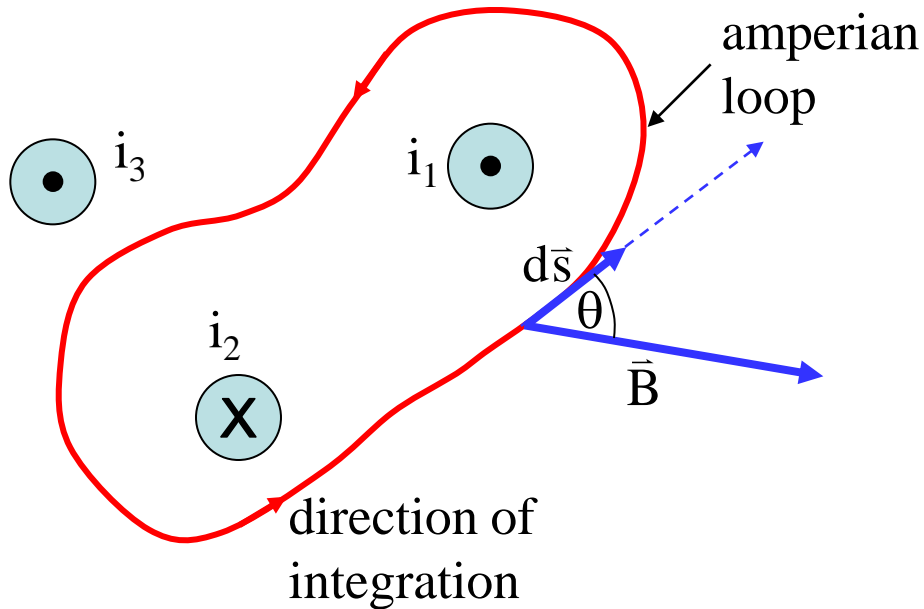


Lecture 15

- In this lecture we will look at:
 - ◆ Ampere's Law.
 - ◆ Magnetic fields associated with straight wires.
 - ◆ B fields in solenoids and toroids.
 - ◆ The magnetic dipole due to a current loop.
- After this lecture, you should be able to answer the following questions:
 - State Ampere's Law in words and as a mathematical formula.
 - Using Ampere's Law, determine the strength of the magnetic field in the centre of a long, current carrying solenoid.
 - Explain how the direction of the magnetic field due to a current loop can be determined.

Ampere's Law

- Despite the name, this law was first deduced by Maxwell!



- Ampere's law:

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 i_{\text{enc}} \quad [15.1]$$

- Can be written:

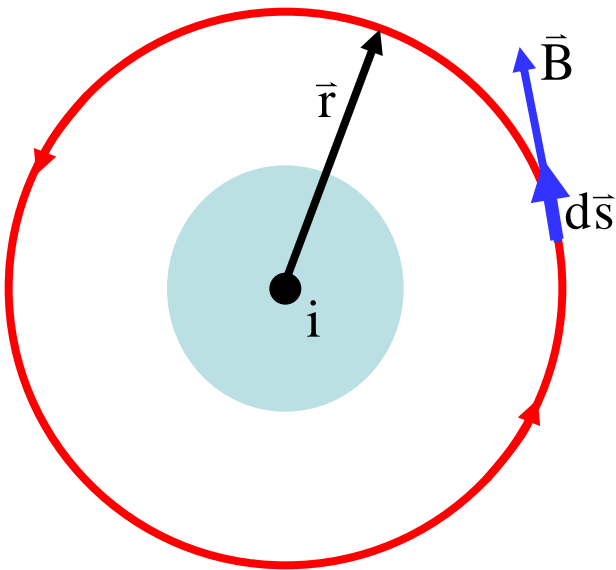
$$\oint B \cos \theta ds = \mu_0 i_{\text{enc}}.$$

- Determine sign of current using yet another right hand rule.
- Curl right hand round amperian loop so fingers point along direction of integration, currents in direction of thumb assigned +ive sign, in opposite direction -ive sign.
- Now have

$$\oint B \cos \theta ds = \mu_0 (i_1 - i_2).$$

Magnetic Field Outside/Inside Long Straight Wire

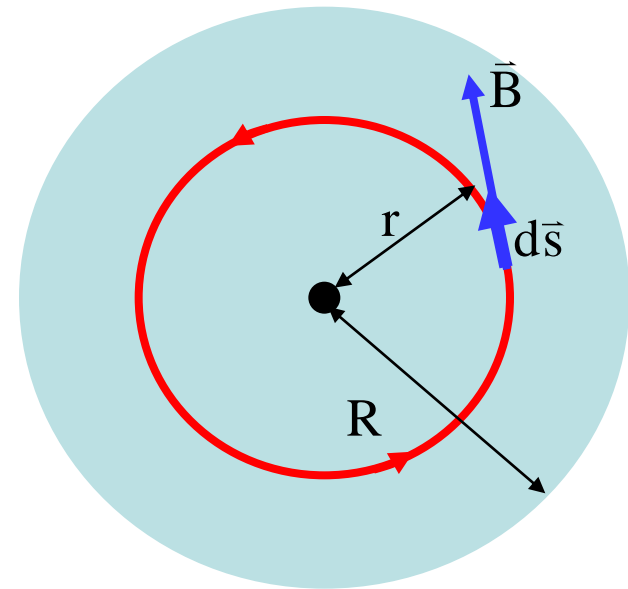
- Outside wire:



- $\oint \vec{B} \cdot d\vec{s} = \oint B \cos \theta ds = B \oint ds = B(2\pi r).$

- Hence: $2\pi r B = \mu_0 i$
 $\Rightarrow B = \frac{\mu_0 i}{2\pi r}$ [15.2]

- Inside wire:

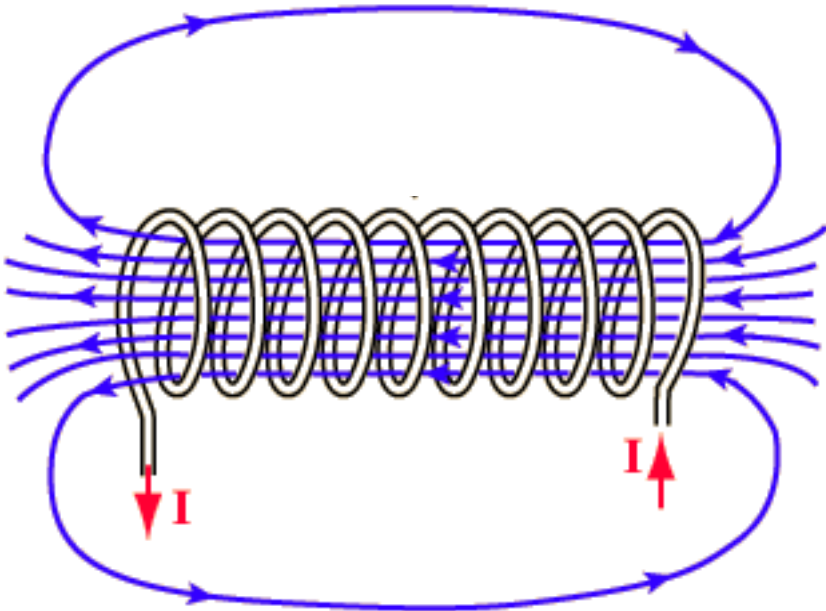


- $2\pi r B = \mu_0 i_{\text{enc}} = i \mu_0 \frac{\pi r^2}{\pi R^2}.$

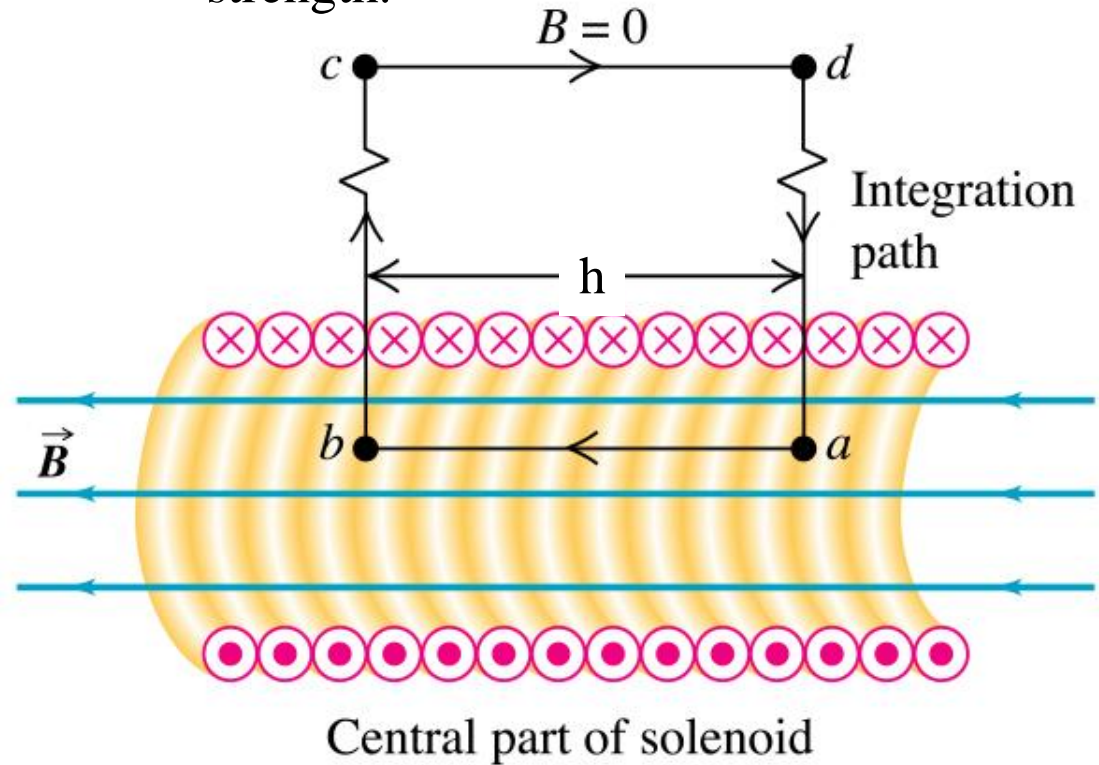
- Hence $B = \frac{\mu_0 i}{2\pi R^2} r$ [15.3]

Field in a Solenoid

- Magnetic field is nearly uniform in the centre of a long solenoid.
- Field outside is relatively weak.



- Use Ampere's law to determine field strength.



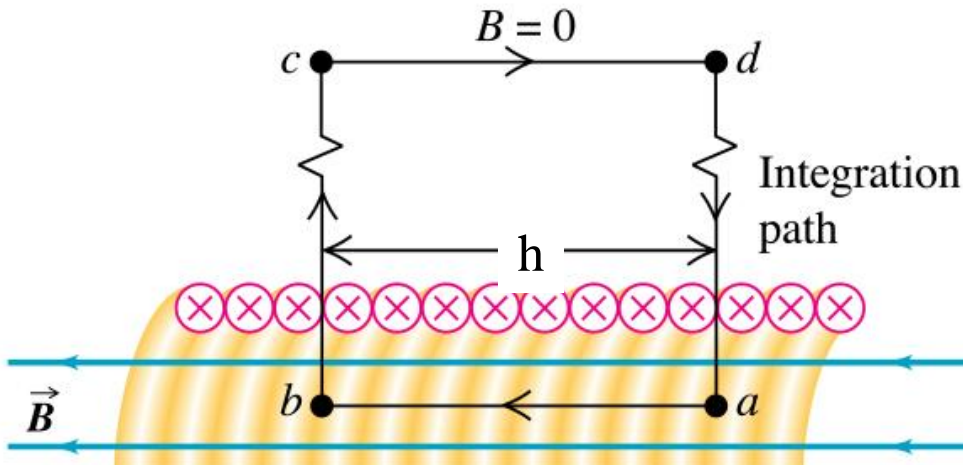
- Using this amperian loop:

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 i_{\text{enc}}.$$

Field in Solenoid

- But the integral can be split up into four parts:

$$\oint \vec{B} \cdot d\vec{s} = \int_a^b \vec{B} \cdot d\vec{s} + \int_b^c \vec{B} \cdot d\vec{s} + \int_c^d \vec{B} \cdot d\vec{s} + \int_d^a \vec{B} \cdot d\vec{s}.$$



- No of turns per unit length n .

- $\int_a^b \vec{B} \cdot d\vec{s} = Bh.$

- $\int_b^c \vec{B} \cdot d\vec{s} = \int_c^d \vec{B} \cdot d\vec{s} = 0.$

- Because the path c to d is a long way from the solenoid, the field is negligible, hence:

- $\int_c^d \vec{B} \cdot d\vec{s} = 0.$

- The result is thus: $\oint \vec{B} \cdot d\vec{s} = Bh.$

- The enclosed current is: $i_{\text{enc}} = i(nh).$

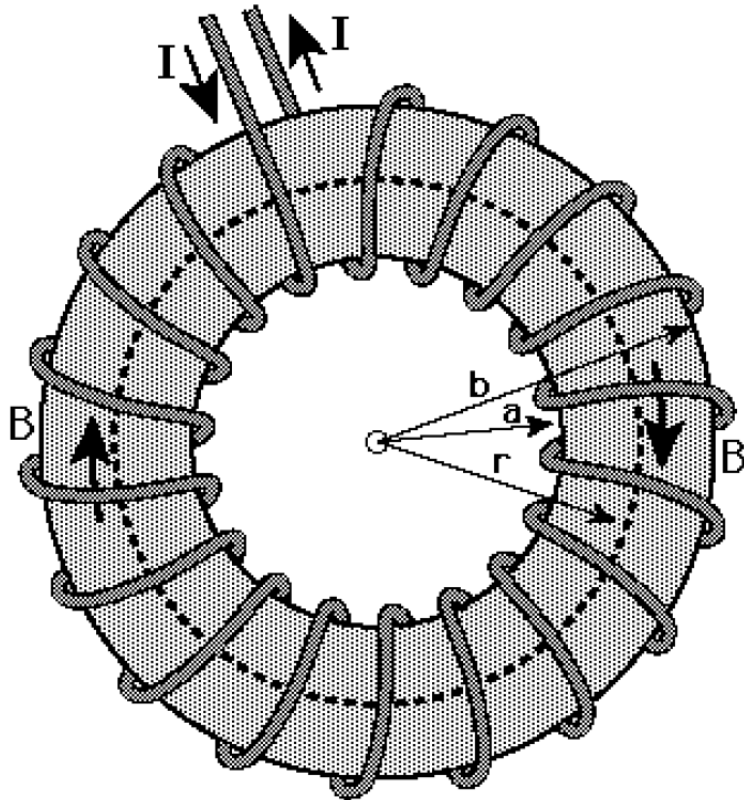
- Hence:

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 i_{\text{enc}} \Rightarrow Bh = \mu_0 inh$$

$$\text{or } B = \mu_0 in \quad [15.4]$$

Field in Toroid

- Toroid is solenoid bent into ring:



- Total number of turns N .

- From Ampere's law (amperian loop at radius r traversed in clockwise direction) we get:

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 i_{\text{enc}}$$

$$\Rightarrow B(2\pi r) = \mu_0 iN$$

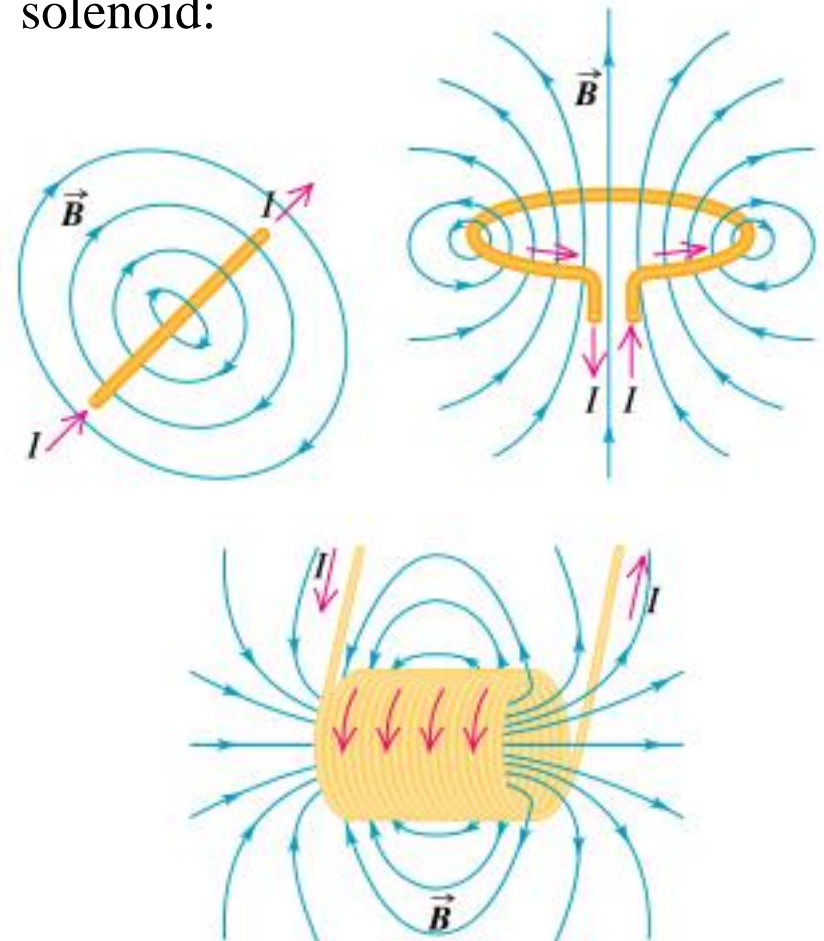
$$\text{or } B = \frac{\mu_0 iN}{2\pi r} \quad [15.5]$$

- Note that field is not uniform as in solenoid, but decreases with increasing radius.
- Field outside toroid (i.e. $r < a$ or $r > b$ or above or below toroid) is zero.

Current Loop as Magnetic Dipole

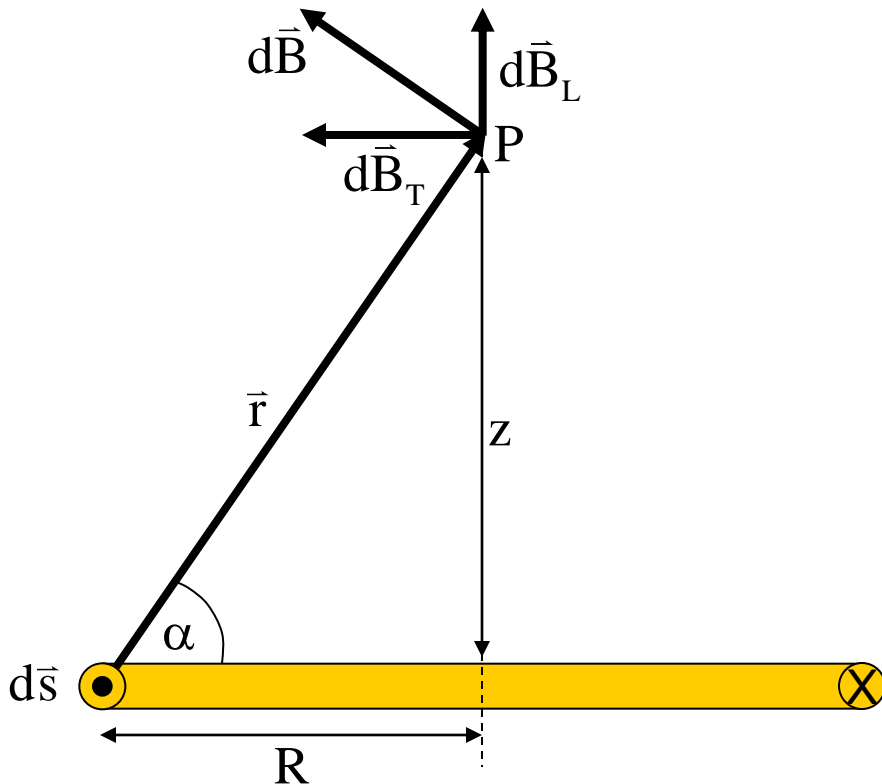
- Have seen torque on current carrying loop in \vec{B} field can be expressed in terms of a magnetic dipole:
$$\vec{\tau} = \vec{\mu} \times \vec{B}.$$
- As “loop” configuration is of importance (electrons in atoms!), now calculate field due to loop.
- Use Biot-Savart law as opposed to Ampere’s law due to limited symmetry.
- (Have already calculated the field at the centre of the loop: $B = \mu_0 i / 2R$.)

- Fields due to wire, current loop and solenoid:



Current Loop as Magnetic Dipole

- Diagram shows rear half of current loop, radius R , loop perpendicular to transparency:



- Angle θ between $d\vec{s}$ and \vec{r} is 90° .

- Biot-Savart law:

$$d\vec{B} = \frac{\mu_0}{4\pi} \frac{i d\vec{s} \sin \theta}{r^2} = \frac{\mu_0}{4\pi} \frac{i d\vec{s}}{r^2}.$$

- Transverse field components cancel.

$$dB_L = \frac{\mu_0 i \cos \alpha ds}{4\pi r^2}.$$

- Use $r^2 = R^2 + z^2$ and $\cos \alpha = \frac{R}{\sqrt{R^2 + z^2}}$.

- Hence:
$$dB_L = \frac{\mu_0 i R}{4\pi (R^2 + z^2)^{\frac{3}{2}}} ds$$

$$\begin{aligned} \text{and } B &= \int dB_L = \frac{\mu_0 i R}{4\pi (R^2 + z^2)^{\frac{3}{2}}} \int ds \\ &= \frac{\mu_0 i R \times 2\pi R}{4\pi (R^2 + z^2)^{\frac{3}{2}}} = \frac{\mu_0 i R^2}{2(R^2 + z^2)^{\frac{3}{2}}}. \end{aligned}$$

Current Loop as Magnetic Dipole

- For points along the z axis far from the loop we have:

$$\mathbf{B} \approx \frac{\mu_0 i R^2}{2z^3}.$$

- Using $A = \pi R^2$: $\mathbf{B} \approx \frac{\mu_0 i A}{2\pi z^3}$.
- If we allow the coil to have N turns this becomes:

$$\mathbf{B} \approx \frac{\mu_0 N i A}{2\pi z^3}.$$

- Writing $\mu = NiA$, we have:

$$\bar{\mathbf{B}} \approx \frac{\mu_0}{2\pi} \frac{\bar{\mu}}{z^3} \quad [15.6]$$

- Curl RH fingers round loop in direction of i , thumb gives direction of B field in loop and of dipole.
- If set $z = 0$ we get agreement with previous calculation using Ampere's Law:

$$\mathbf{B} = \frac{\mu_0 i}{2R}.$$