Lecture 7 – Detectors - I

Early Techniques Collider Detectors

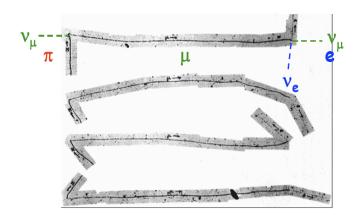
Solenoidal Field

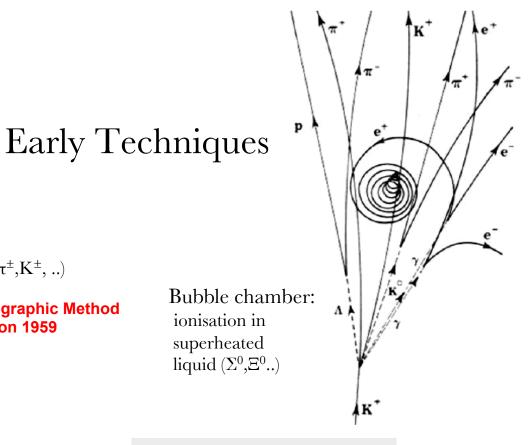
Tracking

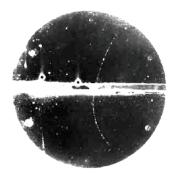
Drift Chamber

Silicon Detector

Introduction to Particle Physics - Max Klein – Lecture 7 - Liverpool University 17.3.14







Emulsion: silver halide crystals in gelatine $(\pi^{\pm}, K^{\pm}, ..)$

C.Powell, P.Fowler, D.Perkins, Pergamon, London 1959

The Study of Elementary Particles by the Photographic Method

Cloud chamber: ionisation in cloud (e^+ , μ^{\pm} , K^0 ,...)

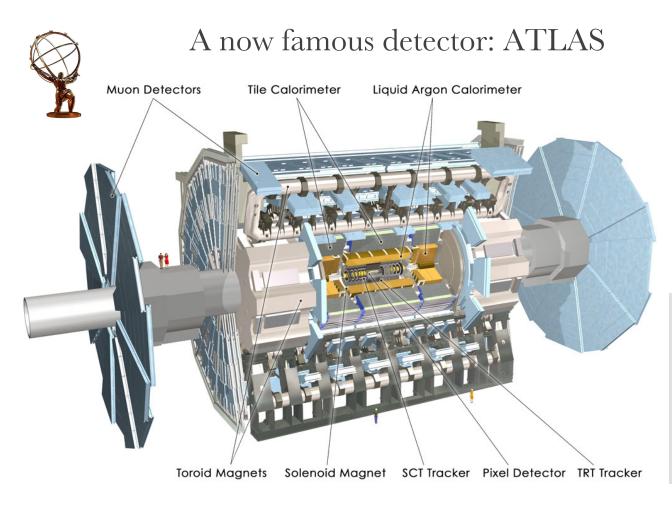
photographic techniques not triggerable small active volumes \rightarrow small event rates

ionisation in

superheated

M, p, Q from ionisation and curvature in magnetic field

A.Bettini, Introduction to Elementary Particle Physics, Cambridge, 2008

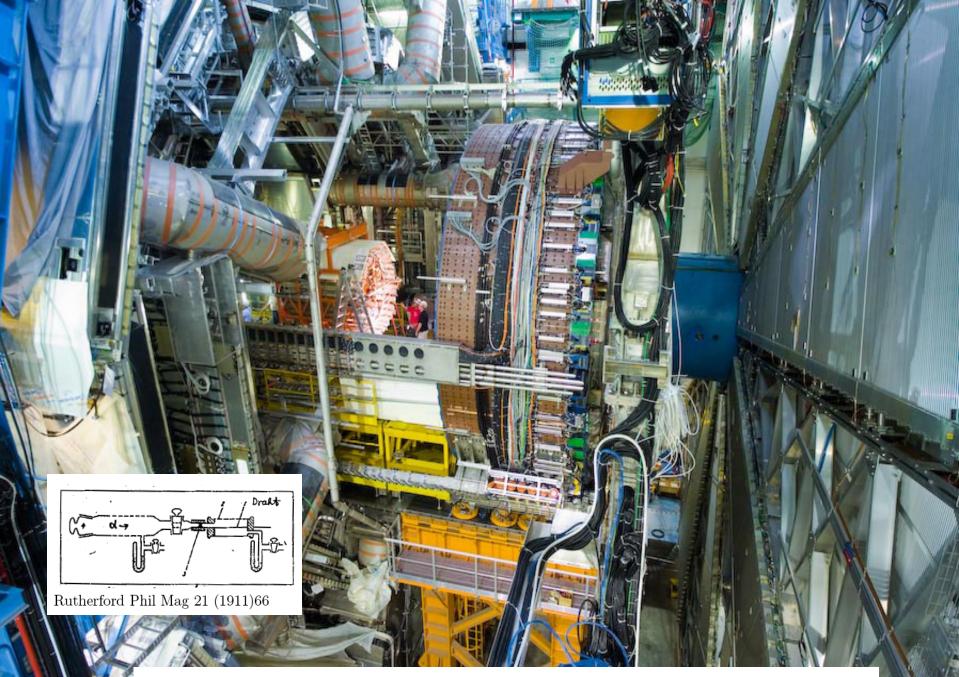


Design 92-97 R&D 94-02 Construction 99-08 Operation 08-33(?)

High energies are nowadays achieved with colliders. Full (4π) acceptance is crucial for searches for new particles. This determines the by now classic detector structure, at LEP, HERA, Tevatron, now the LHC as for future colliders.



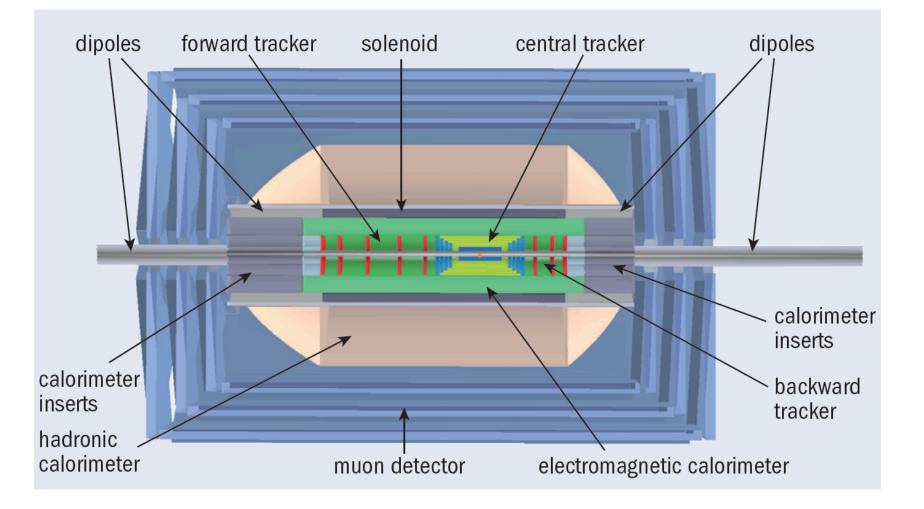
ATLAS: global collaboration of ~ 200 institutes and ~3300 physicists and engineers to explore the physics at the accelerator energy frontier. Detector: housed 100m underground in cavern 25m x 45m D x Z, 7000t, 10⁸ channels, 3000 cables



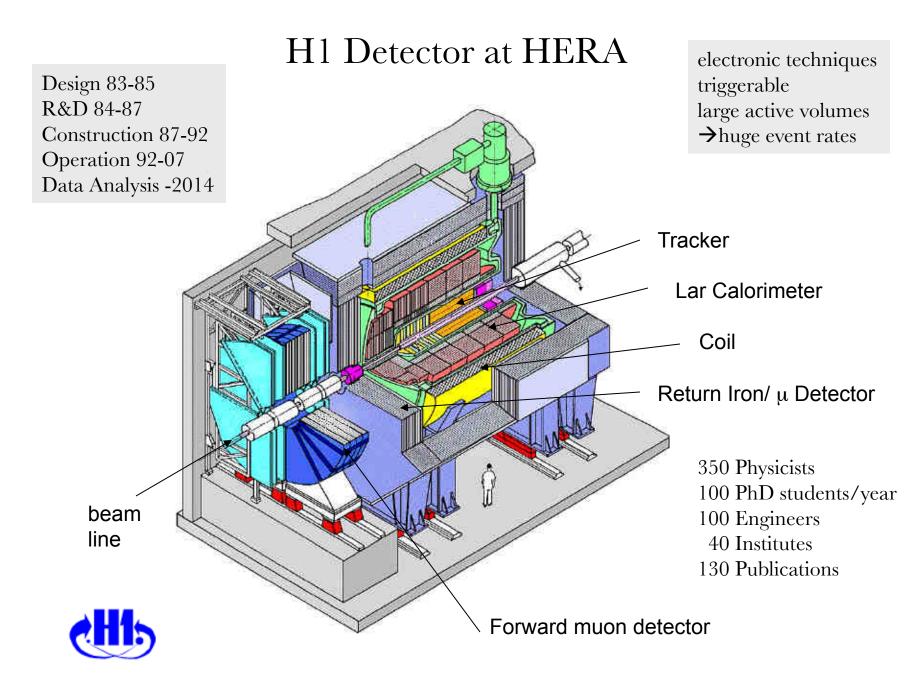
A gigantic evolution of instrumentation, electronics and computing over 100 years

A detector under design

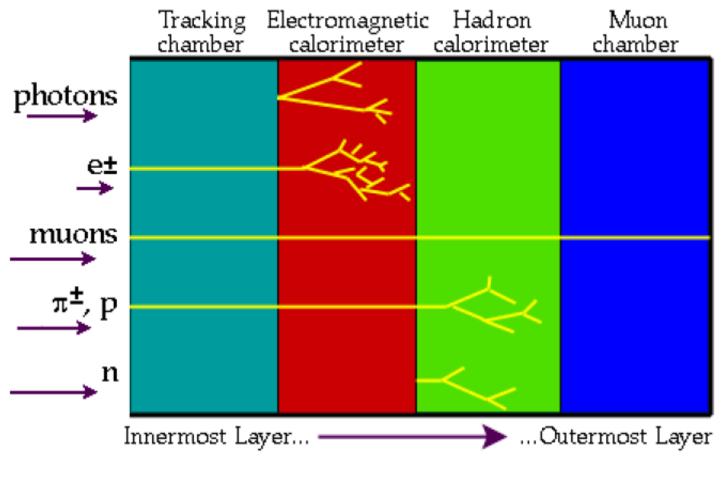




Muon detector: ID with few detector planes, no momentum measurement – unlike ATLAS Detector dimensions: 9m diameter,12m length M.Klein 17.3.2014 L7

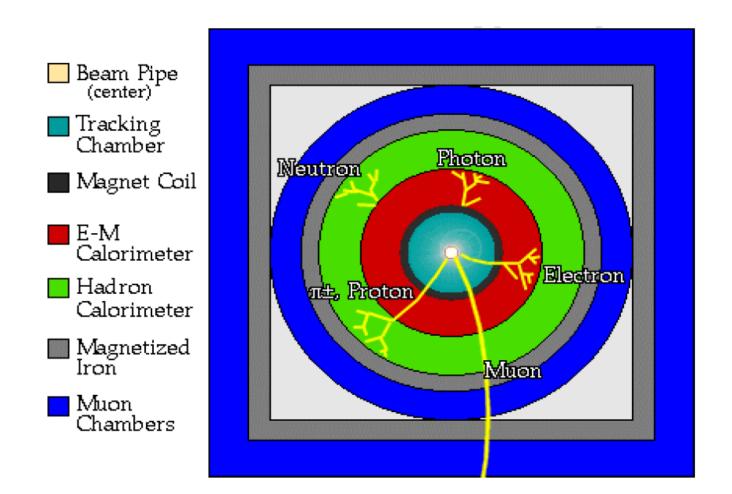


Particle Identification

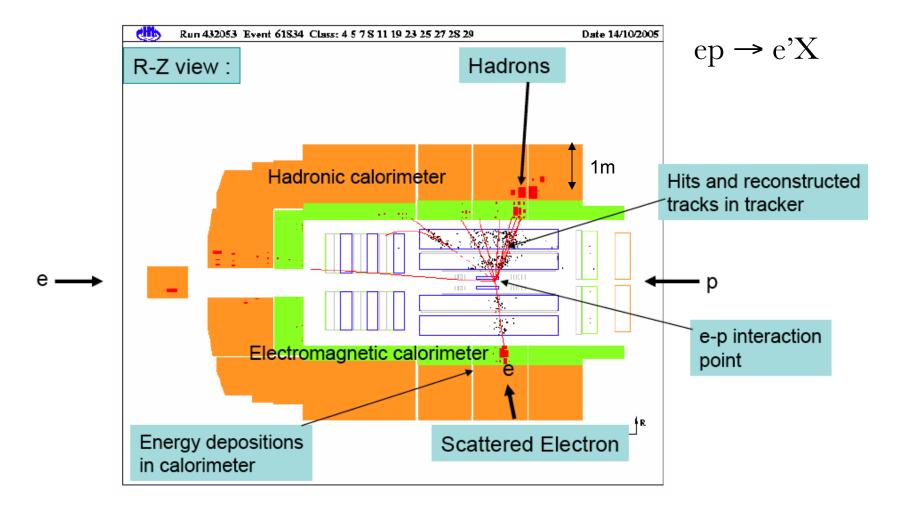


→ v neutrino carries energy away ("missing E") -- Hermetic Detector!

Particle Identification

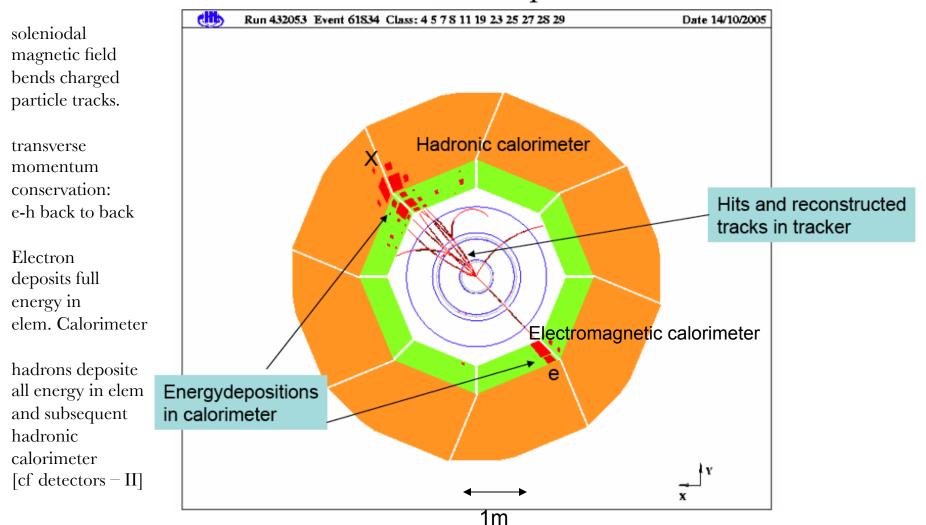


Deep inelastic scattering event in H1



The same event in xy view

 $ep \rightarrow e'X$



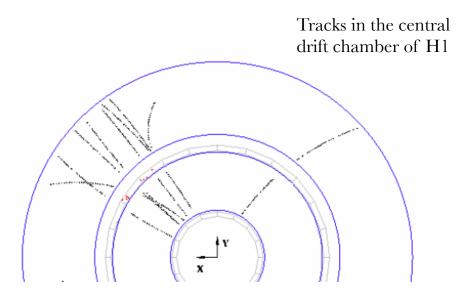
Momentum Measurement

Lorentz force

$$\vec{F} = q \cdot \vec{v} \times \vec{B}$$
$$F = qvB\sin(\vec{v},\vec{B})$$

In solenoidal field: transverse motion perpendicular to B:

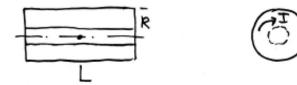
 $F = qvB = \frac{mv^2}{r}$ centrifugal force $mv = p_T$ $p_T = 0.3 \cdot B \cdot r$ [GeV] [T] [m]



Charged particles ionize gas of a tracking chamber: the ionisation is detected on wires and the track is reconstructed using the drift time information. Hits are found to be on a curved track - in the plane transverse to the beam, i.e. the solenoidal field. The curvature is the larger the smaller the transverse momentum, $p_{\rm t}$, of the track is.

Solenoid - Magnet

Magnetic field parallel to the beam axis



$$B(0,0) = \frac{\mu_0 \cdot W \cdot I}{\sqrt{L^2 + 4R^2}} \approx \mu_0 \cdot w \cdot I$$

Stored energy

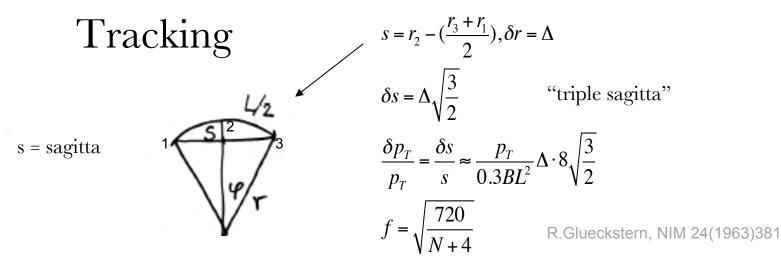
 $E = \frac{1}{2\mu_0} \cdot \int B^2 dV$ $E \approx \frac{1}{2\mu_0} \cdot \pi R^2 \cdot L \cdot B^2$

Keep dimensions small to keep cost down.

Try to arrange the magnet to contain the calorimeter

L = 5.7mW = 1683R = 0.96mI = 10kA $\mu_0 = 4\pi \cdot 10^{-7} N \cdot A^{-2}$ $B = 3.5T, [T = \frac{N}{Am}]$ $E(LHeC) \approx 80MJ, [J = Nm]$ $prize \sim \frac{1}{2} (E / MJ)^{0.66}$ $\sim 9 MioUS$ \$

H1: L=5.7m, R=3m, I=6kA,W=1238, B=1.13T, E=100MJ



$$s = r - r\cos\frac{\varphi}{2} = r \cdot 2\sin^2\frac{\varphi}{4}$$
$$s = \frac{r}{8}\varphi^2 = \frac{L \cdot \varphi}{8}$$
$$\varphi = \frac{L}{r} = \frac{0.3 \cdot B \cdot L}{p_T}$$
$$s = \frac{0.3 \cdot B \cdot L^2}{8p_T}$$

[note clash in notation: r is the radius of curvature, L is the track length, which in a solenoidal field determined the radial extension of the tracking detector ..]

$$\frac{\delta p_T}{p_T^2} = \frac{\Delta}{0.3BL^2} \cdot f$$

→ small Δ , high field B, large L (+ 'no' material)

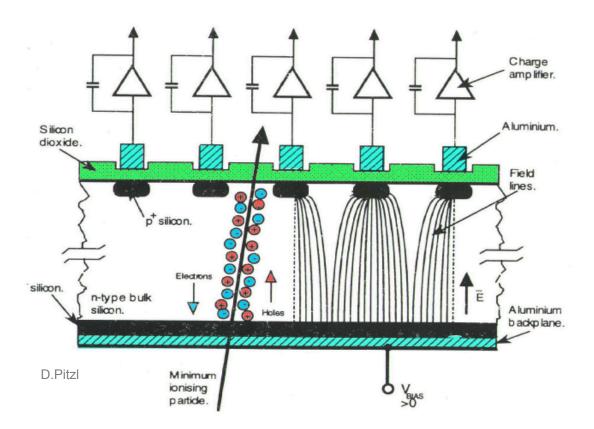
$$CST: \frac{\delta p_T}{p_T^2} \coloneqq 5 \cdot 10^{-4} GeV^{-1}$$

$$B = 3.5T, \Delta \approx 10 \mu m, N = 9, L = 0.4m$$

Given the desired momentum resolution, the field, the point resolution and the number of layers, the radial extension of the tracker is fixed to be ~ 0.5 m

Silicon Detectors

"p in n"



Planar sensor from pure Si wafer segmented with implants (pn)

Strip pitch p~50 μ m, δ ~p/ $\sqrt{12}$ Resolution worsens with inclination

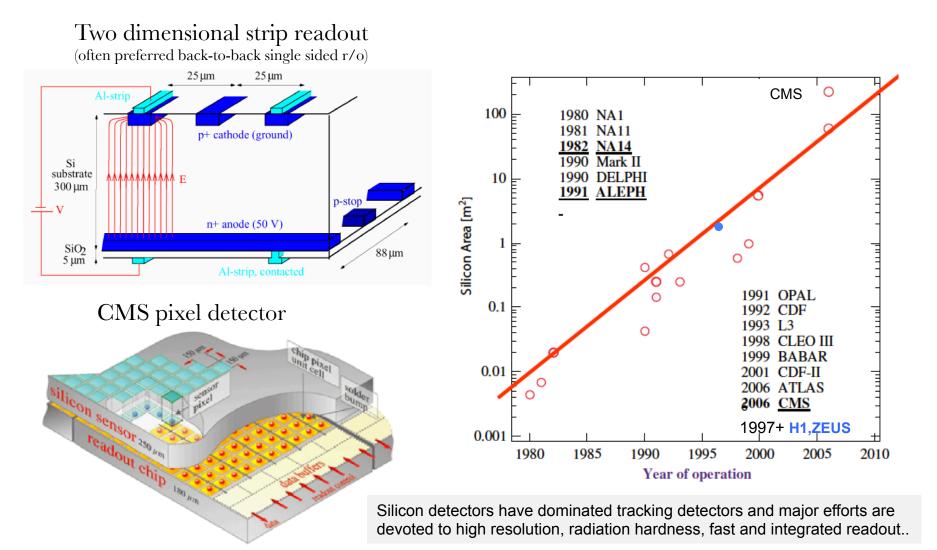
Bulk Si 'depleted' by reverse bias V Ionisation creates e-h pairs e (h) drift time 10 (25) ns

Charge amplified for strip Pedestal defined by random noise

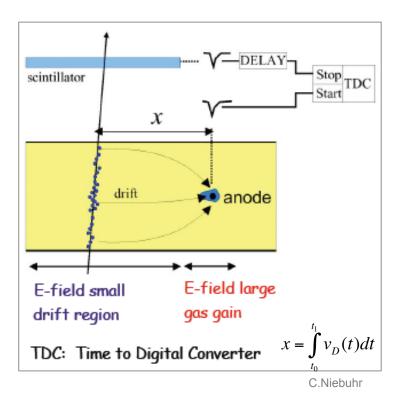
Visit ground floor at "Oliver Lodge"

Modern semiconductor detectors use charge deposited by minimum ionising particle for high precision measurement of charged particle tracks to ~10µm hit reconstruction precision. Signal: ~20ke in 250µm Si, Noise: interstrip C~ 1k, leakage current 0.1k, bias resistors 0.2k [BST H1]

Segmentation and History

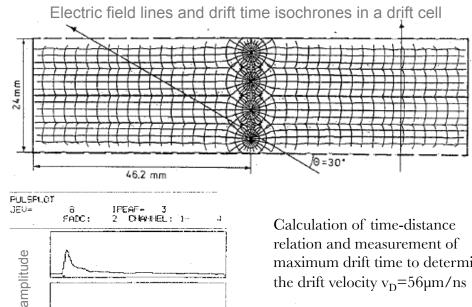


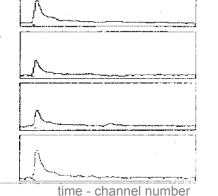
Drift Chambers



Large sensitive volume Point resolution of order 200µm Typical: 1m radius, 10⁴ wires positioned to 20µm precision, strung parallel to beam or circular around ("z chambers").

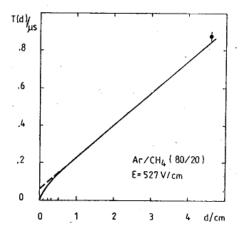
Sensitive to ageing, gas impurities, wire and wire-feedthrough mechanics.. Example: COZ of H1 (PHE 88-02, Zeuthen)

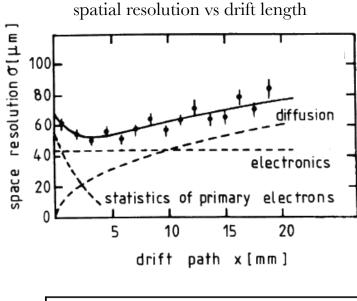


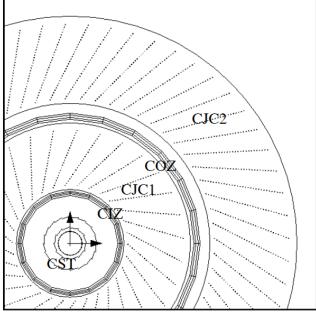


Analog pulses to determine the spatial resolution using different pulse shape analyses. Second coordinate from readout at both ends (charge division)

maximum drift time to determine the drift velocity $v_D = 56 \mu m/ns$







Drift Chambers

Outer H1 central jet chamber (CJC2)



 $15\ 000$ wires,, $6\ tons$ on endflange, separating CJC from forward tracker

Radial view of H1 central tracker

COZ H1 drift chamber built at Zeuthen 1986-91

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