

The Standard Model

1: Overview / introduction

Course Outline

Lectures:

1. Overview of SM framework (2 lectures)
2. Tests of SM with Z^0 bosons
3. Test of SM with W^\pm bosons
4. Quark mixing and CP violation (2 lectures)

Outline

Overview of Standard Model:

- What does it describe?
- What does it need as inputs?
 - What do we use and why
- What can it predict?
- Where can it be tested?

Overview

Standard model describes interactions of fundamental particles:

- **Fermions**; 6 flavours of quarks (3 types - **rgb**), leptons
- **Vector bosons**; γ , $W^{+/-}$, Z^0 , 8 gluons
- **Scalar bosons**; H

Theory describes at least 61 fundamental particles and three forces!

Overview

- SM unites electromagnetic, strong and weak forces
- Represented theoretically by
 - U(1) hypercharge
 - SU(3) colour
 - SU(2) isospin
- Use lagrangian to describe particle field and interactions

Lagrangians

- Describe interactions and fields
- Classically, $L = \text{kinetic energy} - \text{potential energy}$
- Particle physics:
 - Use **Dirac equation** to describe **free** spin-1/2 particle:

$$L = \bar{\Psi}(i\gamma^\mu\partial_\mu - m)\Psi$$

Ψ = wavefunction

m = mass

$\gamma^\mu = \mu^{\text{th}}$ gamma matrix

∂_μ = partial derivative

Symmetries

Note: symmetries in physics imply conservation laws

symmetry	invariant
Space translation	momentum
Time translation	energy
Rotation	Angular momentum
Global phase; $\Psi \rightarrow e^{i\theta}\Psi$	Electric charge
Local phase; $\Psi \rightarrow e^{i\theta(x,t)}\Psi$	Lagrangian + gauge field (\rightarrow QED)

Lagrangians

Apply local gauge symmetry to Dirac equation:

$$\bar{\Psi} \rightarrow e^{i\theta(x,t)}\bar{\Psi}, \Psi \rightarrow e^{-i\theta(x,t)}\Psi$$

Consider very small changes in field:

$$\Psi \rightarrow \Psi + \delta\Psi = \Psi - i\theta(x,t)\Psi \quad \text{ie. } \delta\Psi = -i\theta(x,t)\Psi$$

$$L = \bar{\Psi}(i\gamma^\mu \partial_\mu - m)\Psi \Rightarrow \delta L = \bar{\Psi}\gamma^\mu \partial_\mu \theta(x,t)\Psi$$

If lagrangian is invariant then $\delta L=0$

Lagrangians

Invariant lagrangian $\Rightarrow \delta L=0$

Satisfied:

- 1) if we introduce gauge field A_μ to interact with fermion, and A_μ transforms as;

$$A_\mu + \delta A_\mu = A_\mu + 1/e \partial_\mu \theta(x,t)$$

- 2) If we replace $\partial_\mu \rightarrow D_\mu = \partial_\mu + ieA_\mu$

Hence $L = \bar{\Psi}(i\gamma^\mu D_\mu - m)\Psi$ is invariant

You can check this ...

Lagrangians

Not the whole story – need to add term for field strength (kinetic term):

Define $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$

Add term $-1/4 F_{\mu\nu} F^{\mu\nu}$ (Lorentz invariant, matches Maxwell's equations)

Final lagrangian (for QED!):

$$L = -1/4 F_{\mu\nu} F^{\mu\nu} + \bar{\Psi}(i\gamma^\mu D_\mu - m)\Psi$$

Nb. No mass term for A_μ allowed; then L is not invariant

QED from gauge invariance requirement + massless γ !

QED

Electromagnetic interactions; abelian

- U(1): 1 gauge field \mathbf{B}_μ , coupling g'
 - In SM, source is hypercharge Y
 - Field strength $\alpha_{em} = g'^2/4\pi$ (fine structure constant)
 $\alpha_{em} = 1/137$

$$\mathcal{L}_Y = -1/4 F_{\mu\nu} F^{\mu\nu} + \bar{\Psi}(i\gamma^\mu D_\mu - m)\Psi$$

where $\mathbf{D}_\mu = \partial_\mu + ie\mathbf{B}_\mu$
to conserve invariance under $\Psi \rightarrow e^{-iY(\Psi)}\Psi$

Strong, weak lagrangians

Strong, weak forces are described by **non-abelian theories**:

- In non-abelian theories gauge bosons can self interact
- In non-abelian theories gauge invariance achieved by adding n^2-1 gauge bosons for SU(n)
 - SU(2) – 3 gauge bosons (W^1, W^2, W^3) for weak force
 - SU(3) – 8 gluons for QCD

QCD

Strong interactions; non-abelian

- SU(3): 8 massless gauge bosons (gluons), coupling g_s
 - Source is colour
 - Field strength $\alpha_s = g_s^2/4\pi$
- Fields represented by quark triplets (3 colours)
 - Hadrons (observable states) colourless
 - Leptons, neutrinos do not couple to gluons

QCD

$$L_{\text{QCD}} = -1/4 F_{\mu\nu}^{(a)} F^{(a)\mu\nu} + \bar{\Psi}_j (i\gamma^\mu (D_\mu)_{jk} - m_{jk}) \Psi_k$$

where:

$$D_\mu = \delta_{jk} \partial_\mu + ig_s \Sigma_a (\lambda^a_{j,k}/2) G^a_\mu$$

λ^a are 8 3x3 matrices (analogous to Pauli spin matrices in SU(2))

a is sum over 8 gauge bosons (section 35, PDG)

j,k is colour index (=1,2,3)

G^a_μ is (a^{th}) gluon gauge field

$F_{\mu\nu} F^{\mu\nu}$ contains triple, quartic couplings

$$F_{\mu\nu}^{(a)} = \partial_\mu G^a_\nu - \partial_\nu G^a_\mu - g_s f_{abc} G^b_\mu G^c_\nu$$

f_{abc} structure constants for SU(3) (section 35, PDG)

Weak force

Weak interactions: non-abelian

- SU(2): 3 gauge bosons W^1, W^2, W^3
 - coupling strength g
 - Source is weak charge
- **Observed** to violate parity so left, right handed fermions treated separately in theory

$$\Psi = \Psi_L + \Psi_R$$

Ψ_L **doublet**; interacts under weak force

Ψ_R **singlet**, does not interact under weak force

nb: $\Psi_L = \frac{1}{2} (1 \mp \gamma^5) \Psi$ in theory

Weak force

$$\mathcal{L}_{\text{WEAK}} = -\frac{1}{4} F_{\mu\nu}^{(a)} F^{(a)\mu\nu} + i\bar{\Psi}_L (\gamma^\mu (D_\mu) - m) \Psi_L$$

where:

$$D_\mu = \partial_\mu + igT^a W_\mu^a$$

a is sum over 3 gauge bosons

$T^a = 0.5 \times$ Pauli spin matrices (PDG!)

W_μ^a is weak gauge boson field

$$F_{\mu\nu}^{(a)} = \partial_\mu W_\nu^a - \partial_\nu W_\mu^a - g_s \epsilon_{abc} W_\mu^b W_\nu^c$$

$\epsilon_{abc} =$ SU(2) structure constants

$F_{\mu\nu} F^{\mu\nu}$ contains triple, quartic couplings

Combine forces in SM

Left handed fermions:

- Interact with weak force
- Interact with electromagnetic force
- Quarks interact with strong force

Right handed fermions:

- Interact with electromagnetic force
- Quarks interact with strong force

nb: strong quark eigenstates are superposition of weak quark eigenstates: \rightarrow CKM matrix (later)

Full SM lagrangian

SM: $U(1) \times SU(2) \times SU(3)$

Substitute in for D_μ , $F_{\mu\nu}$ for each interaction, and fermion field $\Psi = \Psi_L + \Psi_R$

$$\mathcal{L}_{(SM,1)} = \sum_{\text{gauge bosons}} -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \sum_{\text{fermions}} \bar{\psi} \not{D}\psi + (D_\mu \Phi)^\dagger (D^\mu \Phi) - V(\Phi)$$

Free particle term

Gauge boson interaction terms

Higgs field Φ terms (to give mass)

$$V(\Phi) = -\mu^2 \Phi^\dagger \Phi + \lambda |\Phi^\dagger \Phi|^2$$

Note; Masses

- Everything in SM lagrangian is **massless**
 - Mass conferred via **Higgs mechanism**
 1. Introduce Higgs field (doublet) $\Phi = (0, v+H)$
 2. Substitute into lagrangian
 3. $D_\mu \Phi D^\mu \Phi$ contains terms of the form

$$(g^2 v^2 / 4) W_\mu^+ W_\mu^- + (v^2 / 8) (g W_\mu^3 - g' B_\mu)^2$$
 where $W_\mu^+ = (W_\mu^1 + i W_\mu^2) / \sqrt{2}$ etc.
 v = vacuum expectation value (246 GeV)
- ⇒ Masses for gauge bosons
 Mixing of U(1) and SU(2) gauge bosons (electroweak unification)

→ Ronan

Electroweak unification

W_μ^1, W_μ^2 mix after Higgs breaking → massive W^+ and W^-

$$W^+ = (W_\mu^1 + i W_\mu^2) / \sqrt{2} \quad (\text{massive})$$

$$W^- = (W_\mu^1 - i W_\mu^2) / \sqrt{2} \quad (\text{massive})$$

W_μ^3, B_μ mix after Higgs breaking → massive Z^0, γ

$$Z^0 = \cos \theta_W W_\mu^3 - \sin \theta_W B_\mu \quad (\text{massive})$$

Unification here!

$$\gamma = \sin \theta_W W_\mu^3 + \cos \theta_W B_\mu \quad (\text{massless})$$

SM relates M_W, M_Z and g', g :

$$\tan \theta_W = g' / g$$

$$M_W = M_Z \cos \theta_W$$

Unification here!

Effect on W^\pm, γ couplings

Subst for W^+, W^-, Z^0, γ in lagrangian;

- Couplings for W^\pm of form (only l.h. fermions)

$$-g/2\sqrt{2} \bar{\nu} \gamma_\mu (1-\gamma^5) e W_\mu^-$$

- Couplings for **photons** of form (l.h. + r.h. charged fermions)

$$g \sin \theta_W \bar{e} \gamma_\mu e B_\mu$$

γ_μ transforms as vector (P odd)

$\gamma_\mu \gamma^5$ transforms as axial vector (P even)

(Vector + axial vector \Rightarrow parity violation for weak force)

You can check this ...

Effect on couplings

Subst for W^+, W^-, Z^0, γ in lagrangian;

- Couplings for Z^0 of form (l.h. (W_μ^3) and r.h. (B_μ) fermions)

$$-g/4 \cos \theta_W \bar{\nu} \gamma_\mu (1-\gamma^5) \nu Z_\mu \quad (\text{l.h. neutrinos})$$

$$g/4 \cos \theta_W \bar{e} (\gamma_\mu (1-\gamma^5) - 4 \sin^2 \theta_W \gamma_\mu) e Z_\mu \quad (\text{l.h. \& r.h. e})$$

\Rightarrow Effective vector and axial couplings g_{vf}, g_{af} for Z^0 decays; $g/4 \cos \theta_W e (\gamma_\mu (g_{vf} - g_{af} \gamma^5)) e Z_\mu$

fermion	g_{vf}	g_{af}
ν_e, ν_μ, ν_τ	0.5	0.5
e, μ, τ	$-0.5 + 2 \sin^2 \theta_W$	-0.5
u, c, t	$0.5 - 4/3 \sin^2 \theta_W$	0.5
d, s, b	$-0.5 + 2/3 \sin^2 \theta_W$	-0.5

Feynman rules

Propagators (free particle L)

	$-i (g_{\mu\nu} - p_\mu p_\nu / M_W^2) / (p^2 - M_W^2)$
	$-i (g_{\mu\nu} - p_\mu p_\nu / M_Z^2) / (p^2 - M_Z^2)$
	$-i g_{\mu\nu} / p^2$
	$i (\gamma \cdot p + m_e) / (p^2 - m_e^2)$
	$i \gamma \cdot p / p^2$
	$i / (p^2 - m_H^2)$

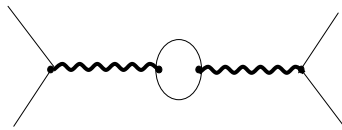
vertex factors (interaction terms of L)

	$-i (g/2\sqrt{2}) \gamma_\mu (1 - \gamma^5)$
	$+i g \sin \theta_W \gamma_\mu$
	$+ \frac{1}{2} i (g / \cos \theta_W) \gamma_\mu (1 - 4 \sin^2 \theta_W - \gamma^5)$
	$- \frac{1}{2} i (g / \cos \theta_W) \gamma_\mu (1 - \gamma^5)$

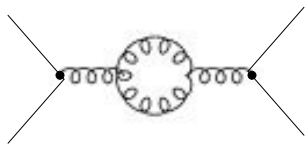
Allow us to calculate cross-sections

Running coupling constants

- Arise from loop diagrams “screening” charge at large distance scales:
 \Rightarrow Coupling constant values are function of E



Screening of charge by vacuum polarisation;
 High E \Rightarrow smaller distance scale
 \Rightarrow see more charge
 Coupling constant increases with E



Non-abelian forces also include these “extra” charge loops
 Net effect: coupling constant decreases with E (asymptotic freedom)

How many free parameters?

SM does not predict:

- Magnitude of gauge couplings g , g' , g_s
- Masses of fermions (3 leptons, 6 quarks, 3 neutrinos)
- Weak-strong eigenstate quark mixing (can express by 4 parameters), ditto for neutrinos
- Higgs related quantities (mass of Higgs and vacuum expectation value)

⇒ **Some 26 unknowns in the theory.**

- These values must be added by hand (experimental measurements)

SM predicts relationships

All observables can be predicted in terms of 26 free parameters

- If we have > 26 measurements of these observables, we overconstrain SM
- Overconstrain \Rightarrow we don't have any more ad hoc inputs AND we can test the consistency of the model
- **Best plan:**
 - pick well measured set of observables
 - Calculate other observables in terms of these well known quantities
 - Test predictions; measure observable, compare to theory

SM predicts relationships

Due to predicted relationships, can reduce number of unknowns

- For example, electroweak boson masses and couplings $M_w, M_z, M_\gamma, g, g'$:

$$g \cdot \sin \theta_w = g' \cdot \cos \theta_w = e$$

$$M_\gamma = 0$$

$$M_w = M_z \cdot \cos \theta_w$$

→ 3 parameters free ($g, \sin \theta_w, M_w$)

- In practise, pick 3 that are best measured experimentally:

$$\alpha_{em}, G_F, M_z$$

Where can the SM be tested?

Particle physics experiments designed to test specific aspects of SM

- Major historical experiments:

- LEP (ALEPH, DELPHI, L3, OPAL) **Electroweak, qcd**

$$(\sqrt{s} = M_z \rightarrow 2M_w)$$

- Major running experiments:

- Babar, Belle ($\sqrt{s} = 2M_B$) **Quark mixing, CP**

- CDF, D0 ($\sqrt{s} = 2 \text{ TeV}$) **electroweak, qcd, quark mixing**

- H1, Zeus ($\sqrt{s} = 0.3 \text{ TeV}$) **qcd**

Review

SM unites electromagnetic, weak, strong forces

SM predicts cross-sections, couplings

SM incomplete – 26 free parameters

- Relations between some free parameters are predicted

Next lectures will discuss how these free parameters are measured and SM is tested