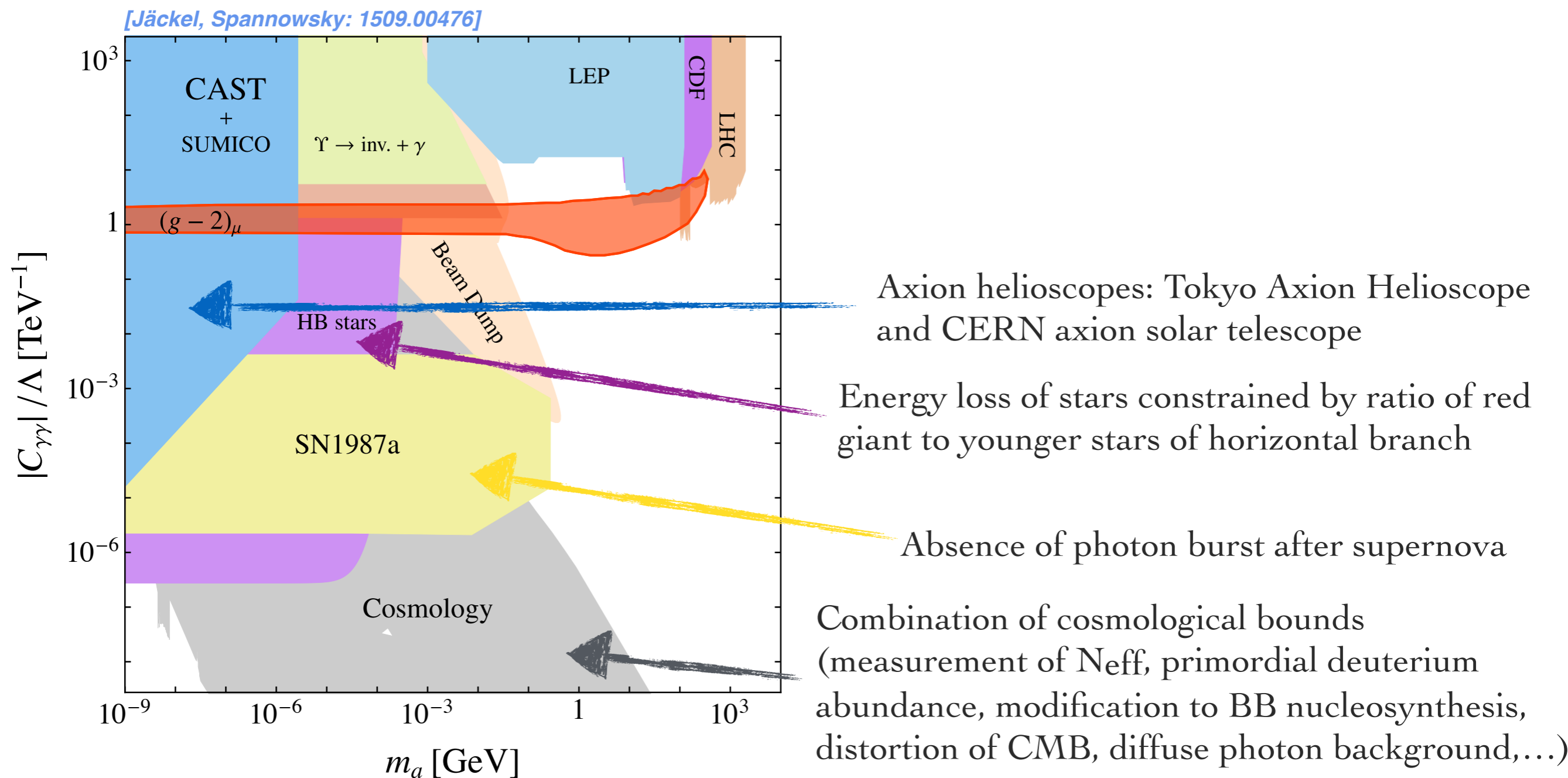
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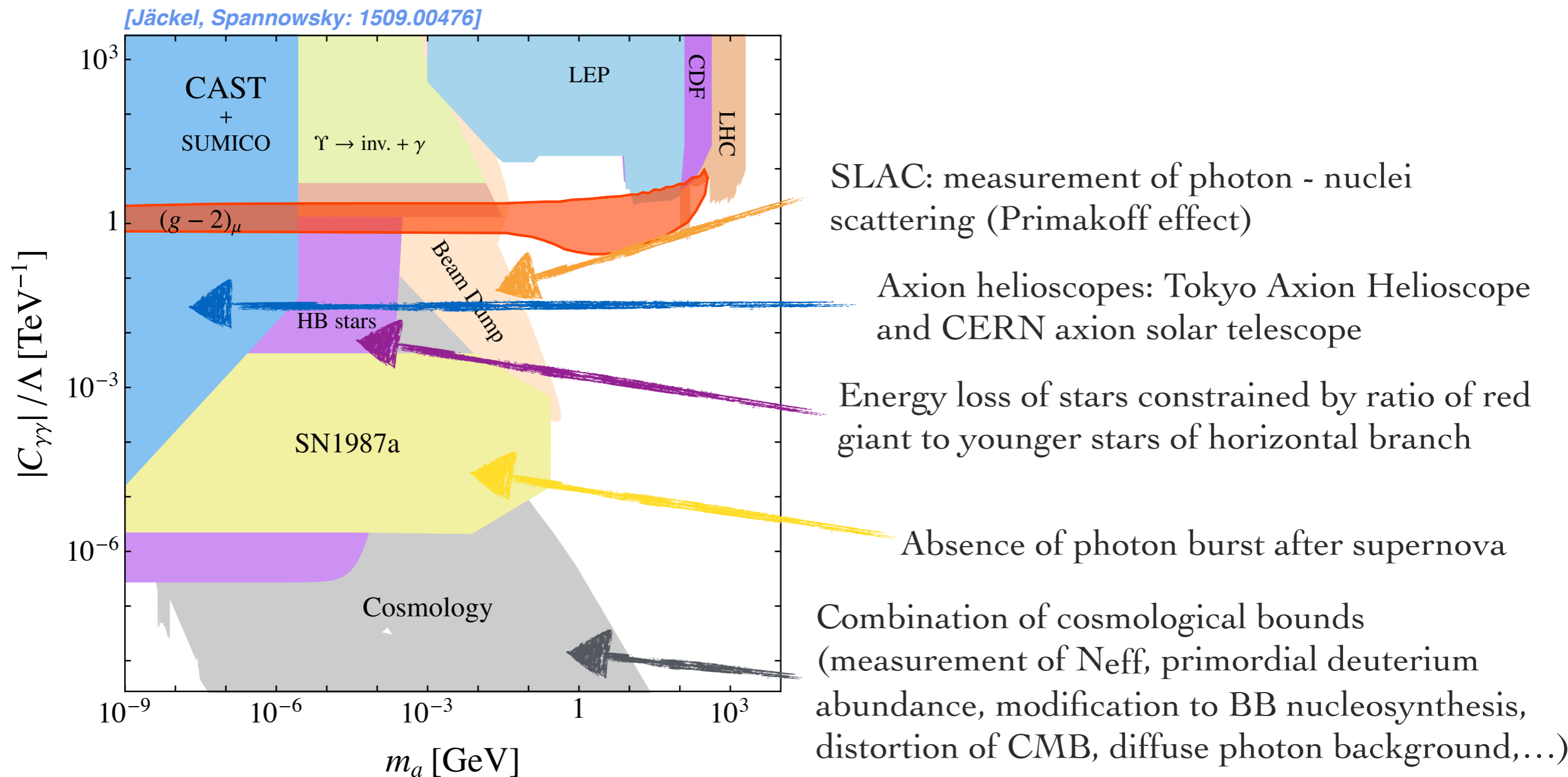
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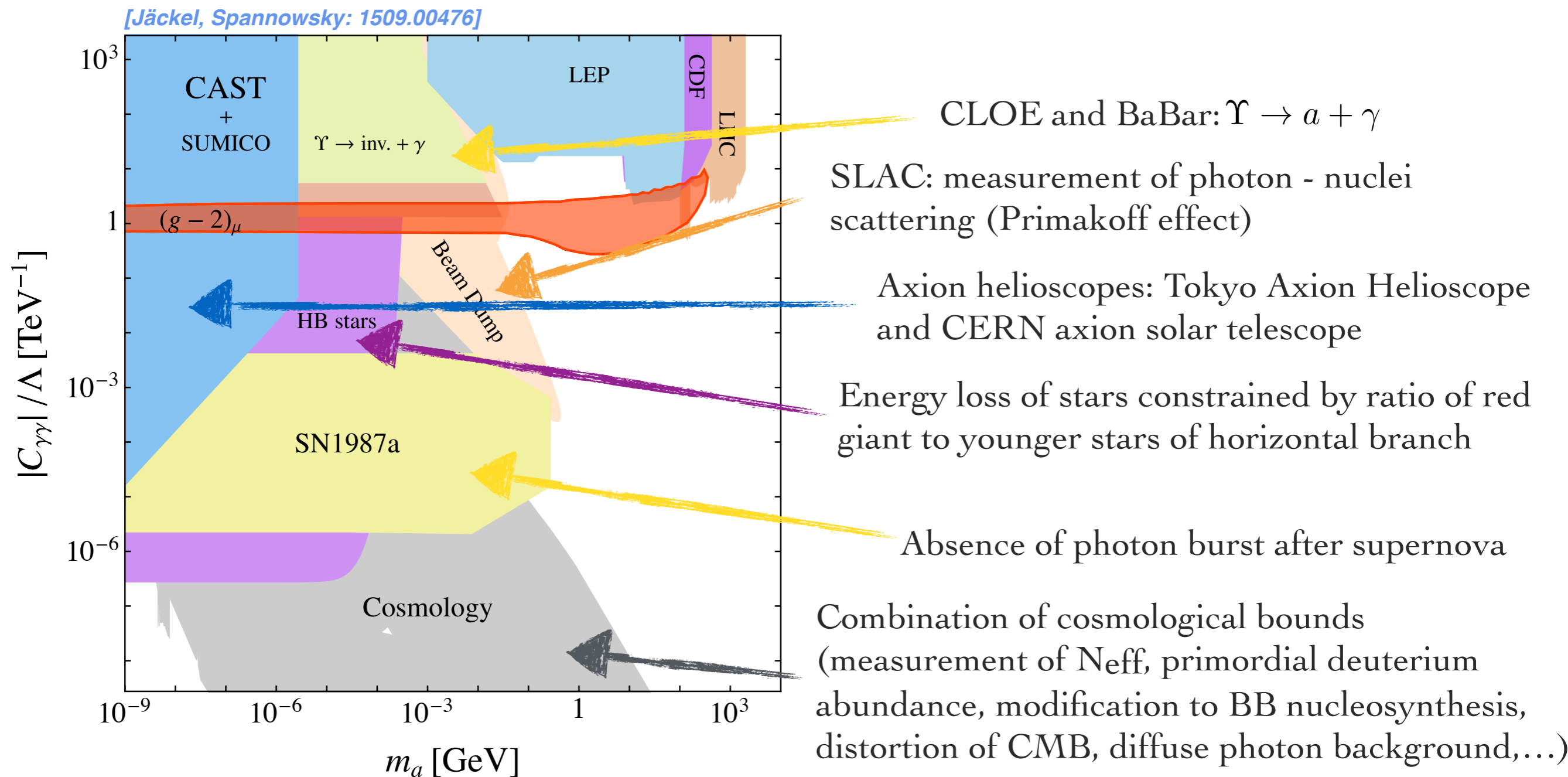
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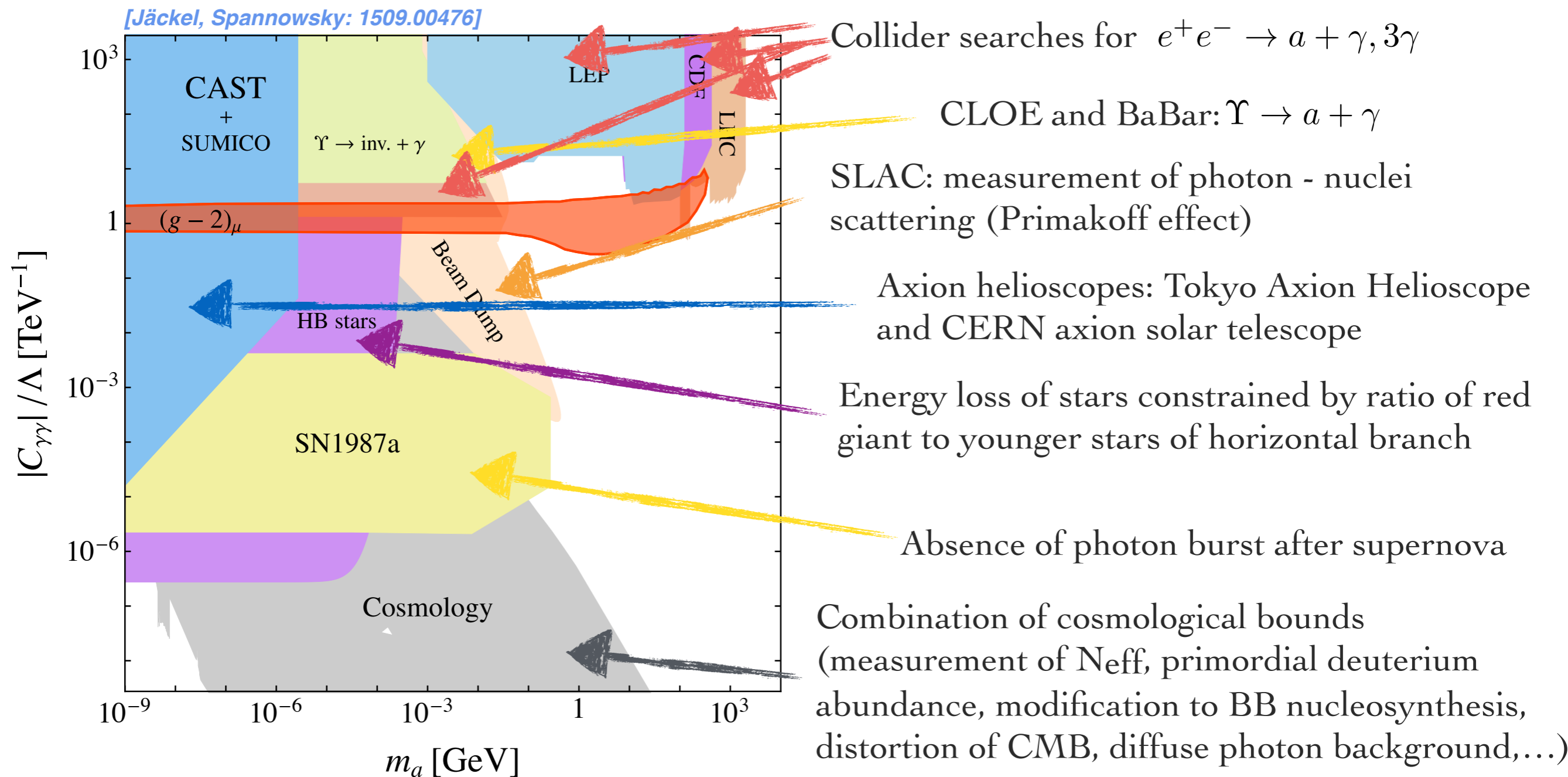
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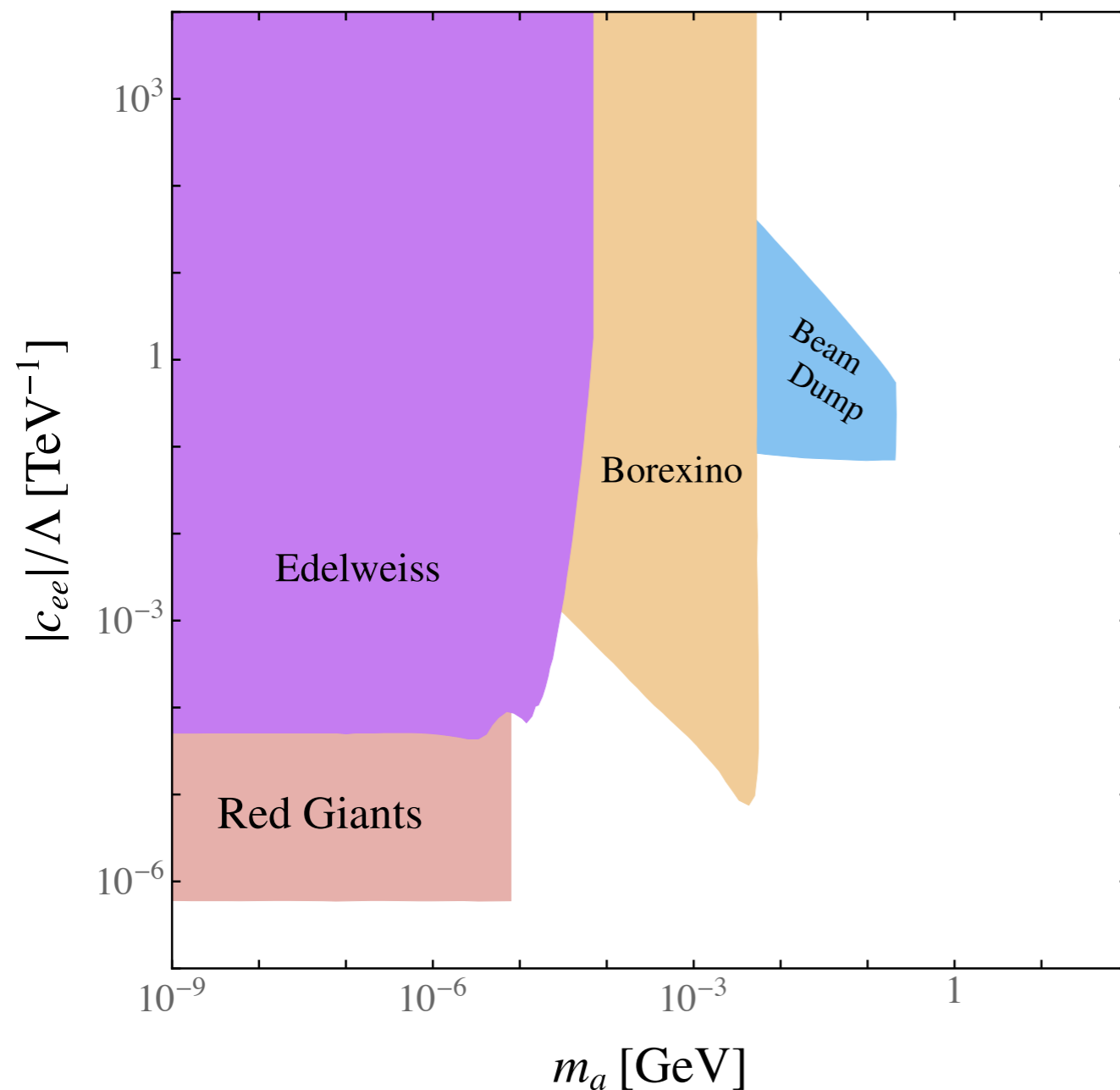
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# Probing the parameter space

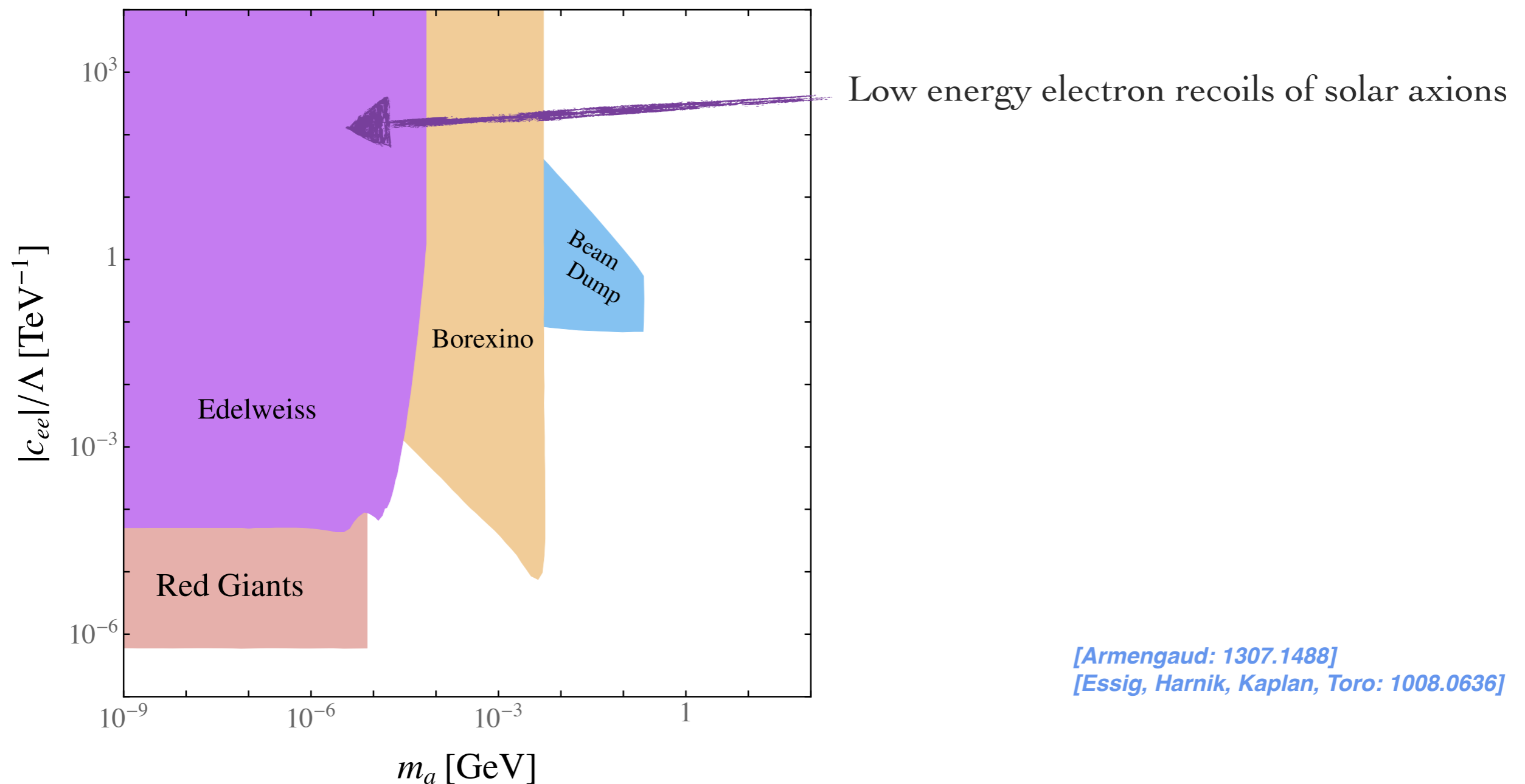
- Constraints on ALP mass and coupling to electrons



[Armengaud: 1307.1488]  
[Essig, Harnik, Kaplan, Toro: 1008.0636]

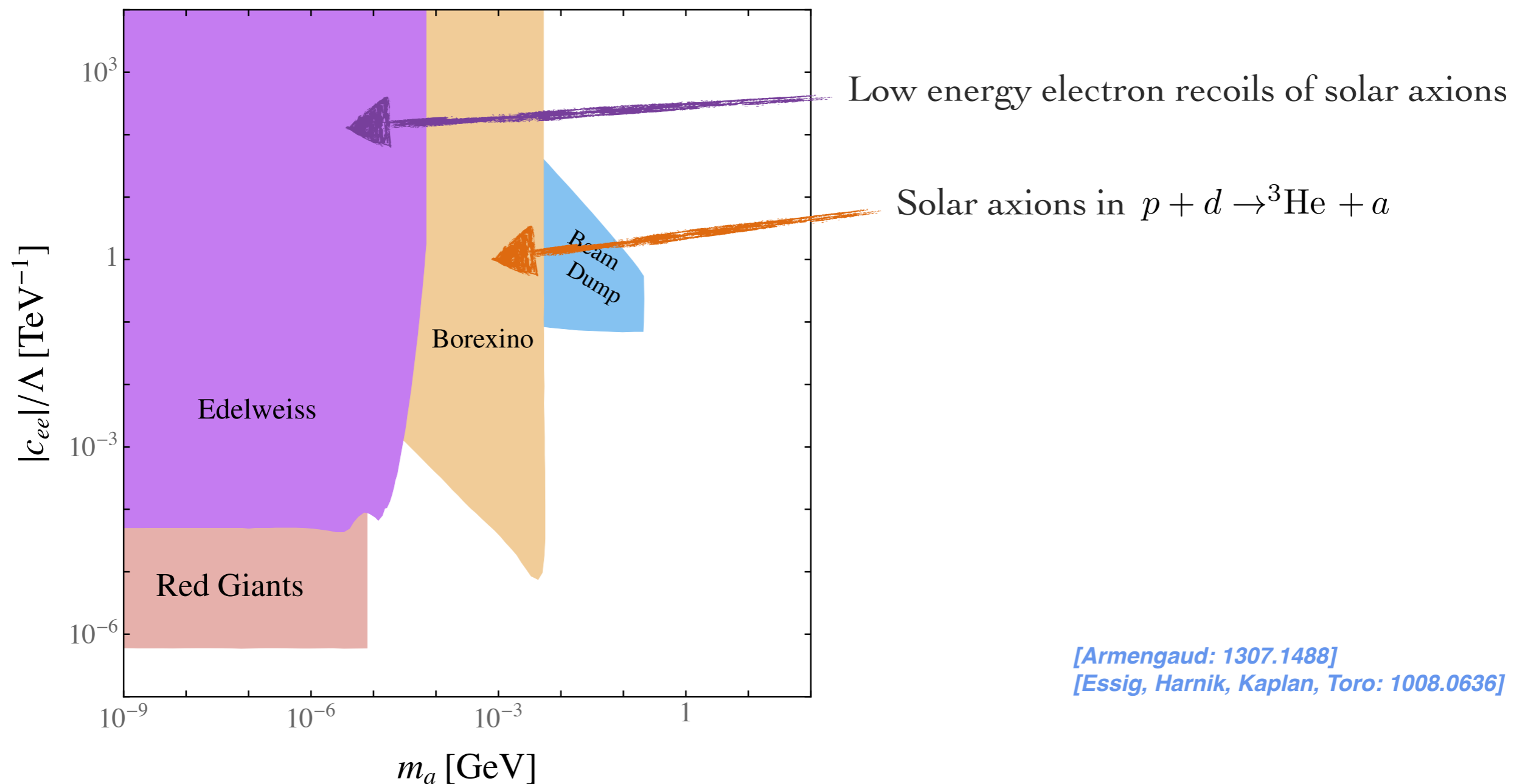
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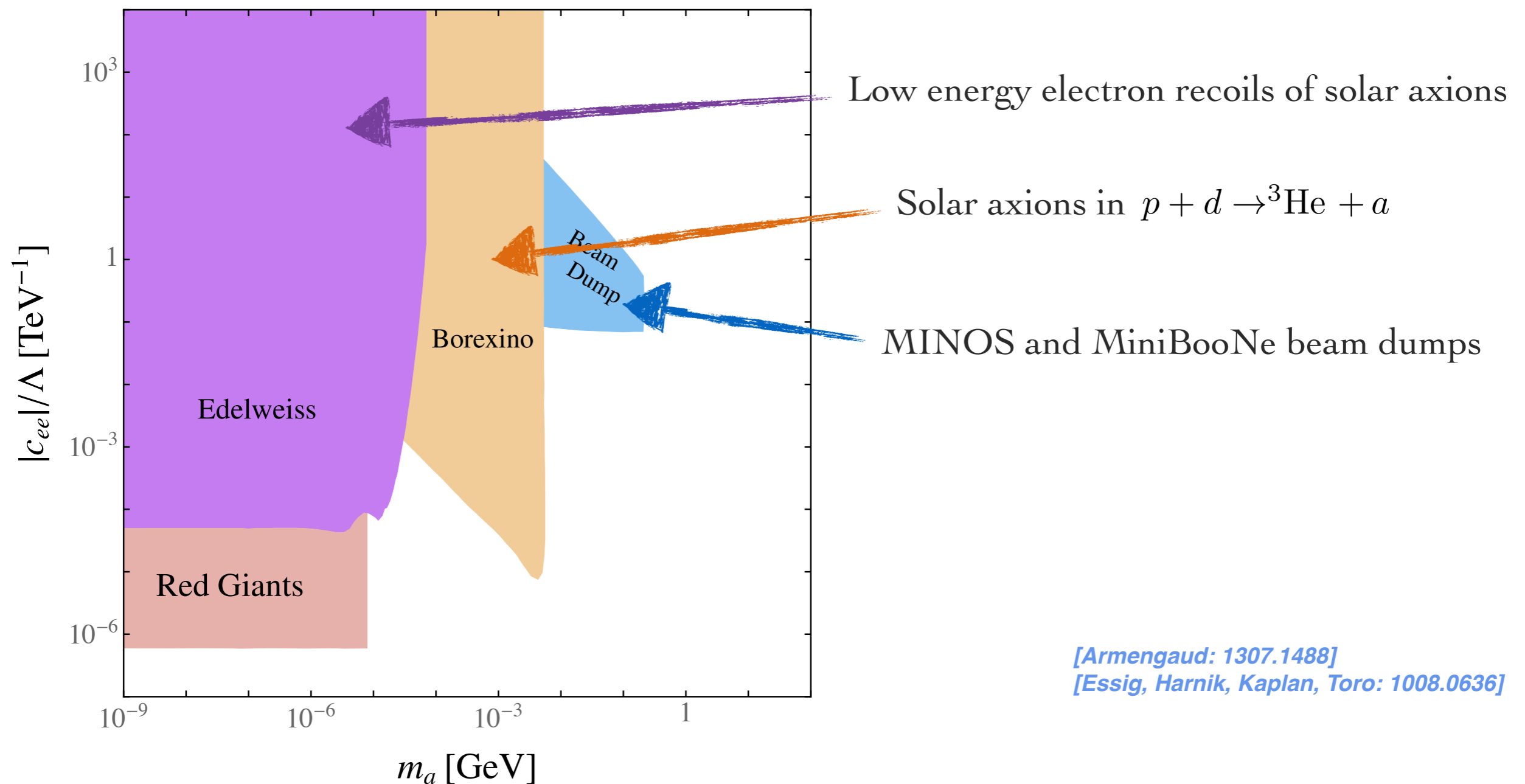
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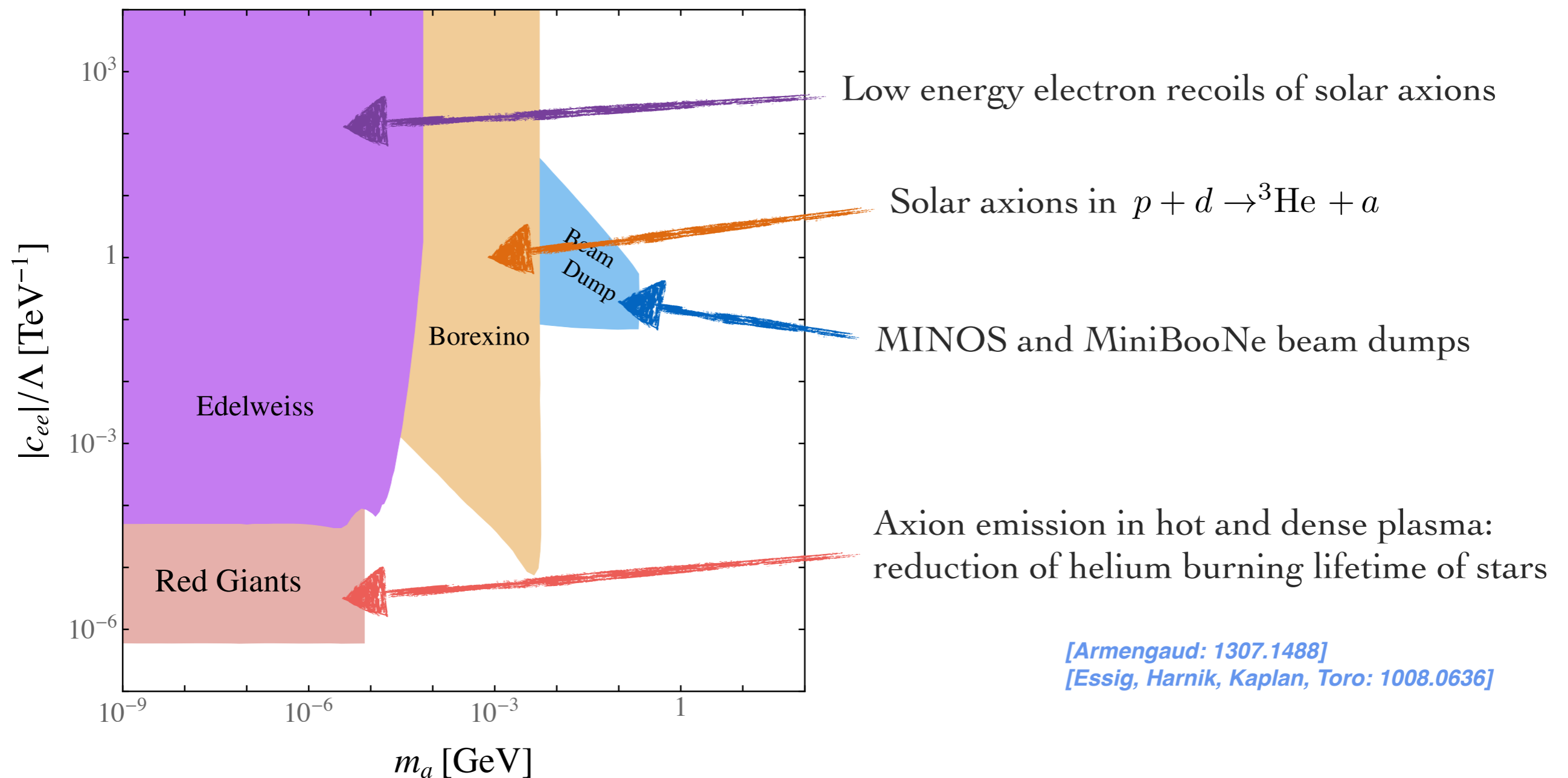
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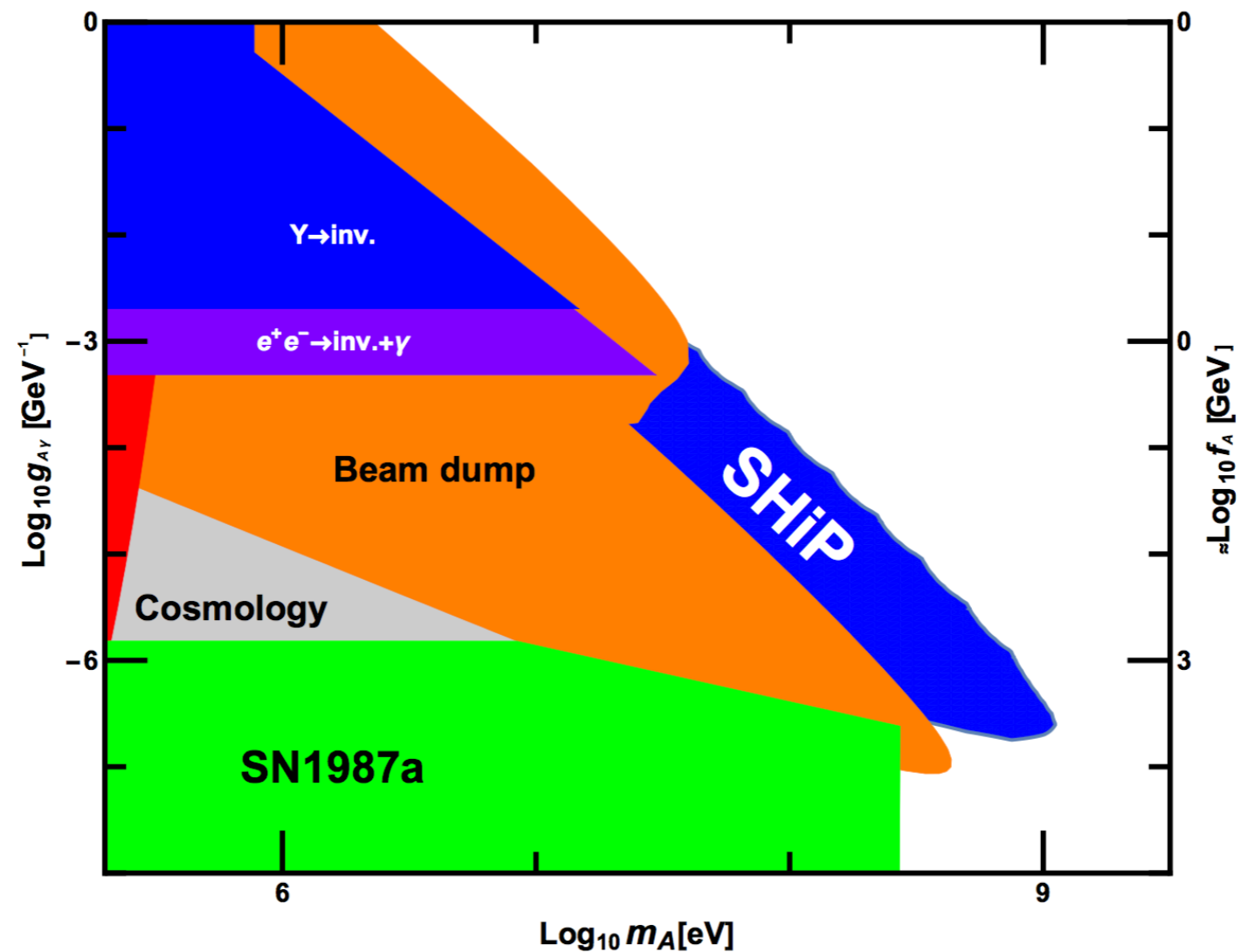
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# SHiP expected reach

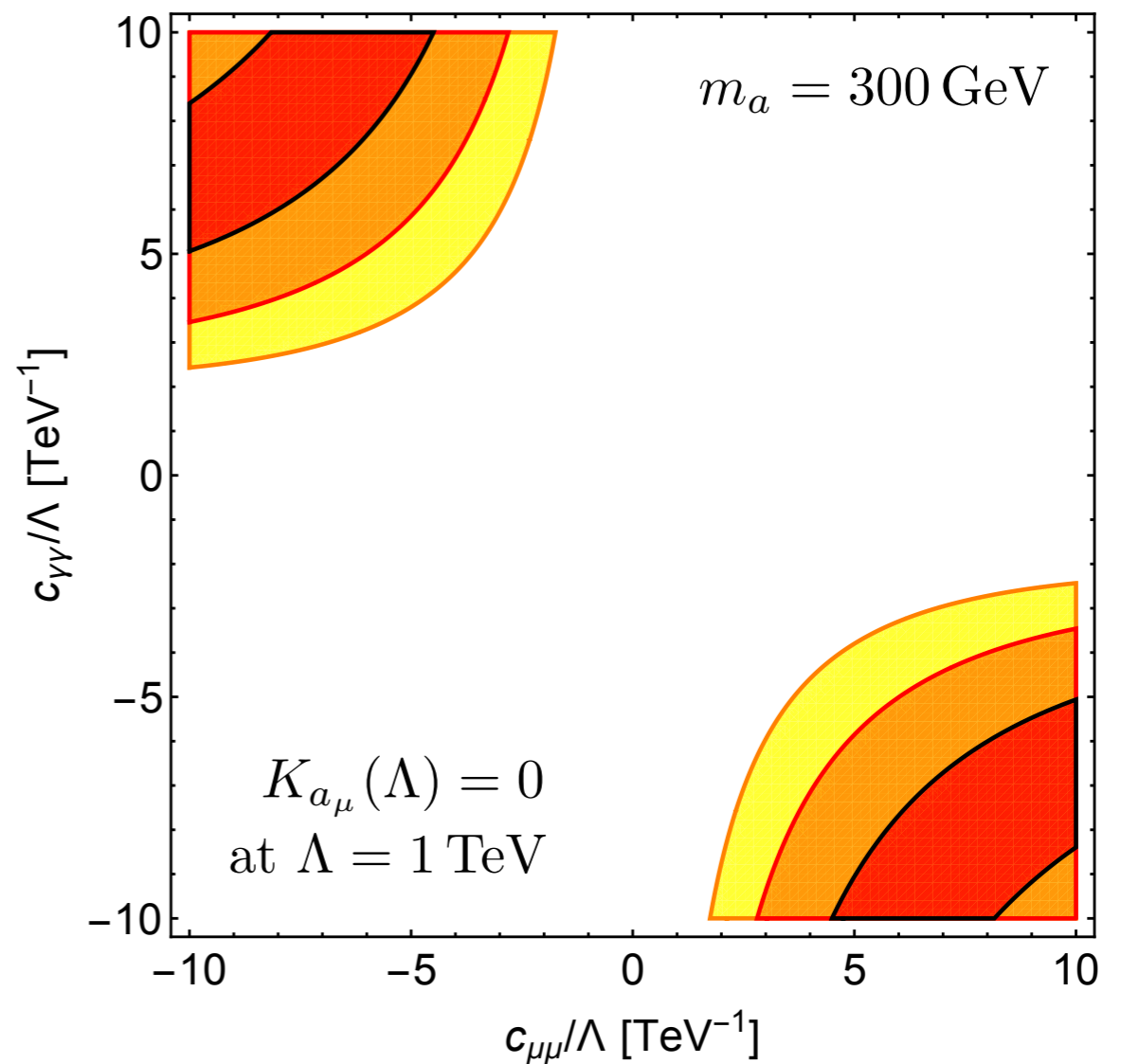
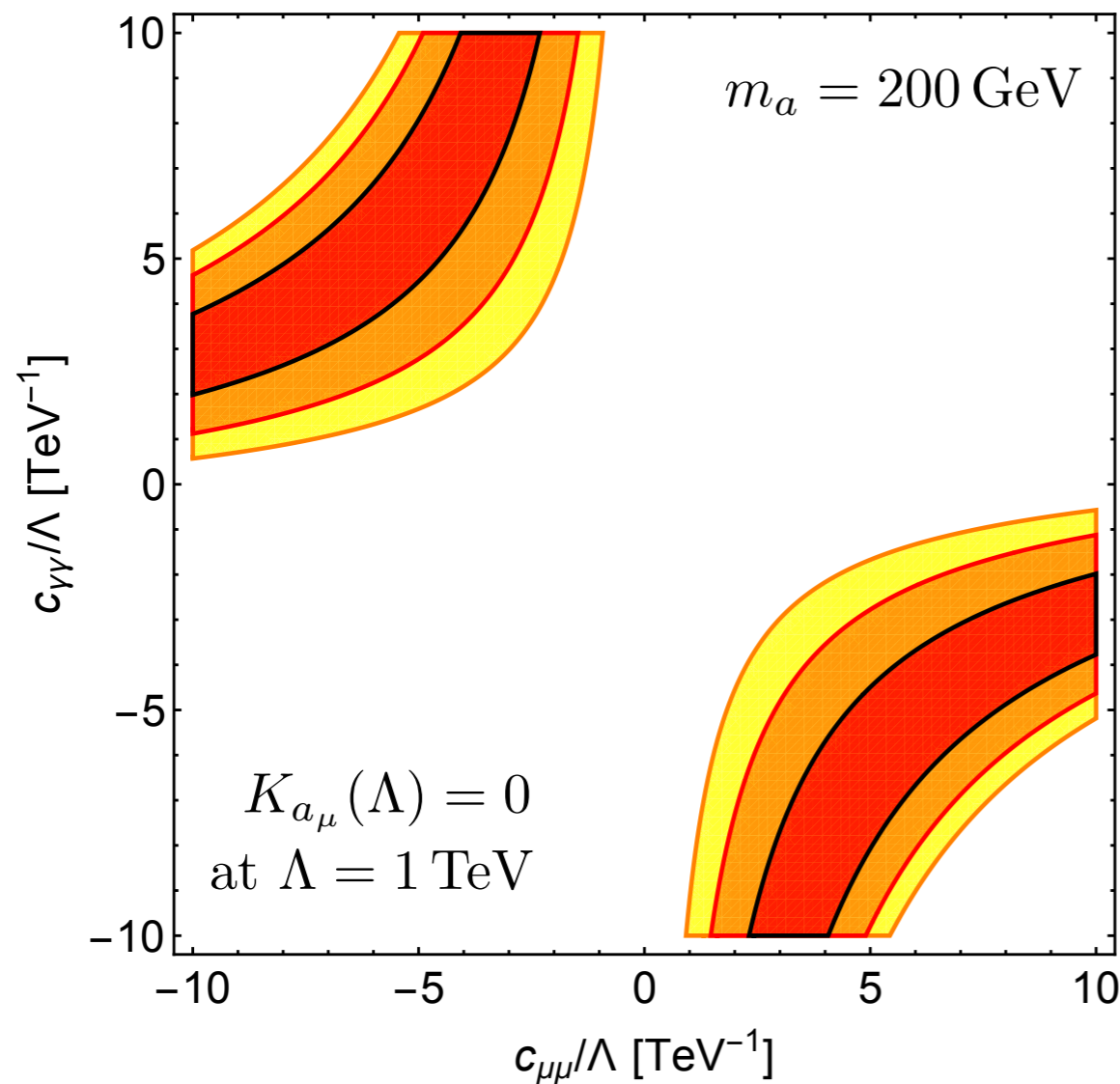
- Fixed target facility at CERN SPS  
(Search for Hidden Particles)



[Alekhin et al.: 1504.04855]

# Muon $(g - 2)_\mu$

- Allowed parameter space moves into corners
- Coupling-mass plots require:  $|C_{\gamma\gamma}|/\Lambda \lesssim 2 \text{ TeV}^{-1}$  and  $|c_{\mu\mu}| \geq |C_{\gamma\gamma}|$



# Probing the parameter space

- Reach in  $Z \rightarrow \gamma a$

