

## Lecture 8

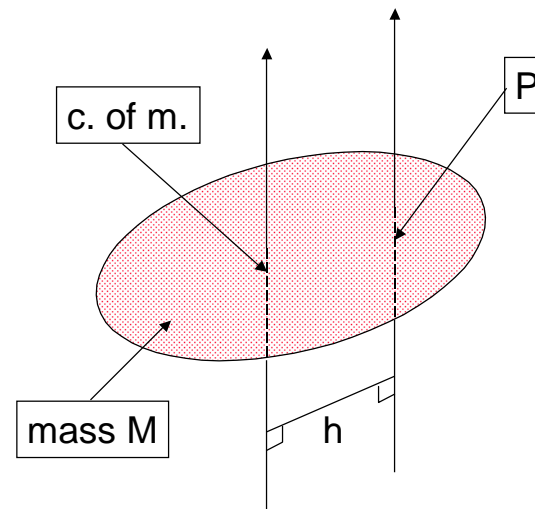
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- ◆ Parallel Axis Theorem
- ◆ Torque
- ◆ Newton's Second Law for Rotation
- ◆ Work, Power and Rotational Kinetic Energy
- ◆ More on Rotational Variables as Vectors
- ◆ More Torque

## Parallel Axis Theorem

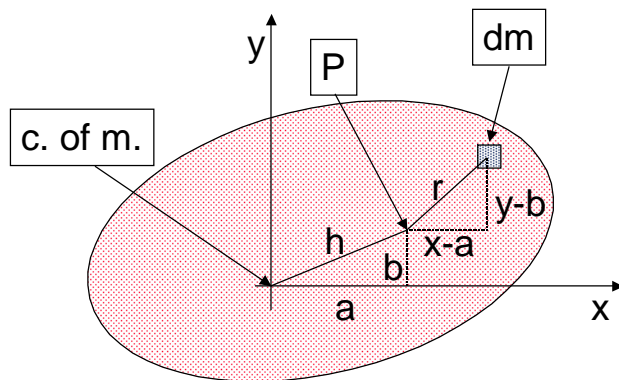
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- ◆ Moment of inertia of body about axis through c. of m. is  $I_{cm}$ . Calc. moment of inertia about parallel axis.



## Parallel axis theorem cont.

- ◆ View from above, place origin at c. of m., z axis along initial axis of rotation.



$$I_P = \int_M r^2 dm$$
$$= \int_M (x - a)^2 + (y - b)^2 dm$$

## Parallel axis theorem cont.

$$I_P = \int_M x^2 + y^2 dm - 2a \int_M x dm -$$
$$2b \int_M y dm + \int_M a^2 + b^2 dm$$

Now 2<sup>nd</sup> and 3<sup>rd</sup> integrals zero from definition of c. of m. and choice of position of origin. First integral just definition of  $I_{cm}$ . We also see:

$$\int_M a^2 + b^2 dm = \int_M h^2 dm = Mh^2$$

Hence parallel axis theorem:

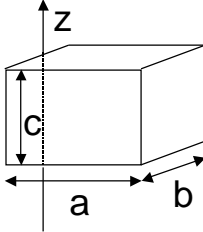
$$I_P = I_{cm} + Mh^2$$

Consequence, minimum moment of inertia for axes through c. of m.

## Parallel Axis Theorem, an Example

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- ◆ Consider again rectangular prism

$$\begin{aligned}
 I_e &= \rho \int_{-\frac{c}{2}}^{\frac{c}{2}} \int_{-\frac{b}{2}}^{\frac{b}{2}} \int_0^a x^2 + y^2 \, dx dy dz \\
 &= c\rho \int_{-\frac{b}{2}}^{\frac{b}{2}} \int_0^a x^2 + y^2 \, dx dy \\
 &= c\rho \int_0^a x^2 y + \frac{y^3}{12} \Big|_{-\frac{b}{2}}^{\frac{b}{2}} dx \\
 &= bc\rho \left( \frac{x^3}{3} + \frac{b^2 x}{12} \right) \Big|_0^a = \frac{abc\rho}{12} (4a^2 + b^2) \\
 &= \frac{M}{12} (4a^2 + b^2)
 \end{aligned}$$


## Parallel Axis Theorem, an Example cont.

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- ◆ Moment of inertia about c. of m.

$$I_{cm} = \frac{M}{12} (a^2 + b^2)$$

- ◆ Using parallel axis theorem

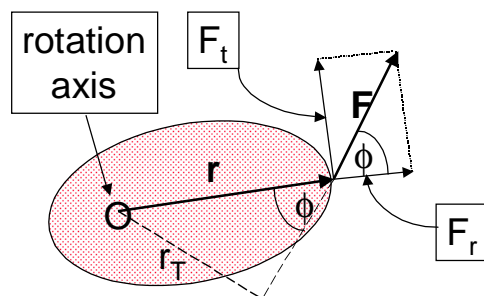
$$\begin{aligned}
 I_e &= I_{cm} + M \left( \frac{a}{2} \right)^2 \\
 &= \frac{M}{12} (a^2 + b^2) + \frac{M}{4} a^2 \\
 &= \frac{M}{12} (4a^2 + b^2)
 \end{aligned}$$

- ◆ Exercise, repeat check for axis along corner of prism and for other shapes.

## Torque

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- ◆ A force applied to a body may tend to rotate the body about an axis. Quantify using concept of torque (from latin “to twist”)
- ◆ Force in plane normal to axis



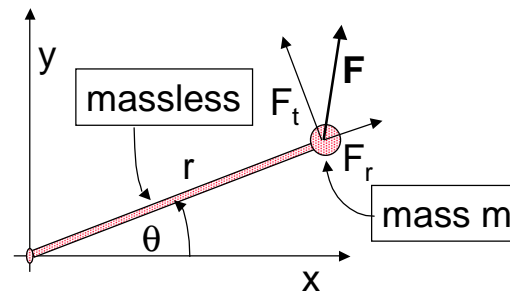
$$\begin{aligned}\tau &= F_t r \\ &= F r_T \\ &= F r \sin \phi \quad \text{units Nm}\end{aligned}$$

Note  $\vec{\tau} = \mathbf{r} \times \mathbf{F}$

## Newton's Second Law

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- ◆ Consider rotation of a (simple) rigid body



- ◆ Relate tangential acceleration to torque

$$\begin{aligned}F_t &= m a_t \\ \tau &= F_t r \\ &= m a_t r\end{aligned}$$

## Newton's second law cont.

- ◆ Consider rotational acceleration

$$\alpha = \frac{d\omega}{dt} = \frac{1}{r} \frac{dv_t}{dt} = \frac{a_t}{r}$$

$$a_t = \alpha r$$

- ◆ Substitute for  $a_t$   
 $\tau = m(\alpha r)r = mr^2\alpha$

- ◆ Recall definition of moment of inertia  
 $I = mr^2$

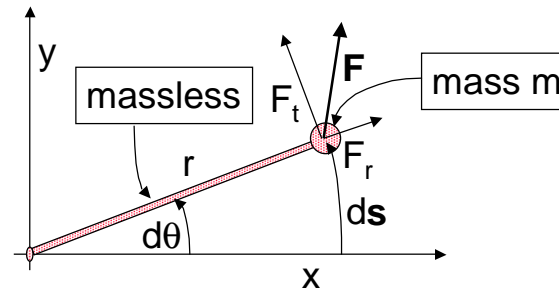
- ◆ Hence Newton's Second Law for rotation  
 $\tau = I\alpha$

- ◆ If many forces applied

$$\sum \tau = I\alpha$$

## Work and Rotation

- ◆ Consider same rigid body rotating through  $d\theta$  under influence of force  $F$



- ◆ Calculate work done

$$dW = \mathbf{F} \cdot d\mathbf{s} = F_t r d\theta = \tau d\theta$$

- ◆ For finite angular displacement

$$W = \int_{\theta_i}^{\theta_f} \tau d\theta \quad \text{cf.} \quad W = \int_{x_i}^{x_f} F dx$$

## Power and Work K.E. Relation

- ◆ Consider power

$$P = \frac{d}{dt} W = \tau \frac{d\theta}{dt} = \tau \omega$$

- ◆ Now relate work to K.E. of rotation

$$W = \int_{\theta_i}^{\theta_f} \tau d\theta = \int_{\theta_i}^{\theta_f} I \alpha d\theta = \int_{\theta_i}^{\theta_f} I \frac{d\omega}{dt} d\theta$$

$$= \int_{\omega_i}^{\omega_f} I \frac{d\theta}{dt} d\omega = \int_{\omega_i}^{\omega_f} I \omega d\omega$$

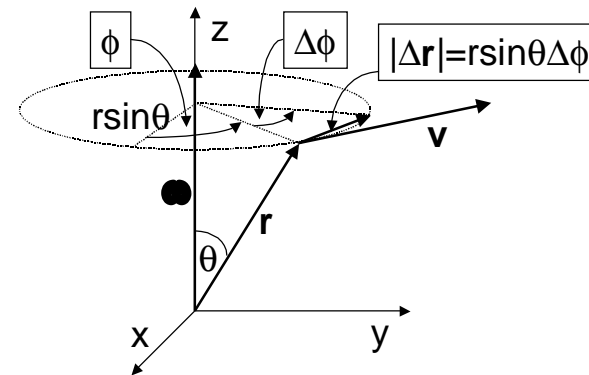
$$= \frac{1}{2} I \omega_f^2 - \frac{1}{2} I \omega_i^2$$

$$= K_f - K_i$$

$$= \Delta K$$

## More on Rotational Variables as Vectors

- ◆ Derive vector angular velocity



$$|\mathbf{v}| = \lim_{\Delta t \rightarrow 0} \frac{|\Delta \mathbf{r}|}{\Delta t} = \lim_{\Delta t \rightarrow 0} \frac{r \sin \theta \Delta \phi}{\Delta t}$$

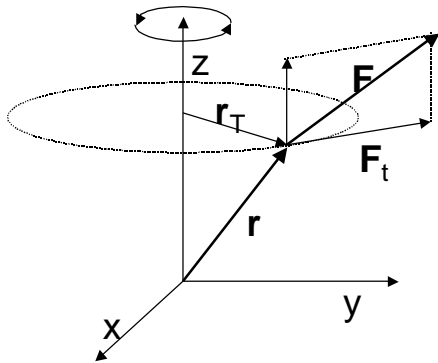
$$= r \sin \theta \frac{d\phi}{dt} = r \omega \sin \theta = |\mathbf{r} \times \boldsymbol{\omega}|$$

As  $\mathbf{v}$  tangential to circle it is normal to plane containing  $\mathbf{r}$  and  $\bullet$ , hence  $\mathbf{v} = \boldsymbol{\omega} \times \mathbf{r}$  (using right hand rule)

## More Torque

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- ◆ For force in plane normal to axis can write torque as  $\boldsymbol{\tau} = \mathbf{r} \times \mathbf{F}$
- ◆ Generalise for all forces.
- ◆ What is torque about given axis?  
Example, torque about z axis.



- ◆ Write  $\mathbf{F} = (f_x, f_y, f_z)$  and  $\mathbf{r} = (r_x, r_y, r_z)$  then  $\mathbf{F}_t = (f_x, f_y, 0)$  and  $\mathbf{r}_T = (r_x, r_y, 0)$ .

## More torque cont.

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- ◆ From previous discussion we know torque about z is  
 $\tau_z = |\mathbf{r}_T \times \mathbf{F}_t|$
- ◆ In terms of vector components this is  
 $\tau_z = |(r_x, r_y, 0) \times (f_x, f_y, 0)|$

$$\begin{aligned} &= \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ r_x & r_y & 0 \\ f_x & f_y & 0 \end{vmatrix} \\ &= r_x f_y - r_y f_x \end{aligned}$$

## More torque cont.

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- ◆ Consider torque as vector and take component of vector along z axis to get torque about z

$$\vec{\tau} = \mathbf{r} \times \mathbf{F}$$

$$= \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ r_x & r_y & r_z \\ f_x & f_y & f_z \end{vmatrix}$$

$$= (r_y f_z - r_z f_y, r_z f_x - r_x f_z, r_x f_y - r_y f_x)$$

$$\tau_z = \vec{\tau} \cdot \hat{\mathbf{k}} = \begin{pmatrix} r_y f_z - r_z f_y \\ r_z f_x - r_x f_z \\ r_x f_y - r_y f_x \end{pmatrix} \cdot \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$= r_x f_y - r_y f_x$$

- ◆ The correct result.