



Fröhlich Lecture Series in Physics 2011/2012

“Particle Physics in Transition”

Professor Max Klein
Department of Physics
University of Liverpool

16.00 Wednesday 26th October 2011
Rotblat Lecture Theatre, Chadwick Building

The Fröhlich Lectures are presentations by research leaders which are intended to be accessible to a general audience at the advanced undergraduate level.

Abstract

Elementary particle physics is now about a hundred years old. A variety of experimental discoveries has supported its Standard Model, a quantum field theory of the interactions of leptons and quarks, as the adequate description of phenomena at the Fermi energy scale. With the Large Hadron Collider at CERN and a number of further energy frontier and high precision accelerator projects we expect to be at the eve of changes and likely surprises in the persisting endeavour to resolve the smallest dimensions of nature.

**Physics
Projects**

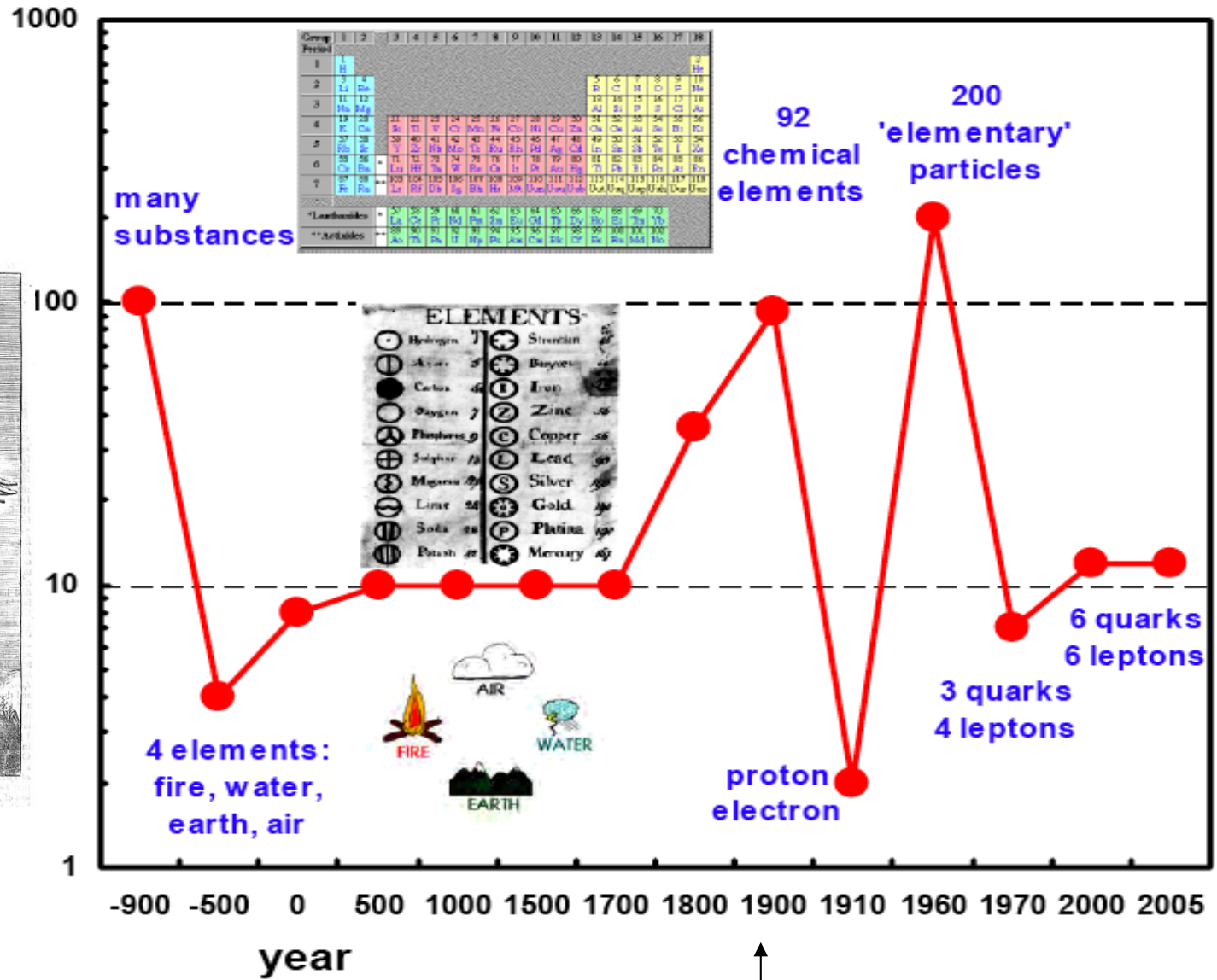
From Empedoclez to today



Empedocle's.

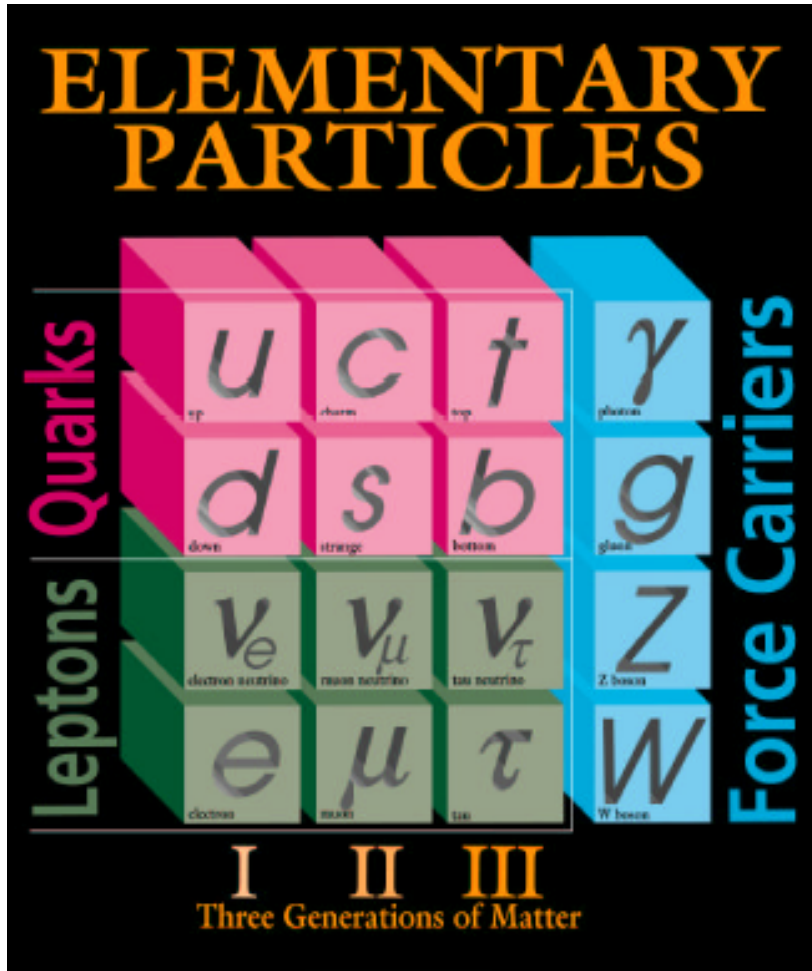
490-430 BC, Sicily

Four elements
and two forces
(love and strife)



Lord Kelvin: There is nothing new to be discovered in physics.
All that remains is more and more precise measurement.

The Standard Model



Renormalisable gauge field theory of the electromagnetic, strong and weak interaction.

$$L_{QED} = \bar{\Psi}D\Psi + m\bar{\Psi}\Psi + (DA)^2 + eA\bar{\Psi}\Psi$$

e propagator e mass photon interaction

virtual particles: quantum theory

perturbative QED

Higgs mechanism to damp divergence of WW cross section. H the “god” particle, a scalar.

Neutrinos

Offener Brief an die Gruppe der Radioaktiven bei der
Gauvereins-Tagung zu Tübingen.

Abschrift

Physikalisches Institut
der Eidg. Technischen Hochschule
Zürich

Zürich, 4. Dez. 1930
Diorastrasse



Wolfgang Pauli

Liebe Radioaktive Damen und Herren,

Wie der Ueberbringer dieser Zeilen, den ich halbvollst
anzuhören bitte, Ihnen das näheren auseinandersetzen wird, bin ich
angesichts der "falschen" Statistik der N- und Li-6 Kerne, sowie
des kontinuierlichen beta-Spektrums auf einen **verzwäifelten Ausweg**
verfallen um den "Wechselatz" (1) der Statistik und den **Energiesatz**
zu retten. Nämlich die Möglichkeit, es könnten elektrisch neutrale
Teilchen, die ich Neutronen nennen will, in den Kernen existieren,
welche den Spin 1/2 haben und das Ausschliessungsprinzip befolgen und
sich von Lichtquanten ausserdem noch dadurch unterscheiden, dass sie
nicht mit Lichtgeschwindigkeit laufen. Die Masse der Neutronen
müsste von derselben Grosseordnung wie die Elektronenmasse sein und
jedenfalls nicht grösser als 0,01 Protonenmasse.- Das kontinuierliche
beta-Spektrum wäre dann verständlich unter der Annahme, dass beim
beta-Zerfall mit dem Elektron jeweils noch ein Neutron emittiert
wird, derart, dass die Summe der Energien von Neutron und Elektron
konstant ist.

$$n \rightarrow p e \nu$$

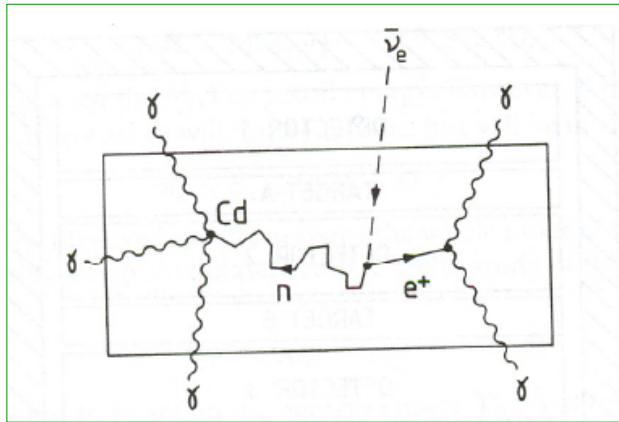
Pauli in a letter to the "radioactive ladies and gentleman" postulates the neutron/ino
Too immature an idea to publish and o busy a man to participate in the workshop..

Early Neutrino Physics



Fred Reines, Clyde Cowan
Savannah River Reactor
 $\nu p \rightarrow n e^+$ 1956

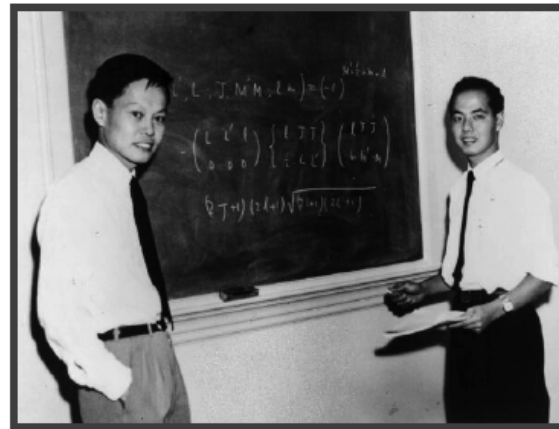
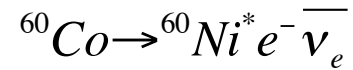
Discovery of the neutrino



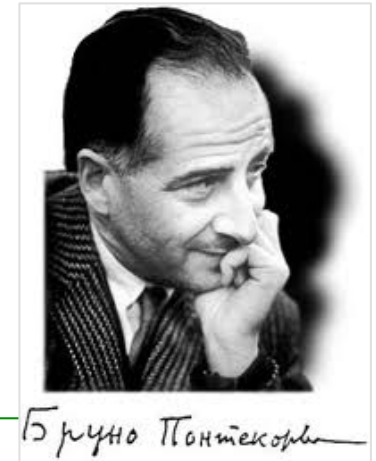
$E_\nu = 1 \text{ MeV}$. 2 events / hour

Parity violation.

Mme Wu et al., 1957



Parity predicted to be violated
by T.D.Lee and C.N. Yang in 1956
("θ-τ puzzle")



Bruno Pontecorvo

Inverse beta processes and nonconservation of lepton charge

Sov.Phys.JETP 7 (1958) 172-173.

The foundation of neutrino oscillation physics.

The standard picture of three neutrino mixing

- flavor oscillation described by PNMS matrix
- parametrized by 3 mixing angles and CP-violating phase δ_{CP}

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \theta_{23} & \sin \theta_{23} \\ 0 & -\sin \theta_{23} & \cos \theta_{23} \end{pmatrix} \cdot \begin{pmatrix} \cos \theta_{13} & 0 & \sin \theta_{13} e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -\sin \theta_{13} e^{-i\delta_{CP}} & 0 & \cos \theta_{13} \end{pmatrix} \cdot \begin{pmatrix} \cos \theta_{12} & \sin \theta_{12} & 0 \\ -\sin \theta_{12} & \cos \theta_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \cdot \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

“atmospheric sector”

θ_{23}

$$\begin{aligned} |\Delta m_{31}^2| & (2.40^{+0.12}_{-0.11}) 10^{-3} \text{ eV}^2 \\ \sin^2 \theta_{23} & 0.50^{+0.07}_{-0.06} \end{aligned}$$

θ_{13}

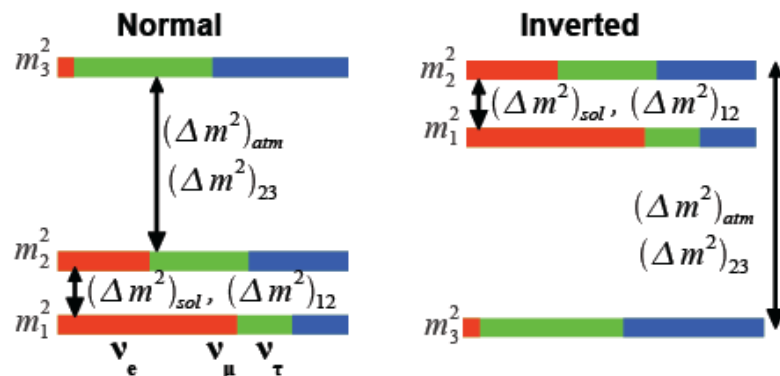
$\sin^2 \theta_{13}$	$0.013^{+0.007}_{-0.005}$ $0.016^{+0.008}_{-0.006}$
δ	$(-0.61^{+0.75}_{-0.65}) \pi$ $(-0.41^{+0.65}_{-0.70}) \pi$

new

“solar sector”

θ_{12}

$$\begin{aligned} \Delta m_{21}^2 & (7.65^{+0.23}_{-0.20}) 10^{-5} \text{ eV}^2 \\ \sin^2 \theta_{12} & 0.304^{+0.022}_{-0.016} \end{aligned}$$

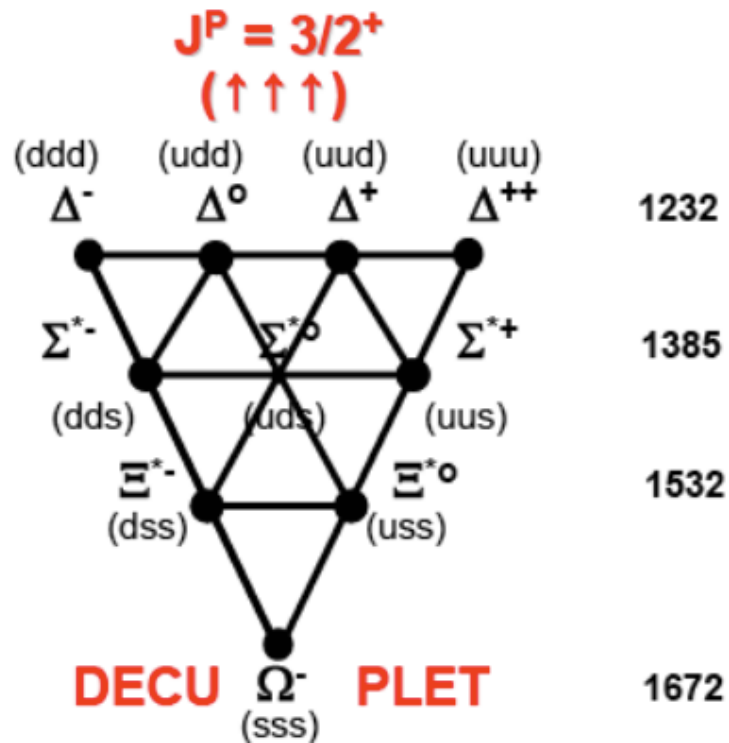


Most urgent points:

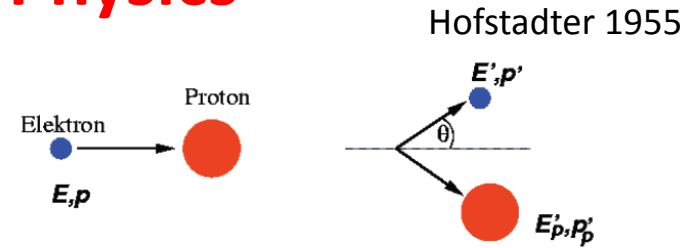
- $\sin^2 2\theta_{13} > 0.01$ at $>5\sigma$ significance ?
- Mass hierarchy $\Delta m_{31}^2 > 0$? , $\Delta m_{31}^2 < 0$?
- CP-phase $\delta \neq 0, \pi$ at $>3\sigma$ significance, δ true ?
- Unitarity ? tri-bimaximal ? differences between quark and lepton sectors ?

Early Quark-Parton Physics

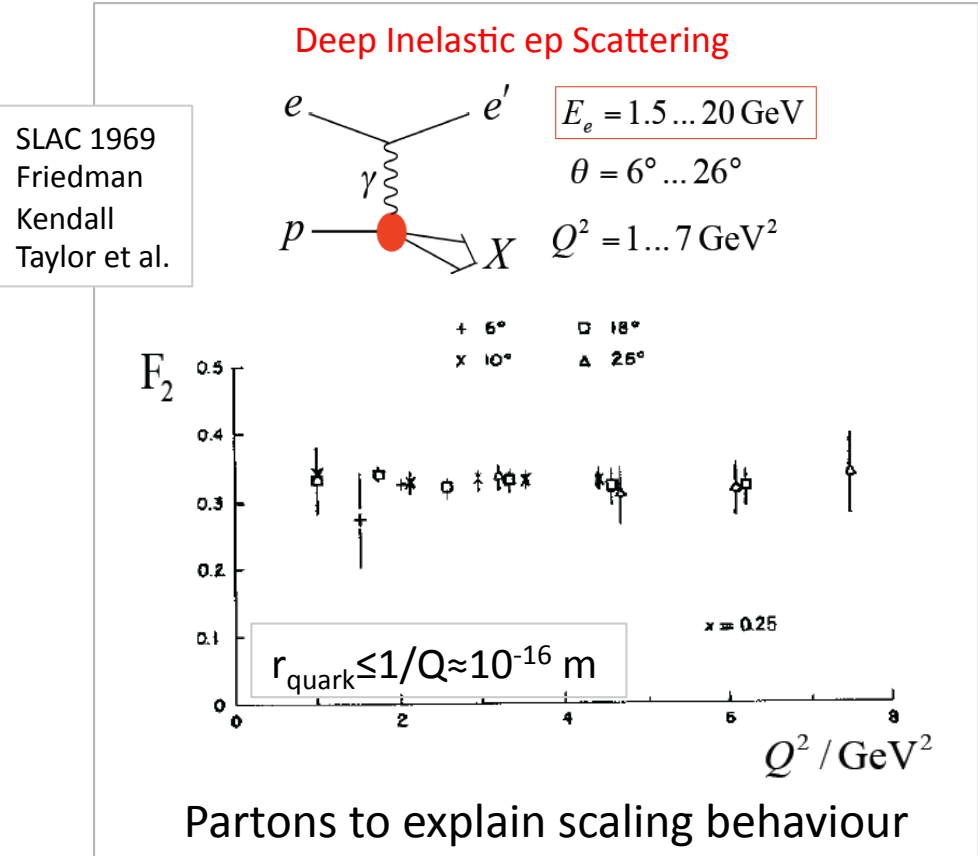
[Three] quarks to explain the proliferation of particles



Ω^- predicted 1962: Gell-Mann $M=1685$ MeV
observed at BNL 1964 at $M=1686 \pm 12$ MeV

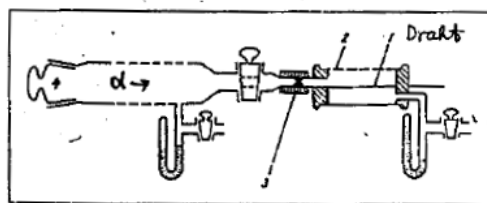


$$r_{\text{Proton}} = (0.74 \pm 0.24) \cdot 10^{-15} \text{ m}$$

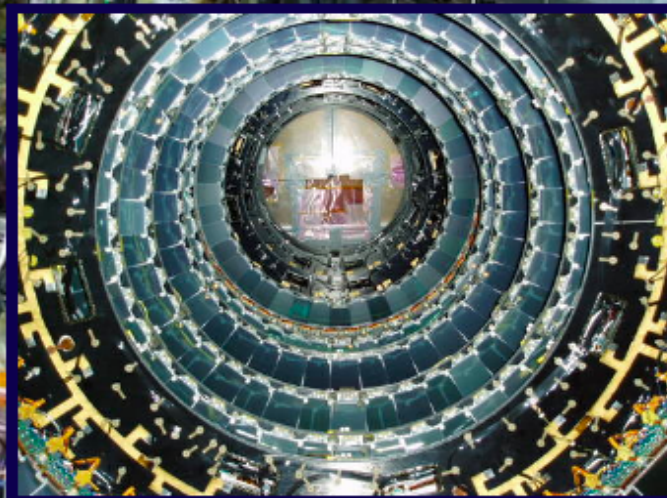


A gigantic evolution of instrumentation, electronics and computing over 100 years

ATLAS

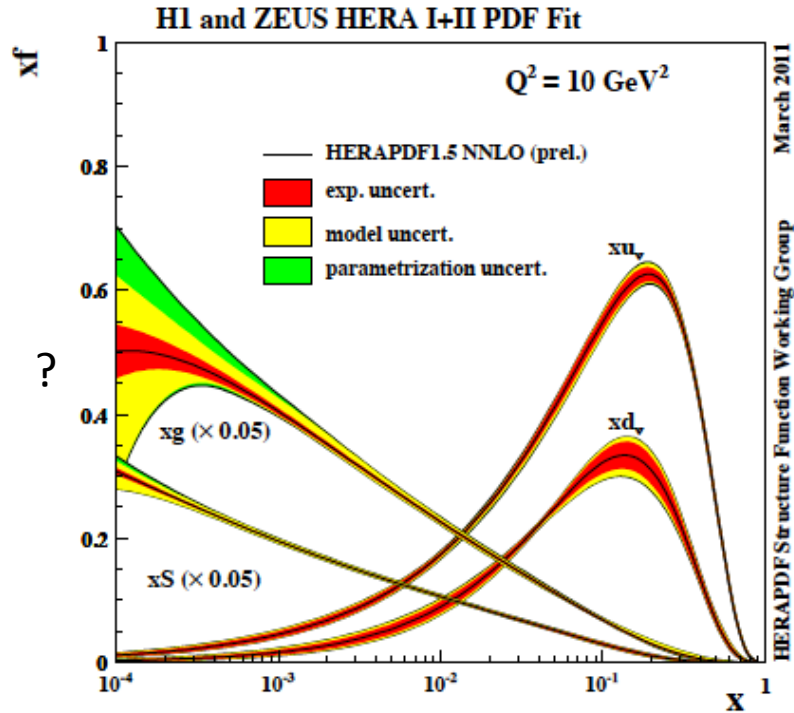


Rutherford Phil Mag 21 (1911)66

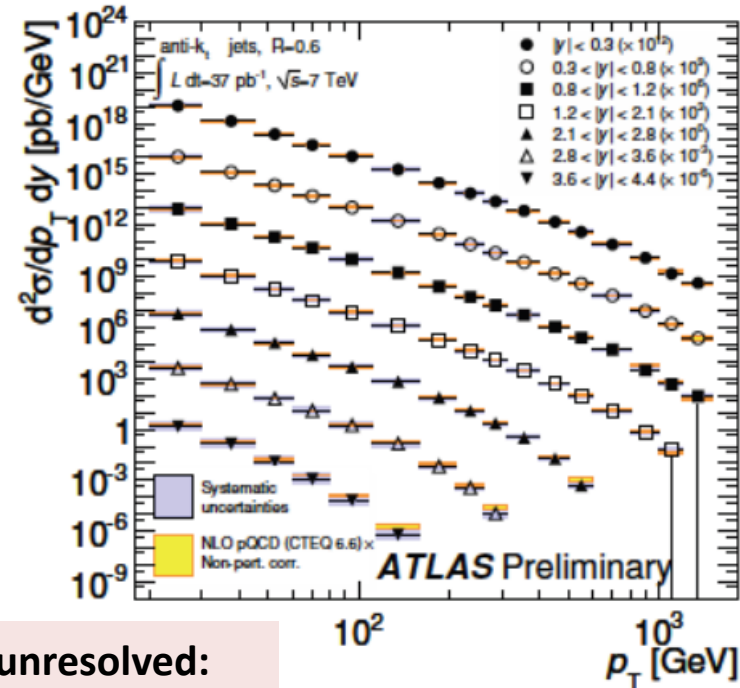


University of Liverpool - SCT.
part of inner track Silicon det.

The standard picture of quark-parton dynamics



Universal? parton distributions



Yet unresolved:

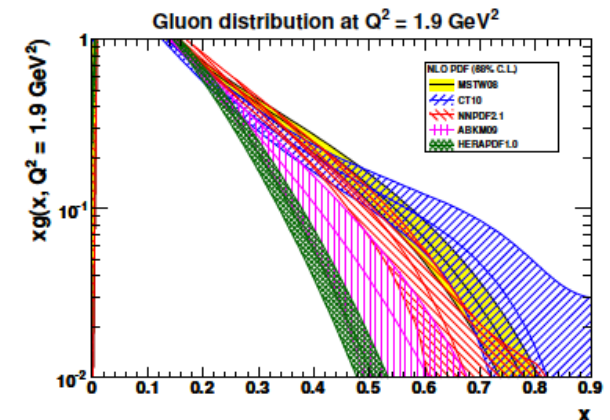
- Gluon at high x/M
- GUT (α_s to 0.1%)
- QPM symmetries
- Saturation
- Axions, Odderons
- Instantons
- Initial QGP state
- N-PDFs
- QCD-String-SUSY
- Substructure..

$$\mu \frac{dg(\mu)}{d\mu} = \beta(g(\mu)) \quad \text{RGE}$$

$$\beta(g) = -g \left[\frac{\alpha_s}{4\pi} \beta_1 + \left(\frac{\alpha_s}{4\pi} \right)^2 \beta_2 + \dots \right]$$

$$\frac{\alpha_s}{4\pi} = \frac{1}{\beta_1 \ln(\mu^2 / \Lambda^2)} - \frac{\beta_2 \ln(\ln \mu^2 / \Lambda^2)}{\beta_1^3 \ln^2(\mu^2 / \Lambda^2)} + \dots$$

Asymptotic Freedom: perturbative QCD



The Higgs Mechanism

Broken Symmetry (as in $M_p \neq M_n$):

non-invariant terms in L or [Nambu] L is symmetric but the ground state isn't

Spontaneous breaking of a global symmetry leads to the existence of massless scalar bosons [Goldstone]

Spontaneous breaking of a locally symmetric gauge theory [Yang Mills] implies:

Goldstone bosons disappear, massless gauge bosons become massive and a massless scalar particle appears [Higgs, Brout, Englert, Kibble, Hagen..]

[“The gauge fields have eaten the Goldstone bosons and grown heavy” Coleman]

BROKEN SYMMETRIES, MASSLESS PARTICLES AND GAUGE FIELDS

P. W. HIGGS

Tait Institute of Mathematical Physics, University of Edinburgh, Scotland

Received 27 July 1964

Phys.Lett.12(1964)132-133

BROKEN SYMMETRIES AND THE MASSES OF GAUGE BOSONS

Peter W. Higgs

Tait Institute of Mathematical Physics, University of Edinburgh, Edinburgh, Scotland

(Received 31 August 1964)

Phys.Rev.Lett.13(1964)508-509

5+5 since 58

In the SM have 4 gauge fields ($A^{1,2,3}, B$) and the Higgs field (ϕ). The gauge fields are transformed to the mass eigenstates: Z, W^\pm and the photon. **The masses of the weak interaction bosons are determined by the VEV of the Higgs potential.**

Remarks on the Higgs particle

$$L = (D_\mu \varphi)^\dagger (D_\mu \varphi) - V(\varphi^\dagger \varphi) - \frac{1}{4} F_{\mu\nu}^a (F^a)^{\mu\nu} - \frac{1}{4} G_{\mu\nu} G^{\mu\nu}$$

$$D_\mu = \left(\partial_\mu + ig A_\mu^a \frac{\epsilon_a}{2} + ig' B_\mu \frac{1}{2} \right), \varphi = \begin{pmatrix} 0 \\ \eta + \sigma(x)/\sqrt{2} \end{pmatrix}$$

$$(D_\mu \varphi)^\dagger (D_\mu \varphi) = \frac{1}{2} (\partial_\mu \sigma) (\partial^\mu \sigma)$$

$$+ \frac{1}{2} \left(\frac{g^2 \eta^2}{2} \right) (A_\mu^1 A^{1\mu} + A_\mu^2 A^{2\mu})$$

$$+ \frac{1}{2} \eta^2 (g A_\mu^3 - g' B_\mu) (g A^{3\mu} - g' B^\mu) + h.c.$$

$$\begin{pmatrix} Z_\mu \\ A_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} A_\mu^3 \\ B_\mu \end{pmatrix}, \tan \theta = \frac{g'}{g}$$

$$W_\mu^\pm = \frac{1}{\sqrt{2}} (A_\mu^1 \pm i A_\mu^2)$$

$$M_{W^+} = M_{W^-} = \frac{g\eta}{\sqrt{2}}$$

$$M_Z = \frac{M_W}{\cos \theta}$$

$$M_\gamma = 0$$

$$M_H = \sqrt{-2\mu^2} = 2\eta \cdot \sqrt{\lambda}$$

Gauge field Lagrangian a la QED

Covariant Derivative including four gauge fields transforming according to SU(2) [A^{1,2,3}] and U(1) B

φ is a specific choice of the Higgs field, η is its VEV. For $V = -\mu^2 \varphi^\dagger \varphi - \lambda (\varphi^\dagger \varphi)^2$ it follows $\eta = \sqrt{-\mu^2/2\lambda}$

The explicit calculation of the 'DD' term yields a squared, i.e. mass term for an A^{1,2} combination and, after rotation, for Z and A³, the photon!

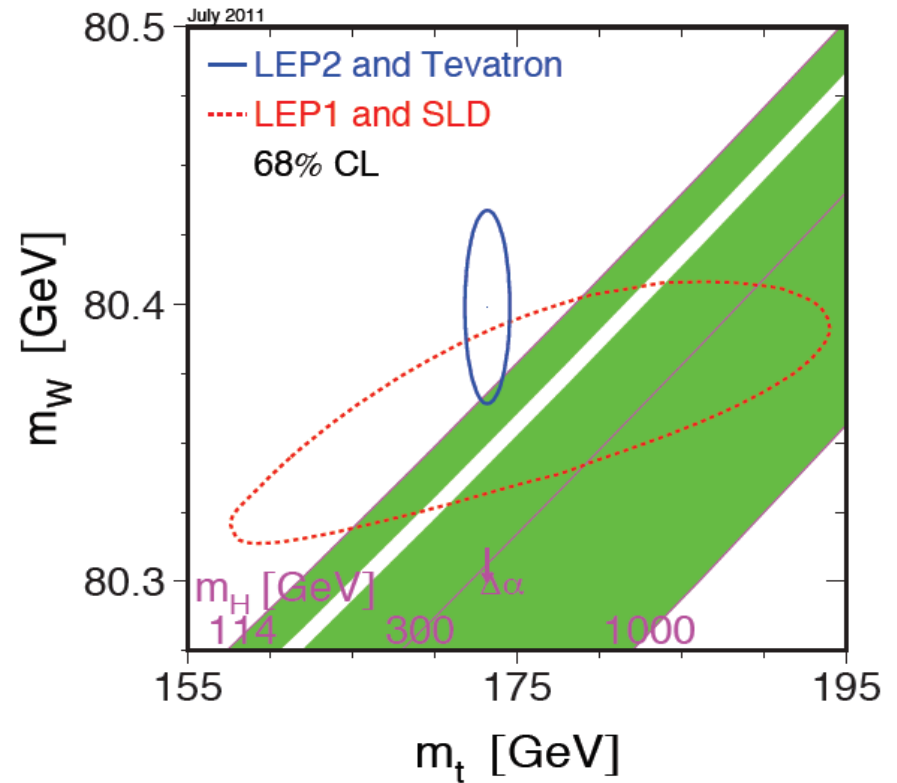
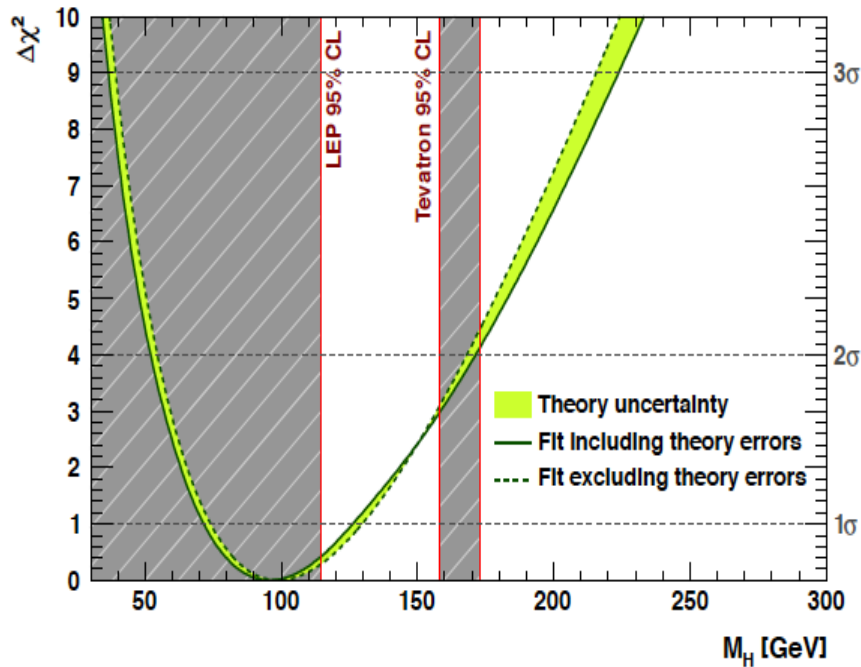
The weak and electromagnetic couplings g' , g get related via the rotation (the "Weinberg") angle. Specifically one has $e = g \sin \theta$. [data: $\sin \theta = 0.23116(13)$]

The masses are related to the VEV of the Higgs field.

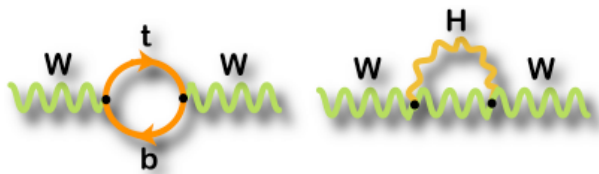
All of this has been beautifully confirmed by experiment: W, Z, $\sin \theta$, mass relations.

We don't know how large λ is and whether this choice of the Higgs mechanism is real.

Search for the Higgs Particle - indirect



Determination of M_H within exploiting loop corrections

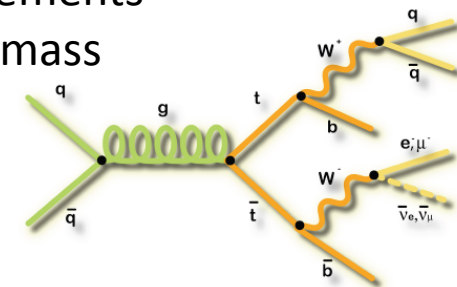


Successful example: prediction of top mass using e^+e^- data

Challenging measurements of W and top quark mass

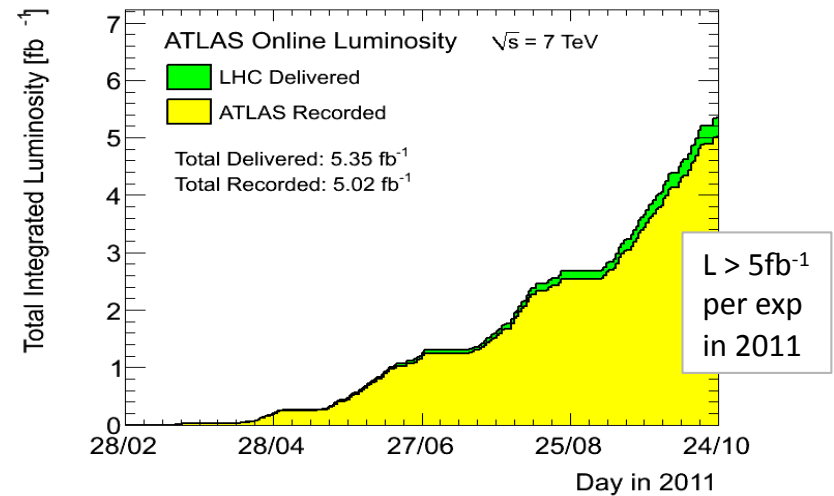
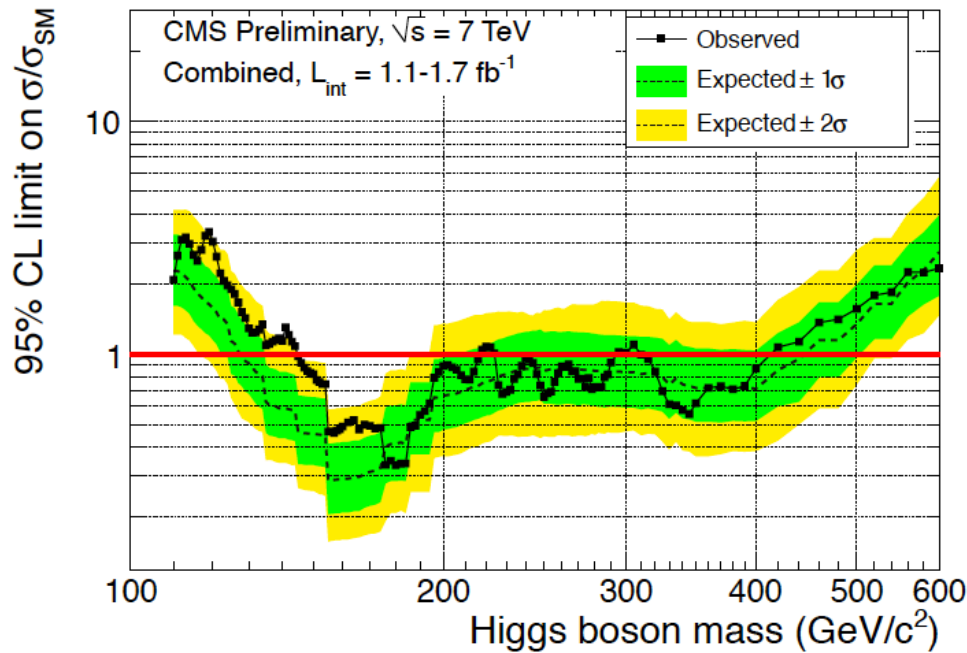
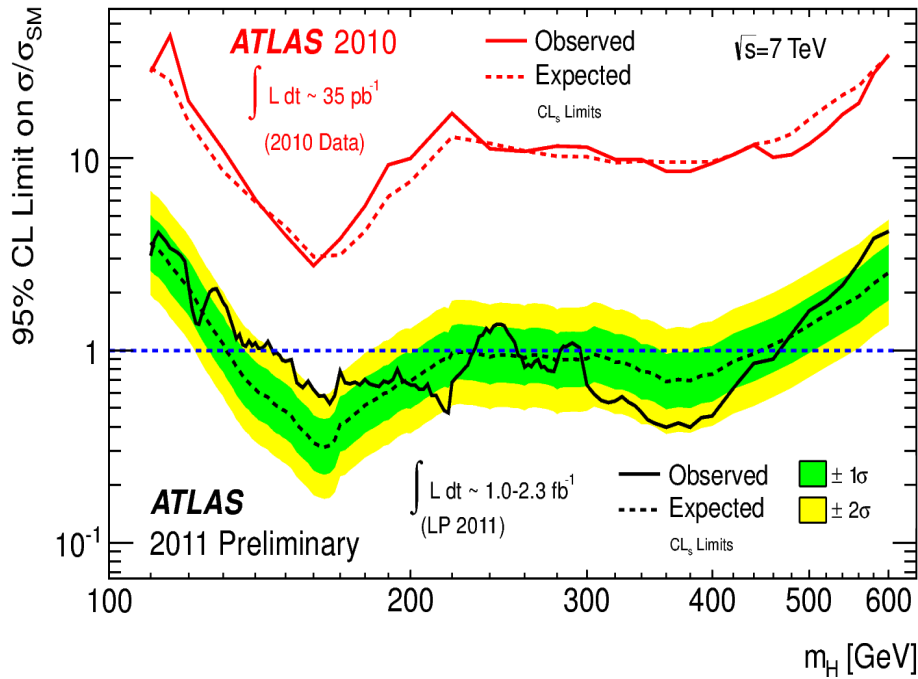
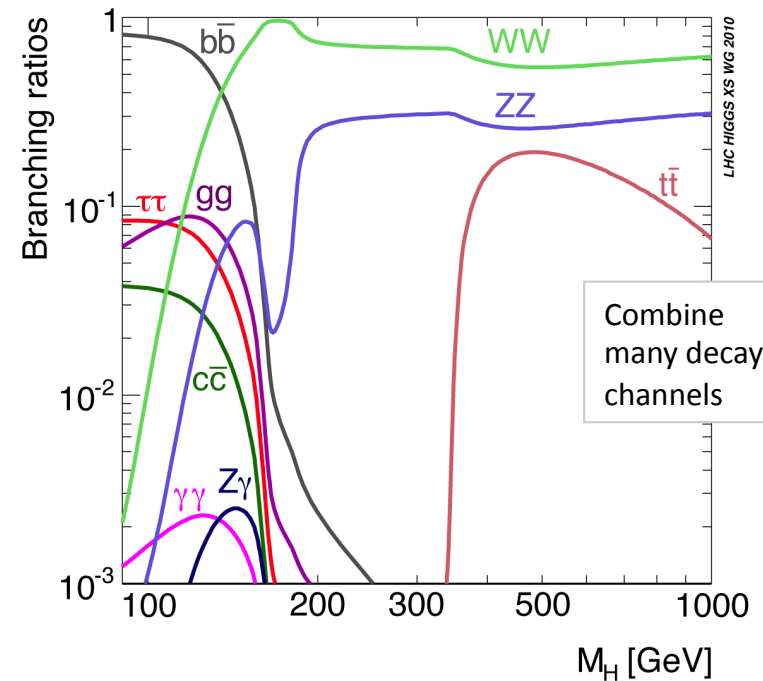
$$m_W = 80.399 \pm 0.023$$

$$m_t = 173.1 \pm 1.3$$



Search for the Higgs Particle

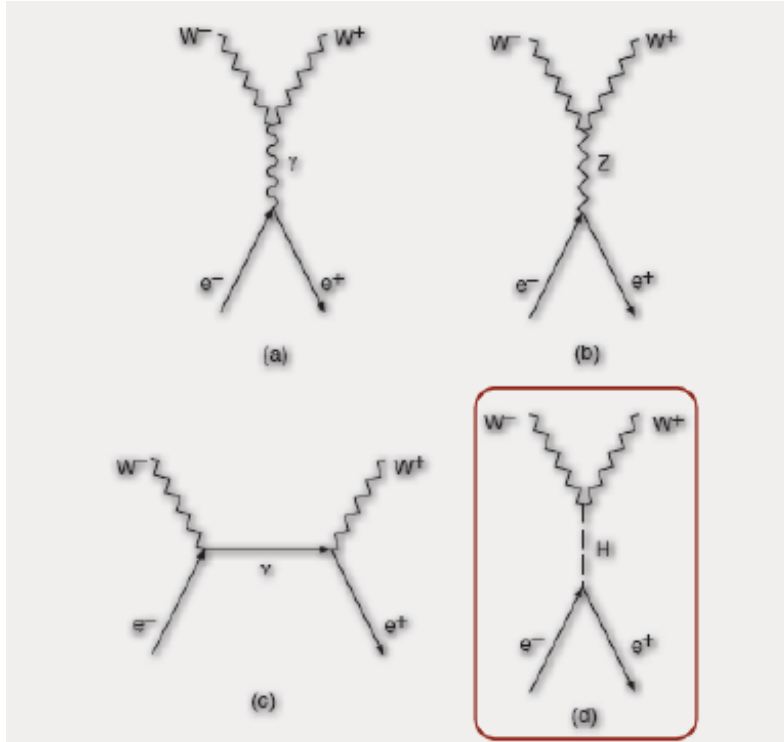
Direct [August 2011]



LHC Tunnel 2002



No Higgs?

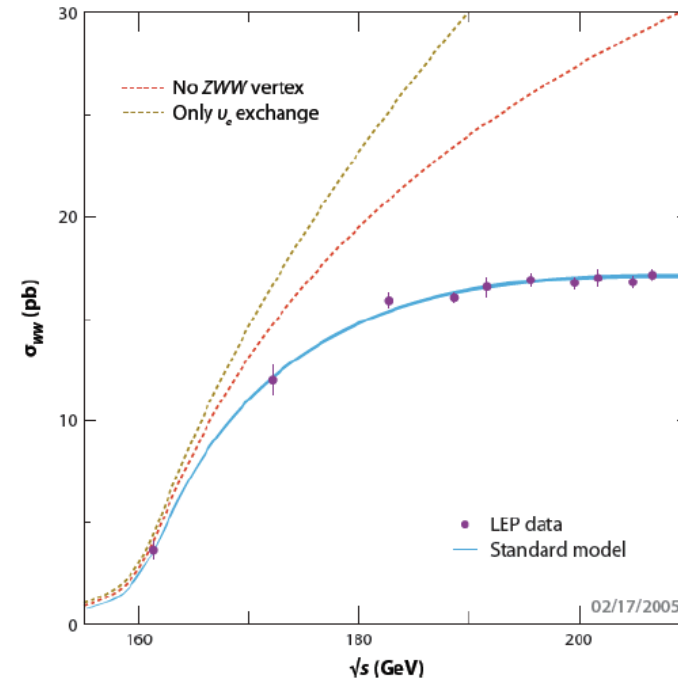


W and Z/ γ behave as predicted

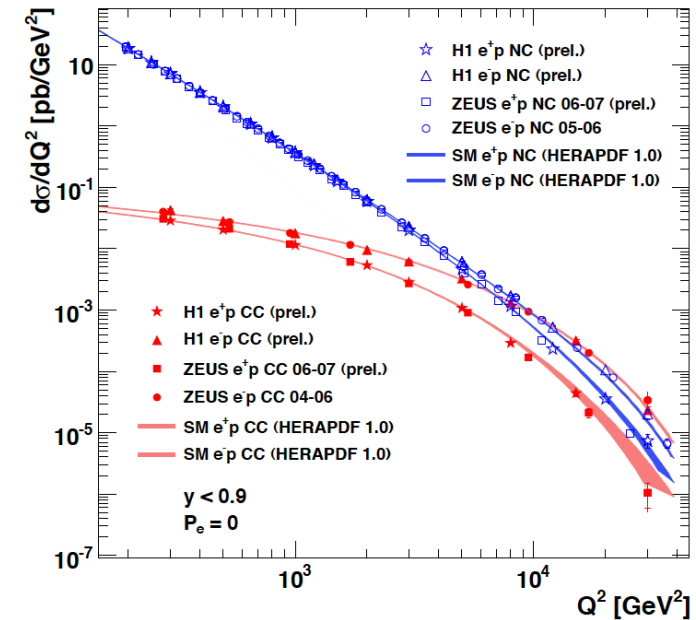
LEP:ZWW

HERA:NC \approx CC

LEP



HERA



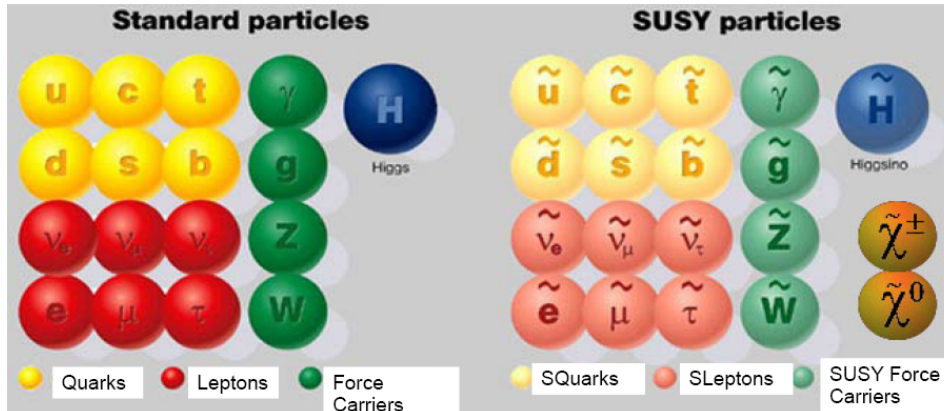
The W damps the rise of the 4fermion cross section.

The H has been expected to give mass to the W,Z but also to damp the rise of the WW cross section.

It is a major Christmas present if H was indeed excluded.

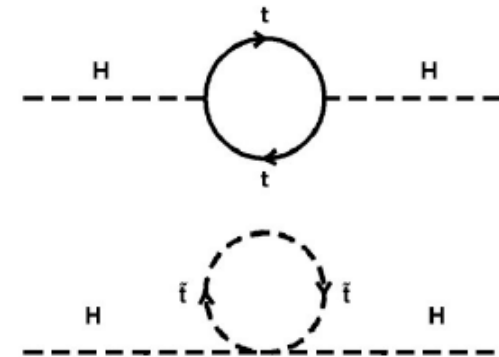
SUperSYmmetry^{*)}

J=1/2 J=1 J=0 J=0 J=1/2 J=1/2



A beautiful theory:

- scalars become normal particles
- divergences are compensated
- matter and field get united
(massless photon and massive photino)
- internal (isospin, colour) symmetries get united with space-time symmetry
(generalisation of Poincaré group)



*) Y.Golfand, E.Likhtman, JETP Letters 13(1971)323
 "Extension of the Algebra of Poincare Group Generators and Violation of P Invariance"
 D.Volkov, V.Akulov, Phys.Lett. 46B (1973)109 "Is the Neutrino a Goldstone Particle?"
 J.Wess, B.Zumino, Nucl.Phys. B70(1974)39 "Supergauge Transformations in 4 Dimensions"

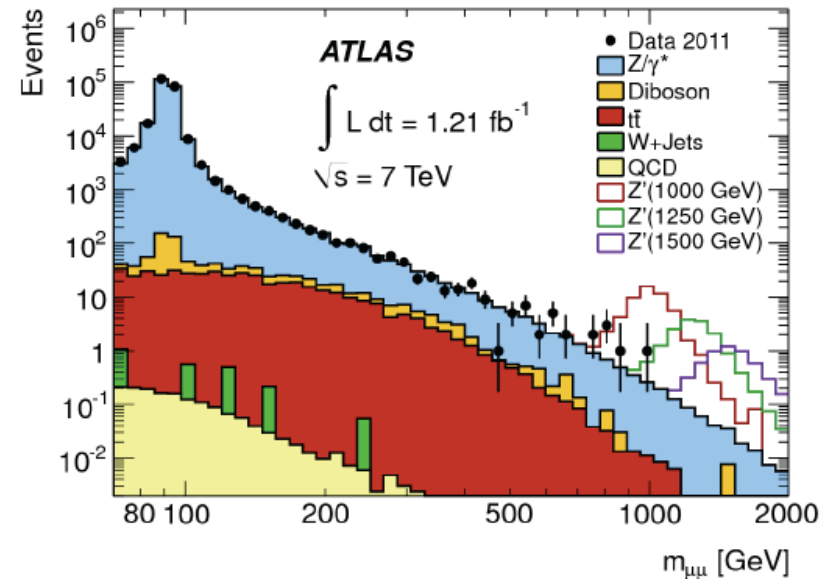
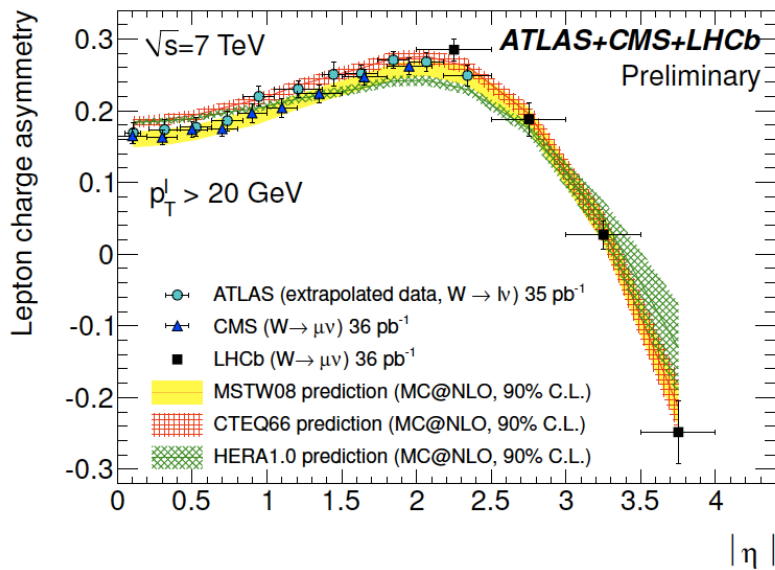
SUSY needs "light" Higgs (most of it)

Technicolor:

"We argue that the existence of fundamental scalar fields constitutes a serious flaw of the Weinberg-Salam theory. A possible scheme without such fields is described. The symmetry breaking is induced by **a new strongly interacting sector whose natural scale is of the order of a few TeV.**"

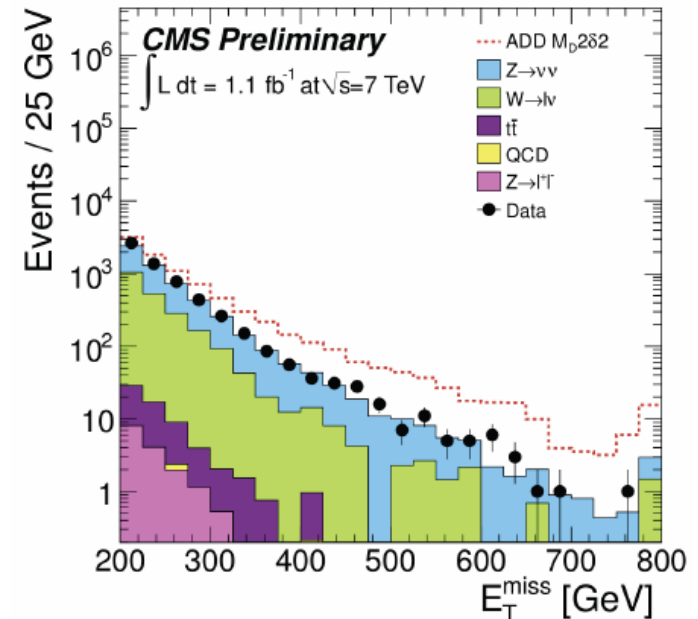
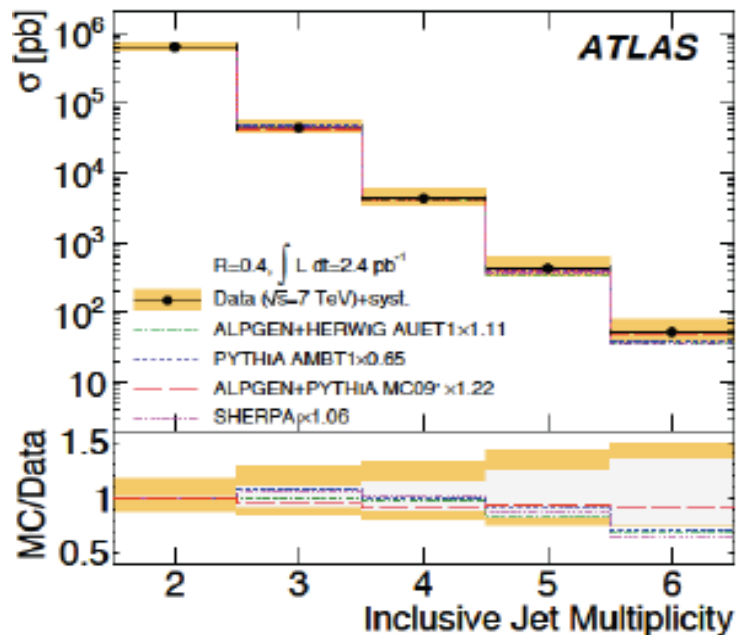
L.Susskind, Dynamics of Spontaneous Symmetry Breaking in the Weinberg Salam Theory. Phys D20 (1979) 2619-2625
 Dimopoulos, Susskind: Mass Without Scalars NP. B155 (1979) 237 Farhi, Susskind: Technicolor Phys.Rept. 74 (1981) 277

LHC- Standard Model Measurements [~400 papers/preliminary results in 2010/11]



No sign for Z' nor W' below about 2 TeV

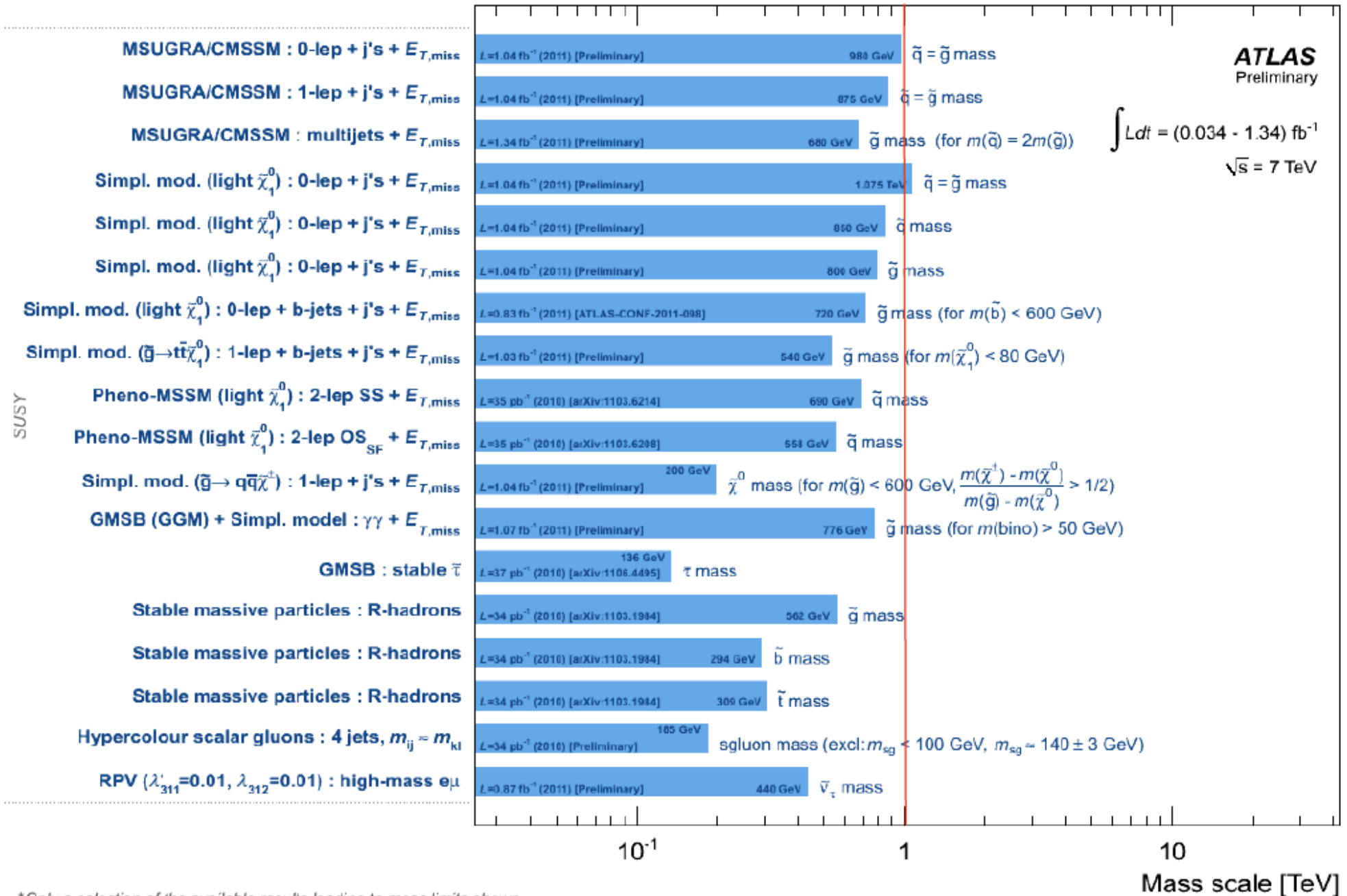
Explore proton+QCD in new region



No large missing energy of unknown origin

ATLAS Limits to SUSY

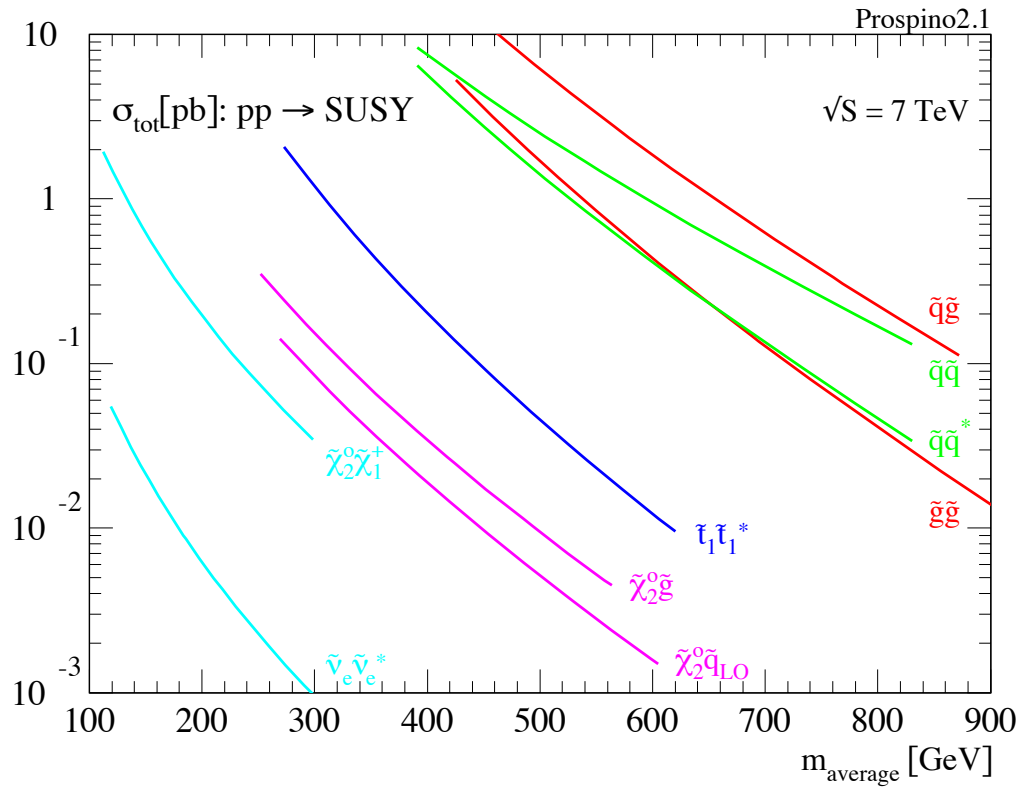
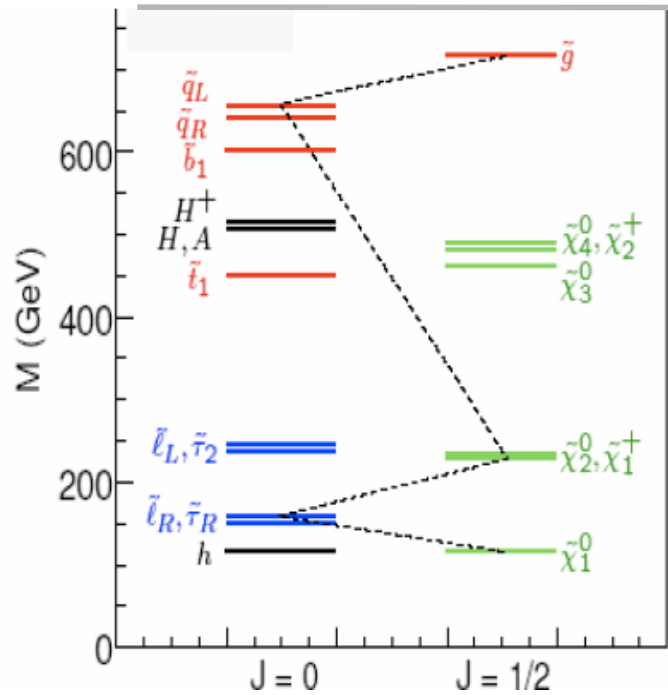
ATLAS Searches* - 95% CL Lower Limits (Status: SUSY 2011)



*Only a selection of the available results leading to mass limits shown

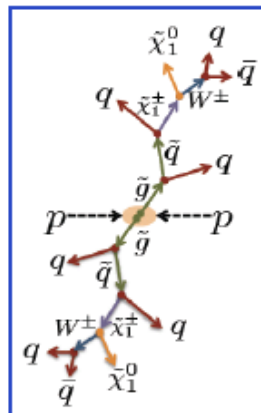
SUSY Before and Now?

A SUSY spectrum before 2011



T.Plehn

Look for
3rd, lighter generation
RPV SUSY [ep, DM?]
Extended decay chains



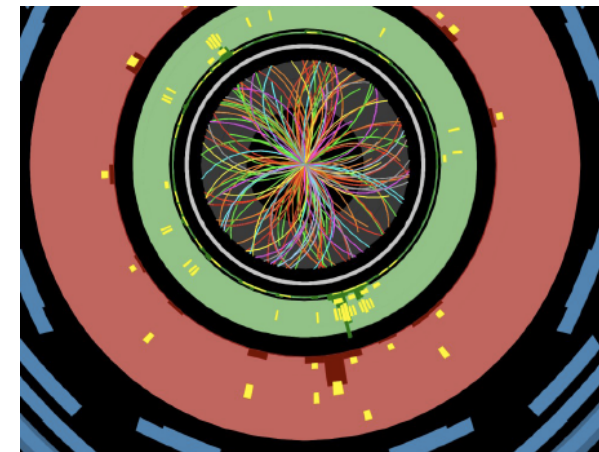
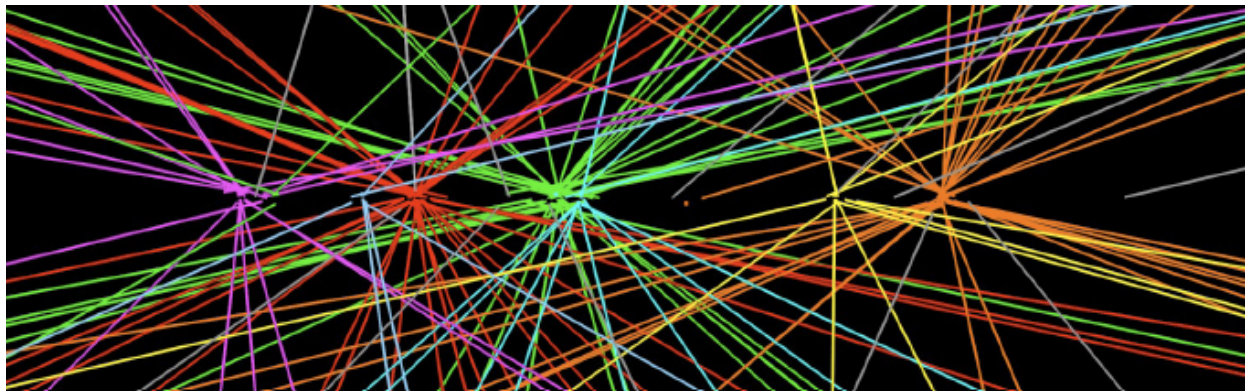
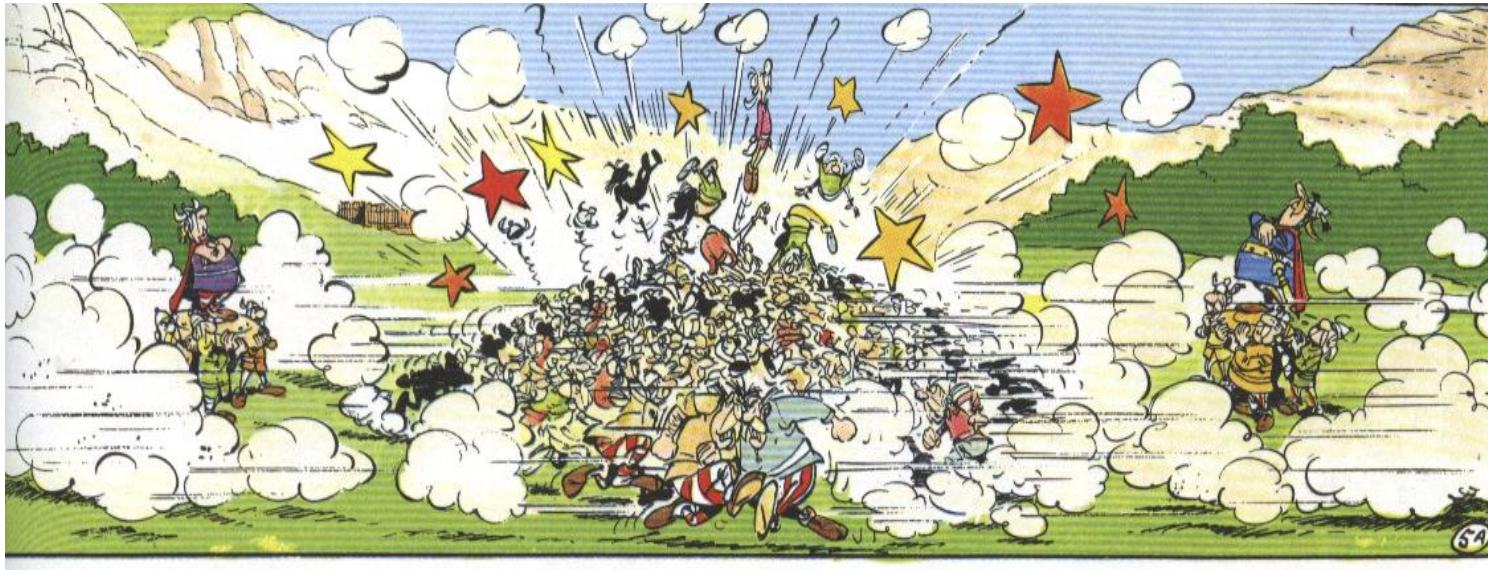
$$R\text{-parity} = (-1)^{2J + 3B + L}$$

= 1 for SM particles, = -1 for SUSY partners

IF SUSY exists and IF R is conserved then the lightest SUSY particles are a candidate for DM

[proton decay violates R parity; R ad hoc in MSSM but natural in SO(10) GUT]

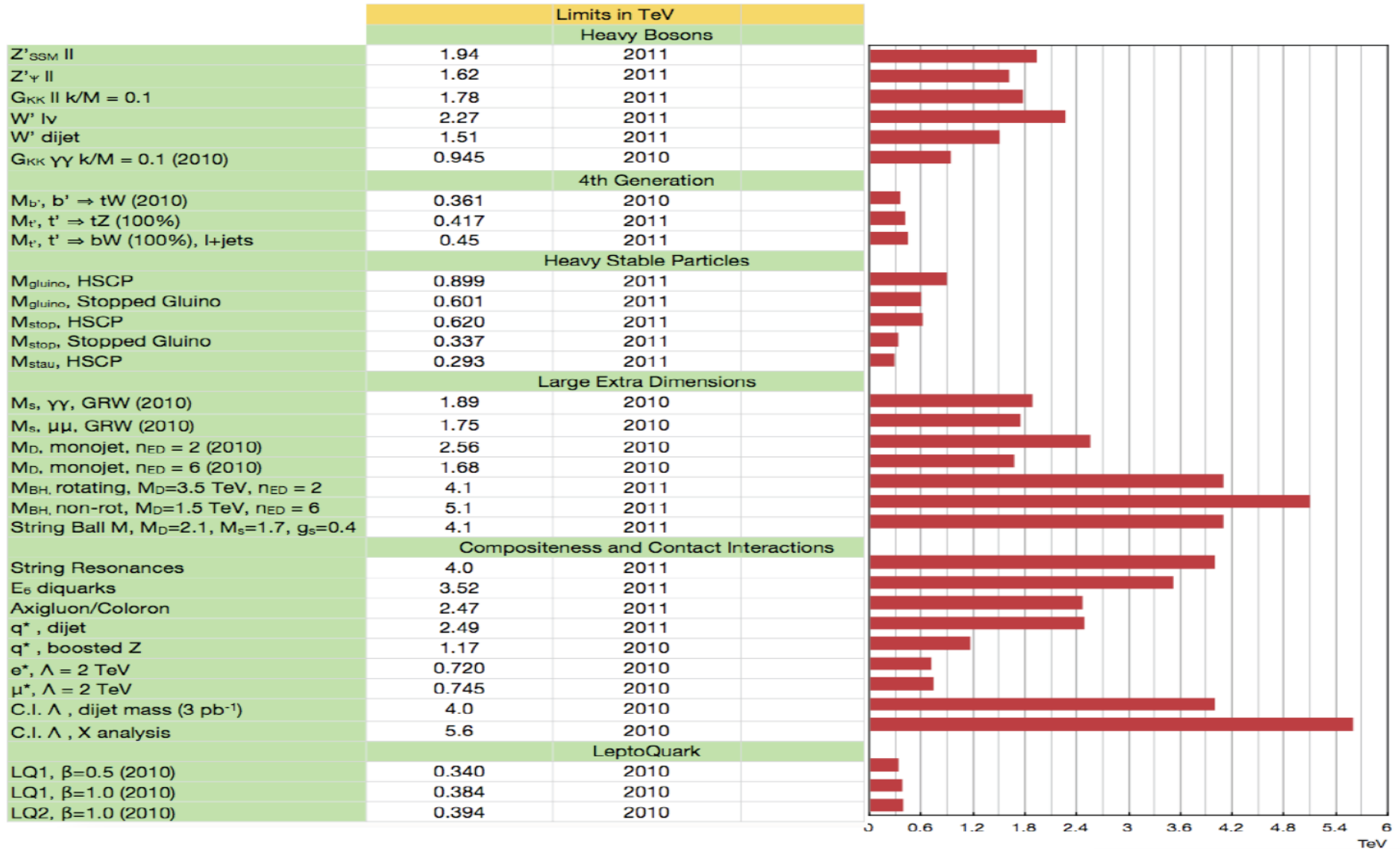
An unprecedented experimental challenge



Mean pileup about 10 at $L=10^{33}\text{cm}^{-2}\text{s}^{-1}$

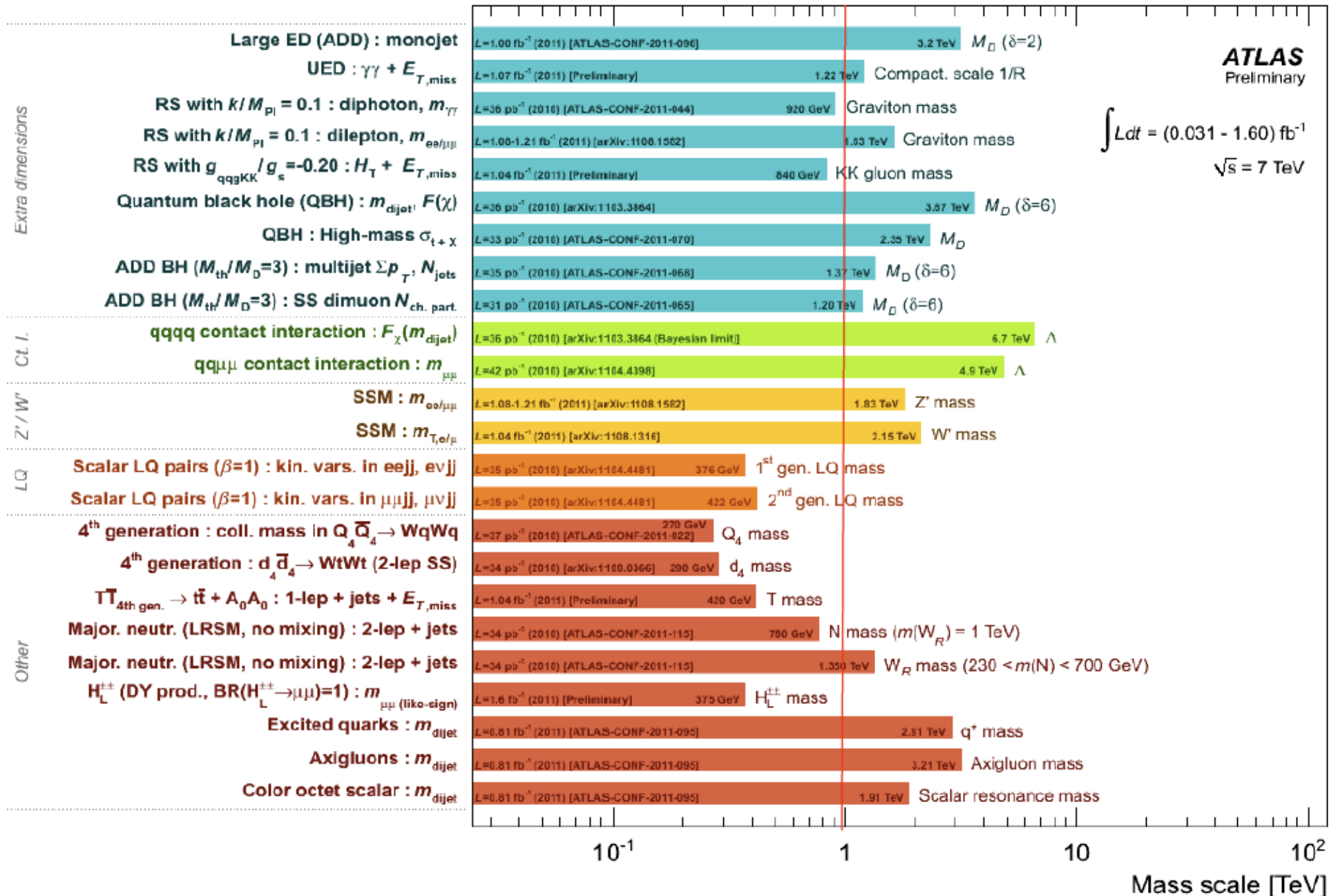
Need full replacement of inner ATLAS tracker for $5 \cdot 10^{34}$

CMS - Limits to Exotic Physics



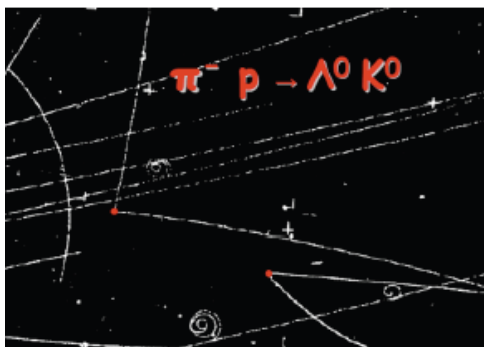
ATLAS Exotics Limits

ATLAS Searches* - 95% CL Lower Limits (Status: SUSY 2011)



Note: LHC is for $\sim 1000 \text{ fb}^{-1}$ at 14 TeV to be compared with 1 fb^{-1} at 7 TeV now

Early Quark Mixing Physics



$$\tau(K_s^0 \rightarrow 2\pi) = 0.9 \cdot 10^{-10} s$$

$$\tau(K_L^0 \rightarrow 3\pi) = 0.5 \cdot 10^{-7} s$$

$$CP|K^0\rangle = |\bar{K}^0\rangle$$

$$|K_S^0\rangle = \sqrt{\frac{1}{2}}(|K^0\rangle + |\bar{K}^0\rangle), CP = +1$$

$$|K_L^0\rangle = \sqrt{\frac{1}{2}}(|K^0\rangle - |\bar{K}^0\rangle), CP = -1$$

Strangeness oscillations: $m_L - m_S = 3.5 \cdot 10^{-6} \text{ eV}$
(Gell-Mann and Pais, 1955)

Regeneration of K_S component (Pais et al. 1956)

Rare decay of K_L into 3 pions: CP violation
(Cronin and Fitch 1964)

Suppression of strangeness changing weak charged currents

Cabibbo 1963

$$\begin{pmatrix} u \\ d' \end{pmatrix}$$

$$d' = d \cos \theta_c + s \sin \theta_c$$

$$K^+ (\bar{s}u) \rightarrow \pi^0 (\bar{u}u) e^+ \nu_e$$

$$\pi^+ (\bar{d}u) \rightarrow \pi^0 (\bar{u}u) e^+ \nu_e$$

$$\frac{\Gamma(K^+ \rightarrow \pi^0 e^+ \nu_e)}{\Gamma(\pi^+ \rightarrow \pi^0 e^+ \nu_e)} \propto \tan^2 \theta_c$$

Suppression of strangeness changing weak neutral currents

Glashow, Iliopoulos, Maiani (GIM) 1970

$$\begin{pmatrix} u \\ d' \end{pmatrix}, \begin{pmatrix} c \\ s' \end{pmatrix}$$

$$d' = d \cos \theta_c + s \sin \theta_c$$

$$s' = -d \sin \theta_c + s \cos \theta_c$$

$$\frac{\Gamma(D^+ \rightarrow K^- \pi^+ \pi^+)}{\Gamma(D^+ \rightarrow K^+ \pi^- \pi^+)} \propto \frac{T(c \rightarrow \bar{s}ud)}{T(c \rightarrow \bar{d}us)} \propto \tan^2 \theta_c$$

Introduction of 4th quark
Charm decays

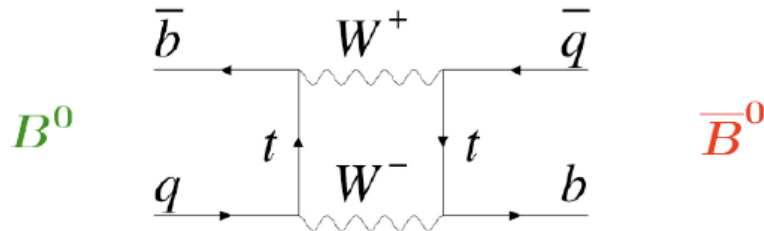
$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix} \equiv \hat{V}_{\text{CKM}} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

CP violating phase \rightarrow 3rd family

Makoto Kobayashi, Toshihide Maskawa

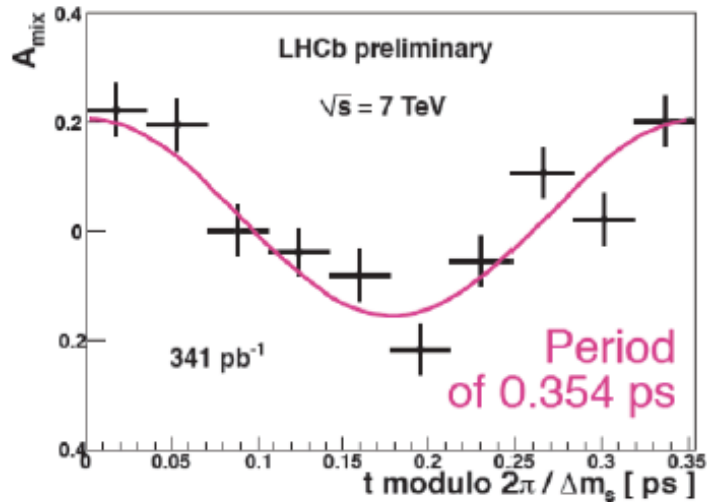


Recent B Physics



Discovered by ARGUS 1987

LHCb: $\delta\rho=300\mu\text{m}$, $\delta t=50\text{fs}$



B_s mixing measured ok with SM
best measurement of Δm_s

$\text{Br}(B_s \rightarrow \mu\mu) < 1.1 \cdot 10^{-8}$ [3.4 * SM]

LHCb+CMS do not confirm CDF excess

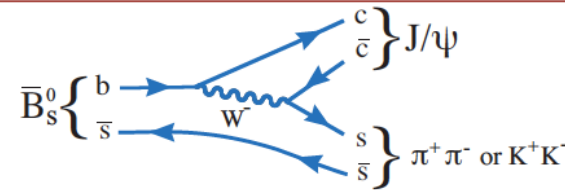
$B \rightarrow K^* \ell \ell$ FB asymmetry: LHCb: SM ok (\neq Belle09)

Like sign di-muon asymmetry

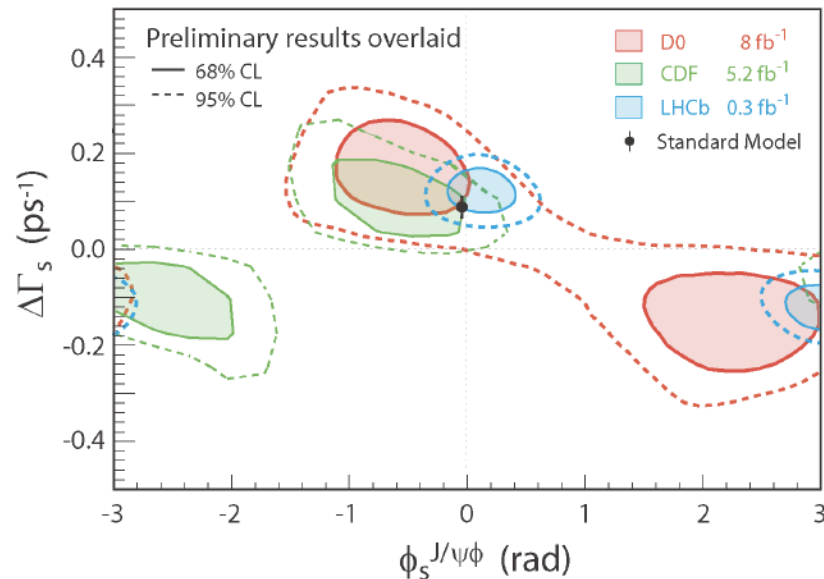
D0: $(-0.787 \pm 0.172 \pm 0.093)\%$ 3.9σ above SM

LHCb: to come

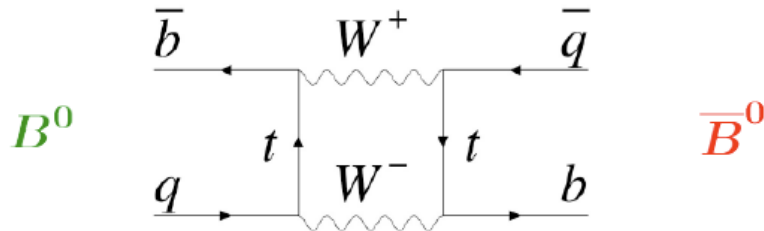
Some tension between $\text{Br}(B \rightarrow \tau\nu)$ and $\sin 2\beta$



CP violating phase Φ in B_s^0 decays

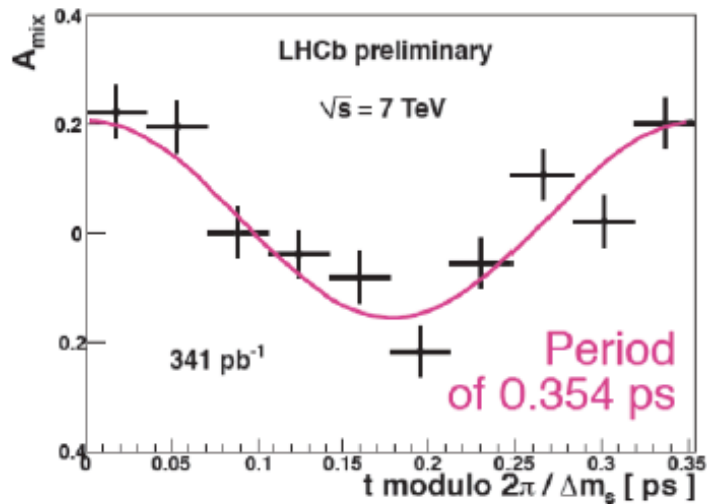


Future B Physics



Discovered by ARGUS 1987

LHCb: $\delta\rho=300\mu\text{m}$, $\delta t=50\text{fs}$



B_s mixing measured ok with SM
 best measurement of Δm_s

NA62

$\text{Br}(B_s \rightarrow \mu\mu) < 1.1 \cdot 10^{-8}$ [$3.4 \cdot \text{SM}$]

LHCb+CMS do not confirm CDF excess

$B \rightarrow K^* \ell \ell$ FB asymmetry: LHCb: SM ok (\div Belle09)

Like sign di-muon asymmetry

D0: $(-0.787 \pm 0.172 \pm 0.093)\%$ 3.9σ above SM

LHCb: to come

Some tension between $\text{Br}(B \rightarrow \tau\nu)$ and $\sin 2\beta$

To be studied (with high precision):

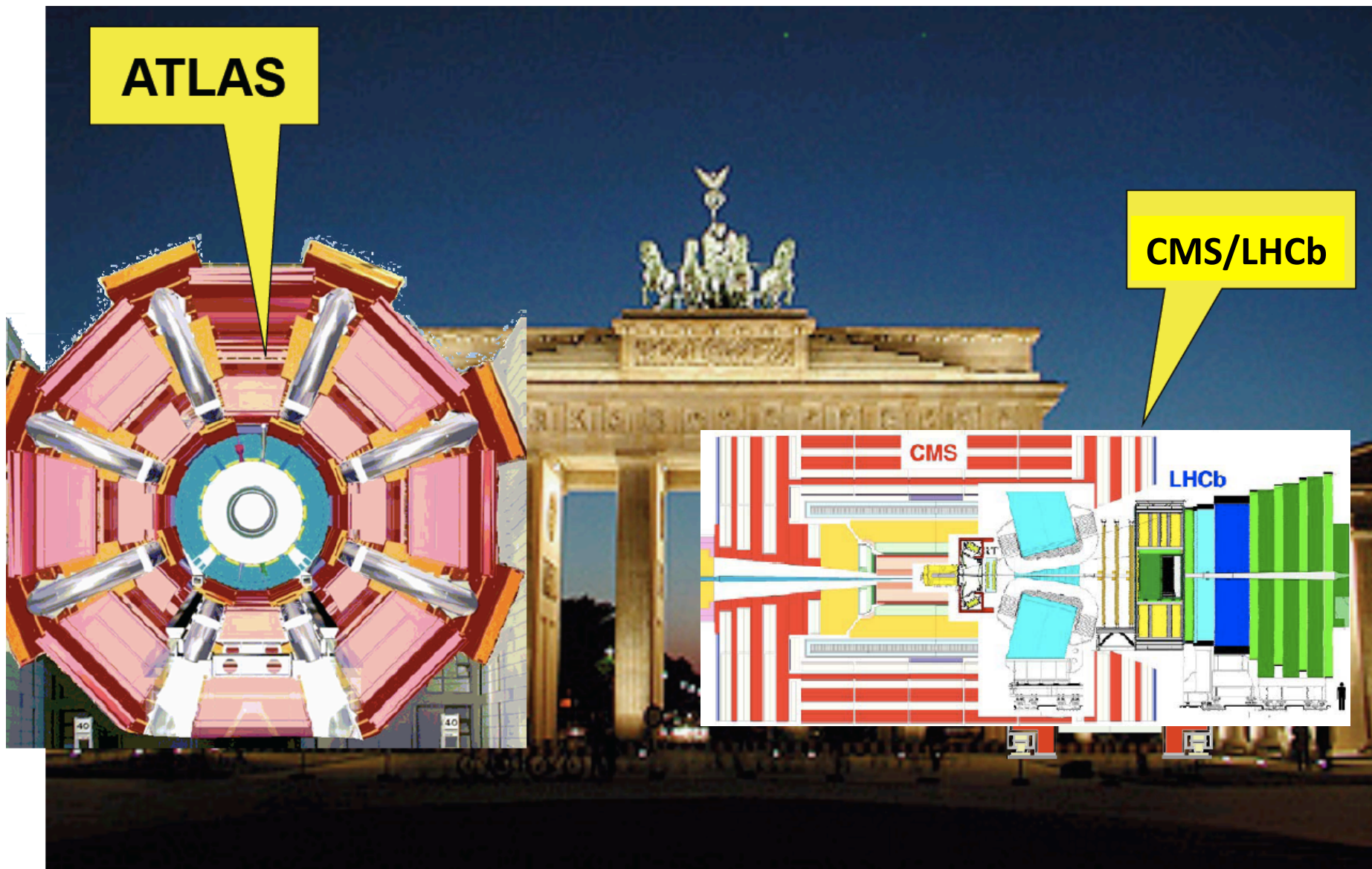
- γ from tree ($B \rightarrow DK, \dots$)
- $|V_{ub}|$ from exclusive semilept. B decays
- $B_{s,d} \rightarrow \mu\mu$
- CPV in B_s mixing
- $B \rightarrow K^* \mu\mu$ (angular analysis)
- $B \rightarrow \tau\nu, \mu\nu$
- $K \rightarrow \pi\nu\nu$
- CPV in D mixing

New B factories
 in Japan and
 near Rome

Projects



Projects



The LHC Upgrades

	2013	2023	2033+
	LHC	HL-LHC	HE-LHC
Collision energy [TeV]	14	14	33
Peak/leveled luminosity [$10^{34} \text{ cm}^{-2}\text{s}^{-1}$]	1.0	7.9/5.0	2.0/2.0
integrated luminosity per year (1900h) [fb^{-1}]	57	250	100
events per crossing	19	150	76
# bunches / beam	2808	2808	1404
bunch population [10^{11}]	1.15	1.7	1.29
Beam current [A]	0.58	0.86	0.32
Luminosity leveling	no	θ_c , V_{crab} or β^*	$\epsilon_{x,y}$
initial transverse normalized emittance [μm]	3.75	3.75	3.75 (x), 1.84 (y)
number of IPs contributing to tune shift	3	3	2
maximum total beam-beam tune shift	0.01	0.01	0.01
IP beta function [m]	0.55	0.14	1.0 (x), 0.43 (y)
full crossing angle [μrad]	285 ($9.5 \sigma_{x,y}$)	0 (509)	175 ($12 \sigma_{x0}$)
dipole field [T]	8.33	8.33	20
dipole coil aperture [mm]	56	56	40-45
stored beam energy [MJ]	362	504	479
SR power per ring [kW]	3.6	5	62.3
longitudinal SR emittance damping time [h]	12.9	12.9	0.98
luminosity lifetime [h]	23	4	13

The LHC Luminosity Upgrade

- New high field insertion **quadrupoles**
- Upgraded **cryo system** for IP1 and IP5
- Upgrade of the intensity in the **Injector Chain (LIU)**
- **Crab Cavities** to take advantage of the small beta*
- **Single Event Upsets**
 - **SC links** to allow power converters to be moved to surface
- **Misc**

S.Myers EPS11 Grenoble

Goals for 25ns bunch crossing

$2 \cdot 10^{11}$ p/bunch

$\beta^* = 0.15$ m

$L = 7 \cdot 10^{34}$ leveled to $5 \cdot 10^{34} \text{ cm}^{-2}\text{s}^{-1}$

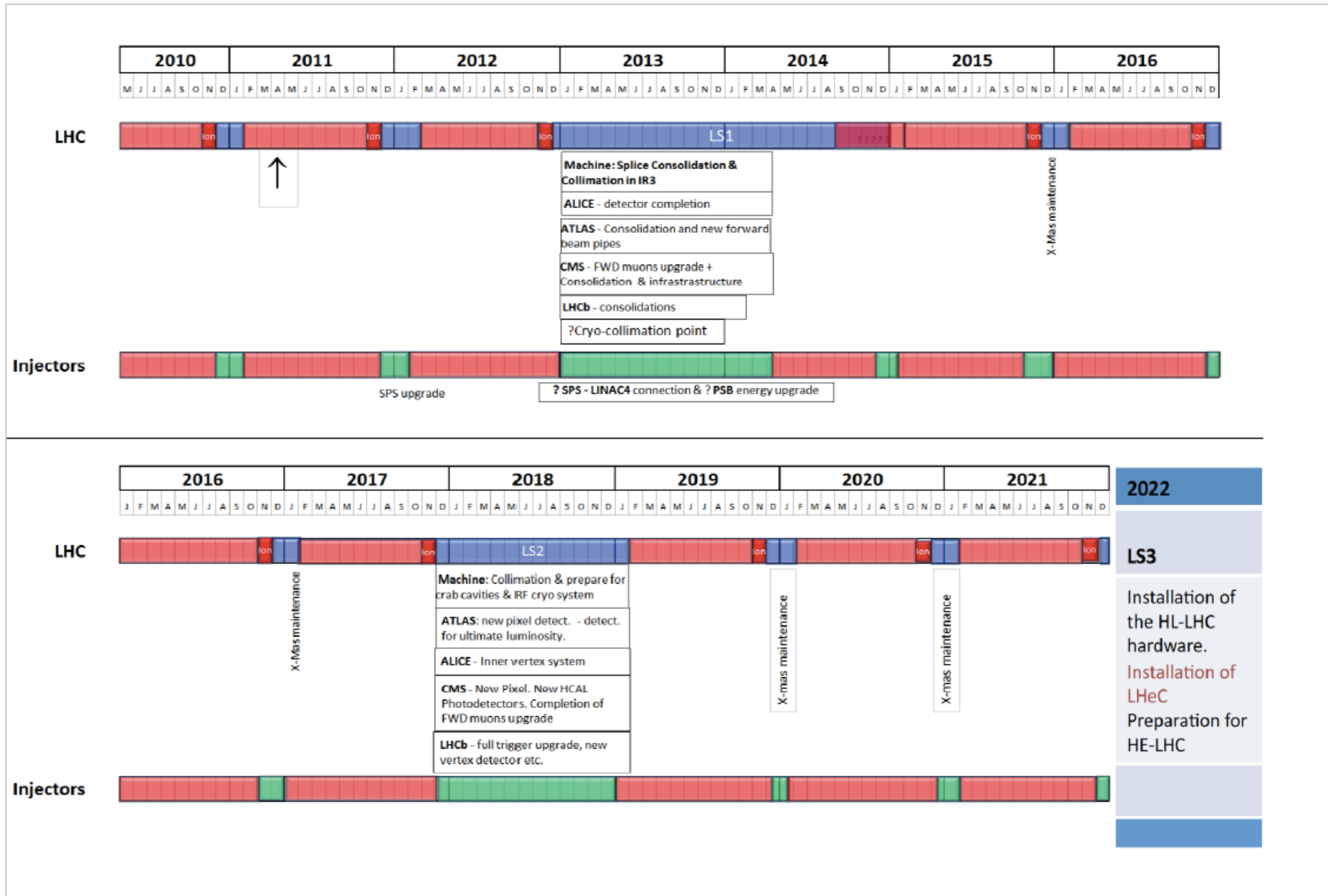
Integrated $L = 3\text{ab}^{-1}$ by 2030+

Corresponding detector upgrades:

ATLAS 150m^2 strips, 10m^2 pixels

Trigger, muons..

The coming LHC decade -DRAFT



The Large Hadron **electron** Collider



$$L = \frac{1}{4\pi} \cdot \frac{N_p}{\epsilon_p} \cdot \frac{1}{\beta^*} \cdot \gamma \cdot \frac{I_e}{e}$$

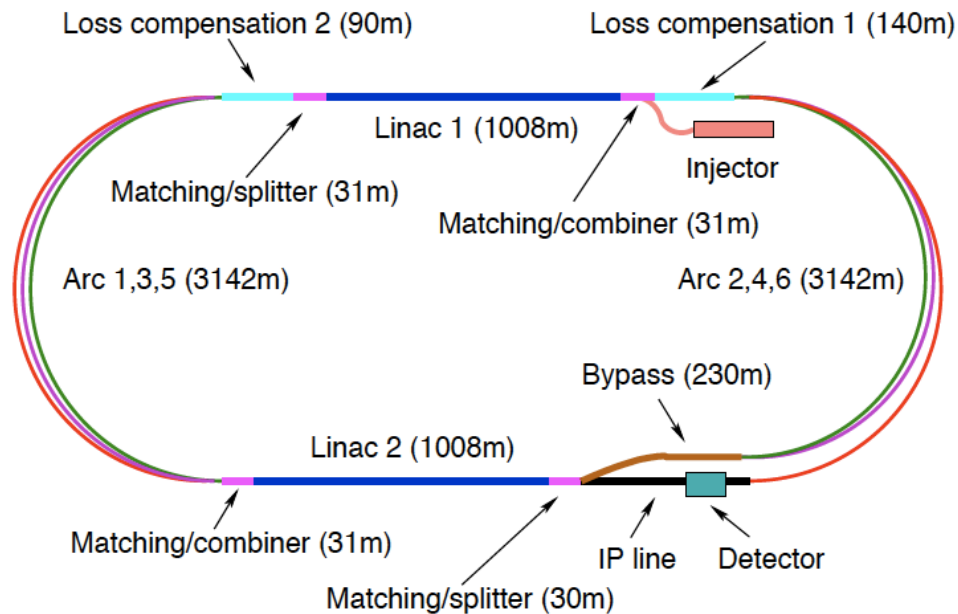
$$N_p = 1.7 \cdot 10^{11}, \epsilon_p = 3.8 \mu\text{m}, \beta^* = 0.2\text{m}, \gamma = 7000/0.94$$

$$L = 8 \cdot 10^{31} \text{cm}^{-2} \text{s}^{-1} \cdot \frac{N_p 10^{-11}}{1.7} \cdot \frac{0.2}{\beta^*/\text{m}} \cdot \frac{I_e/\text{mA}}{1}$$

$$I_e = \text{mA} \frac{P/\text{MW}}{E_e/\text{GeV}}$$

Table 2: Components of the Electron Accelerators

	Ring	Linac
magnets		
beam energy	60 GeV	
number of dipoles	3080	3600
dipole field [T]	0.013 – 0.076	0.046 – 0.264
total nr of quads	866	1588
RF and cryogenics		
number of cavities	112	944
gradient [MV/m]	11.9	20
RF power [MW]	49	39
cavity voltage [MV]	5	21.2
cavity R/Q [Ω]	114	285
cavity Q_0	–	$2.5 \cdot 10^{10}$
cooling power [kW]	5.4@4.2 K	30@2 K



- TeV scale ep collider [60x7000]GeV²
- ep and eA collisions using LHC
- synchronous ep and pp operation
- 100 times HERA luminosity
- 100 MW wall plug power
- Energy recovery allows GW power

Legend:

- CERN existing LHC
- CLIC 500 GeV
- CLIC 3 TeV
- ILC 500 GeV
- LHeC

Potential underground siting

Jura Mountains

IP

Geneva

Lake Geneva

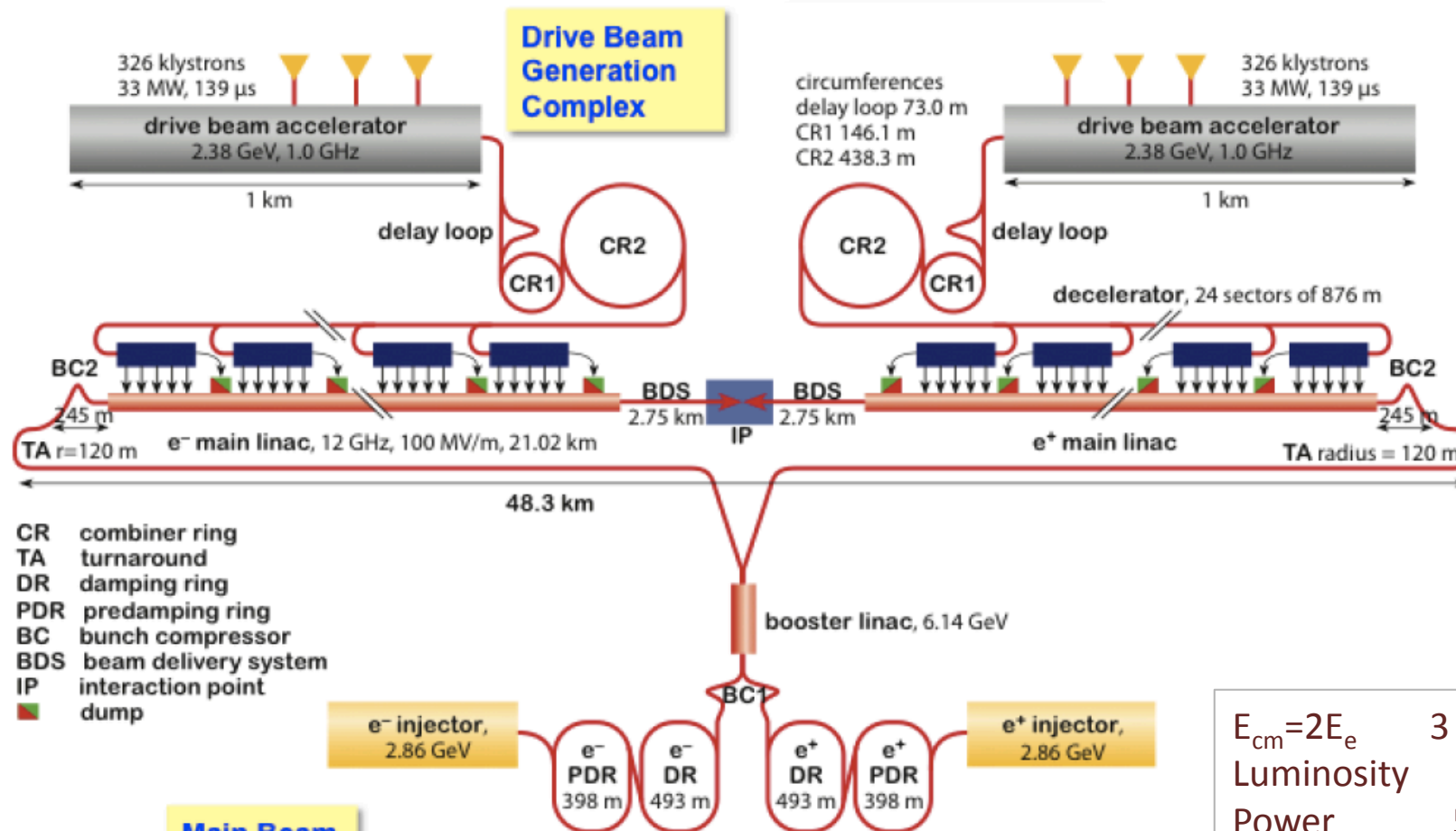
Potential future projects at CERN

©2010 Google

Schematic layouts for several potential future projects are shown on this Google Earth view of the Geneva region around CERN:

- CLIC (Compact Linear Collider) at collision energies of 500GeV and 3 TeV.
- ILC (International Linear Collider) at 500GeV energy
- The Linac-Ring Solution of LHeC (A new electron beam supplied via a 60 GeV

Compact Linear Collider - CLIC



- CR combiner ring
- TA turnaround
- DR damping ring
- PDR predamping ring
- BC bunch compressor
- BDS beam delivery system
- IP interaction point
- dump

Main Beam Generation Complex

booster linac, 6.14 GeV

e⁻ injector, 2.86 GeV

e⁻ PDR 398 m
e⁻ DR 493 m

e⁺ DR 493 m
e⁺ PDR 398 m

e⁺ injector, 2.86 GeV

consider E staging depending on LHC:
0.5 TeV 14km
1-2 TeV 20-34km

$E_{cm} = 2E_e$	3 TeV
Luminosity	$5.9 \cdot 10^{34}$
Power	560 MW
Gradient	100 MV/m
Length	48 km
IP beam size	40/1 nm
Push-pull detectors?	

The International Linear Collider

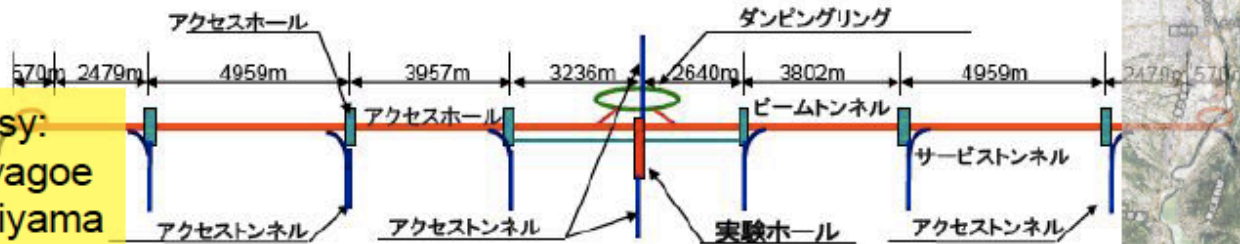


Courtesy: A. Enomoto, M. Miyahara

今年度の検討課題

■ 検討ケースの概要

共通事項: 地下構造物の基本レイアウト

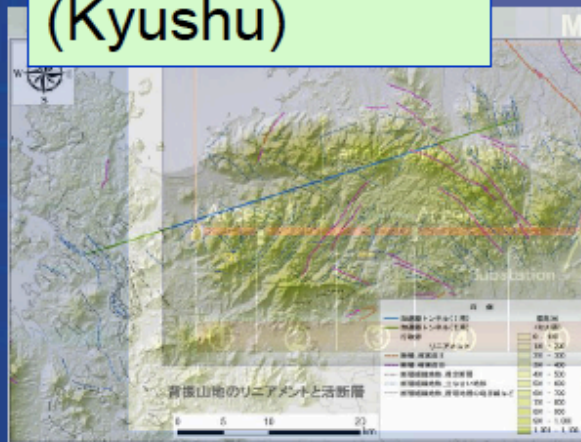


Courtesy:
K. Kawagoe
A. Sugiyama



Northeast
(Tohoku)

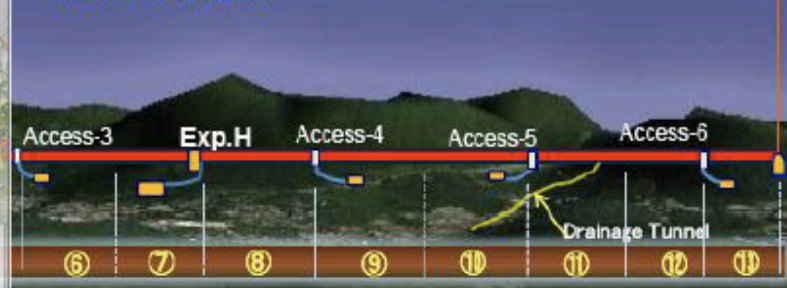
Far west Japan
(Kyushu)



山岳サイトのイメージ/NATM

Main Linac Tunnel, L=31km

Granite Rock Zone



加速器学会・加速器土木

$E_{cm} = 2E_e$	500 GeV
Luminosity	$1.5 \cdot 10^{34}$
Power	215 MW
Gradient	31.5 MV/m
Length	31 km
IP beam size	474/4 nm

Push-pull detectors?

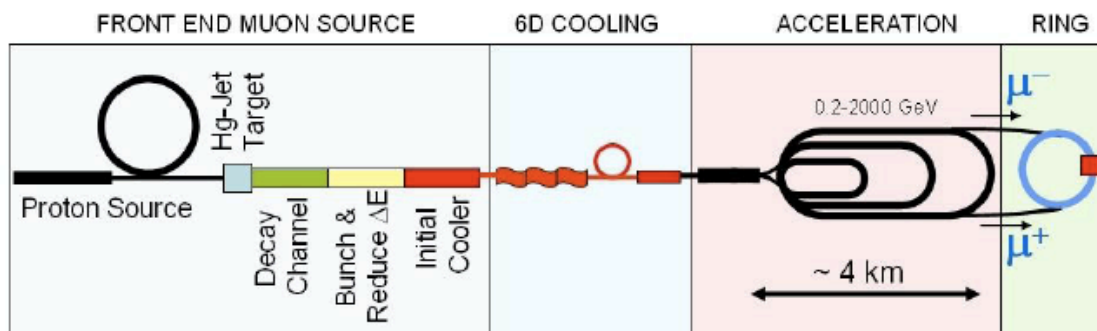
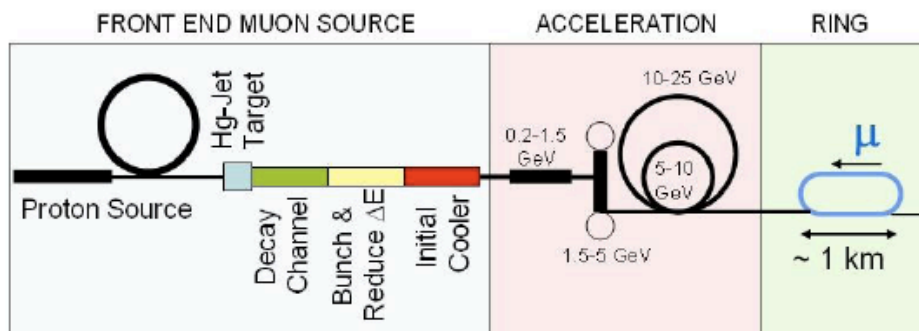
Muon Collider

A Future for Fermilab:

ProjectX, Neutrino Factory, Muon Collider

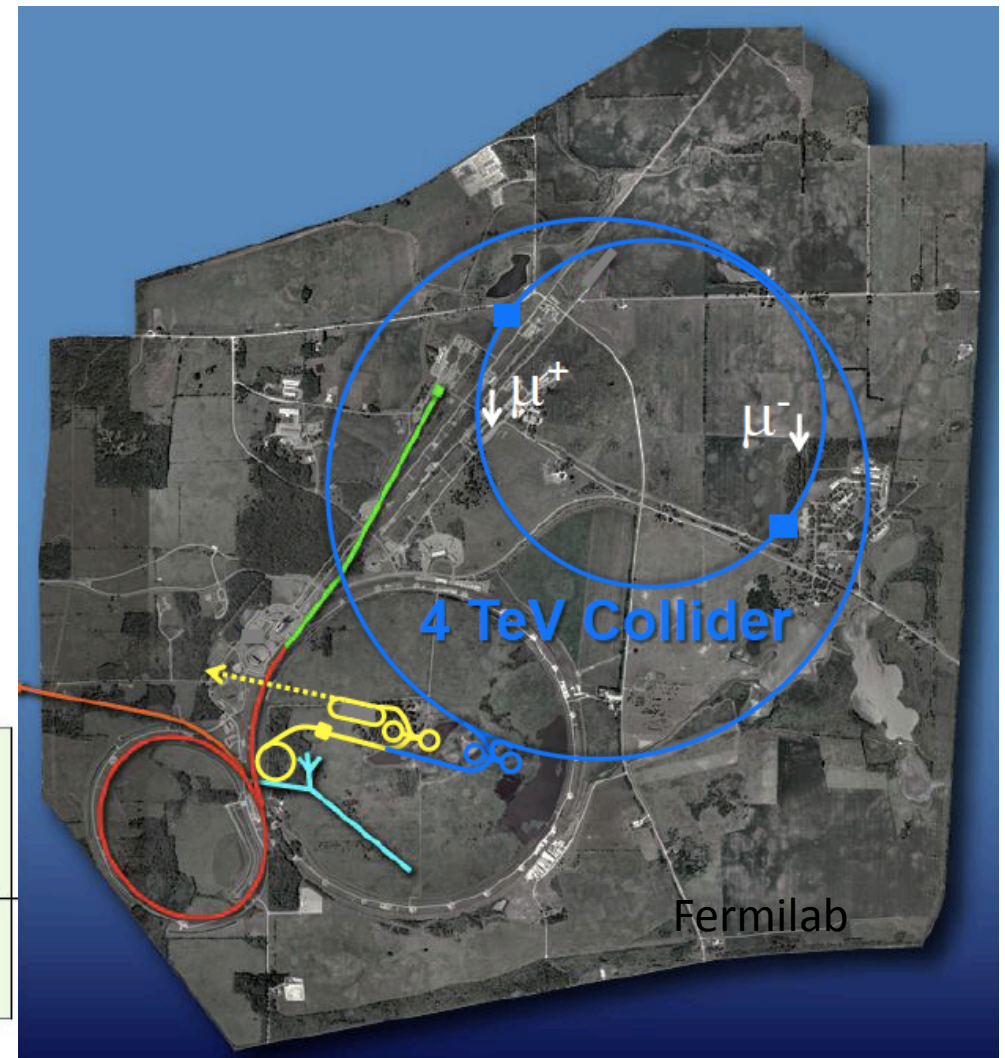
3 Proposals for ν Factories:

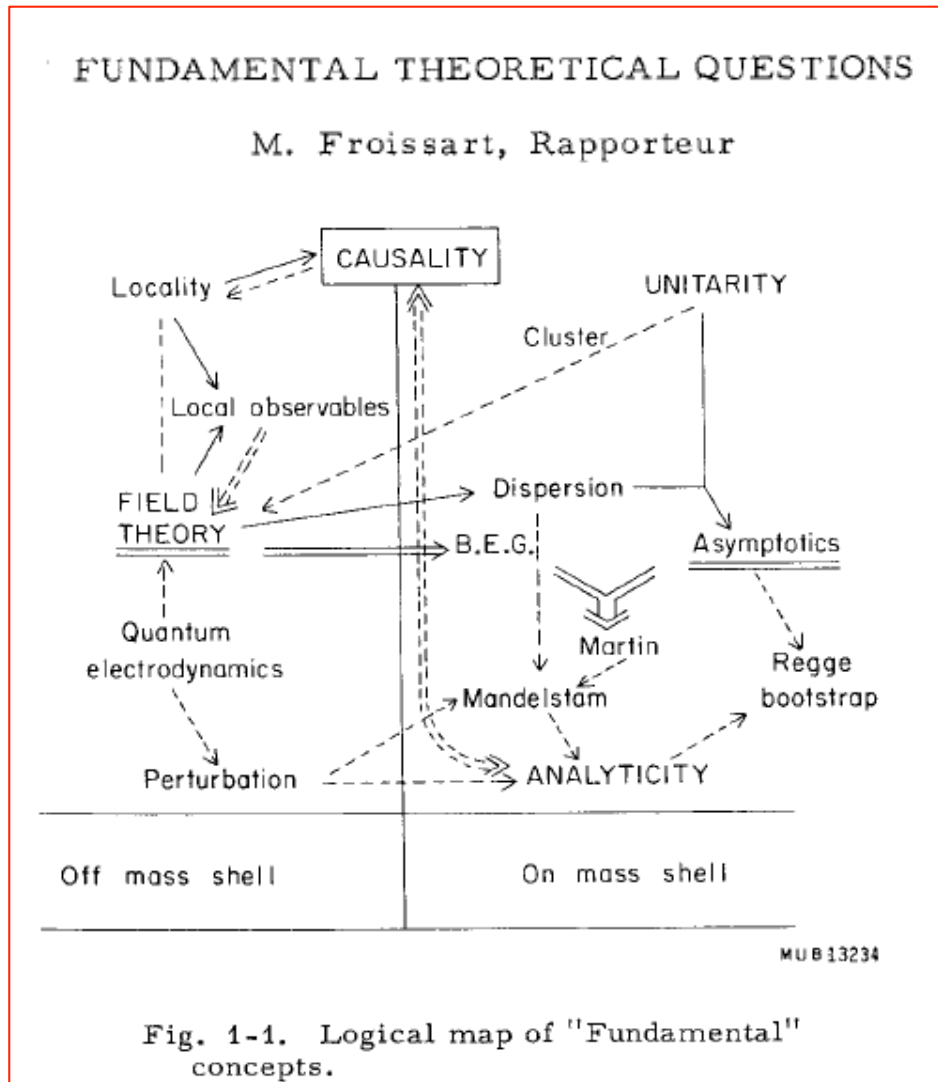
JPARC, CERN, FNAL based. Detector R+D



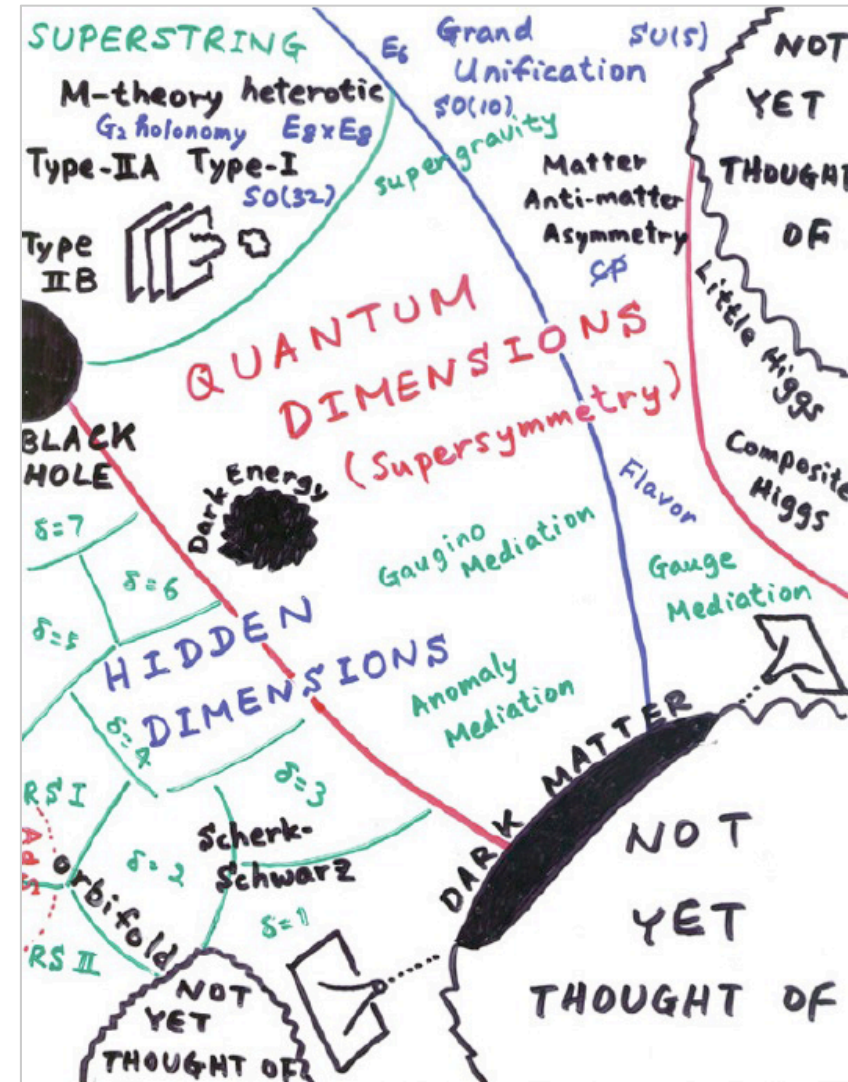
MUON COLLIDER

Circumference = 4.5km
 $L = 3 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$
 $\mu/\text{bunch} = 2 \times 10^{12}$
 $\delta p/p = 0.1\%$





→ Quarks in 1969



→ ?in 2014?



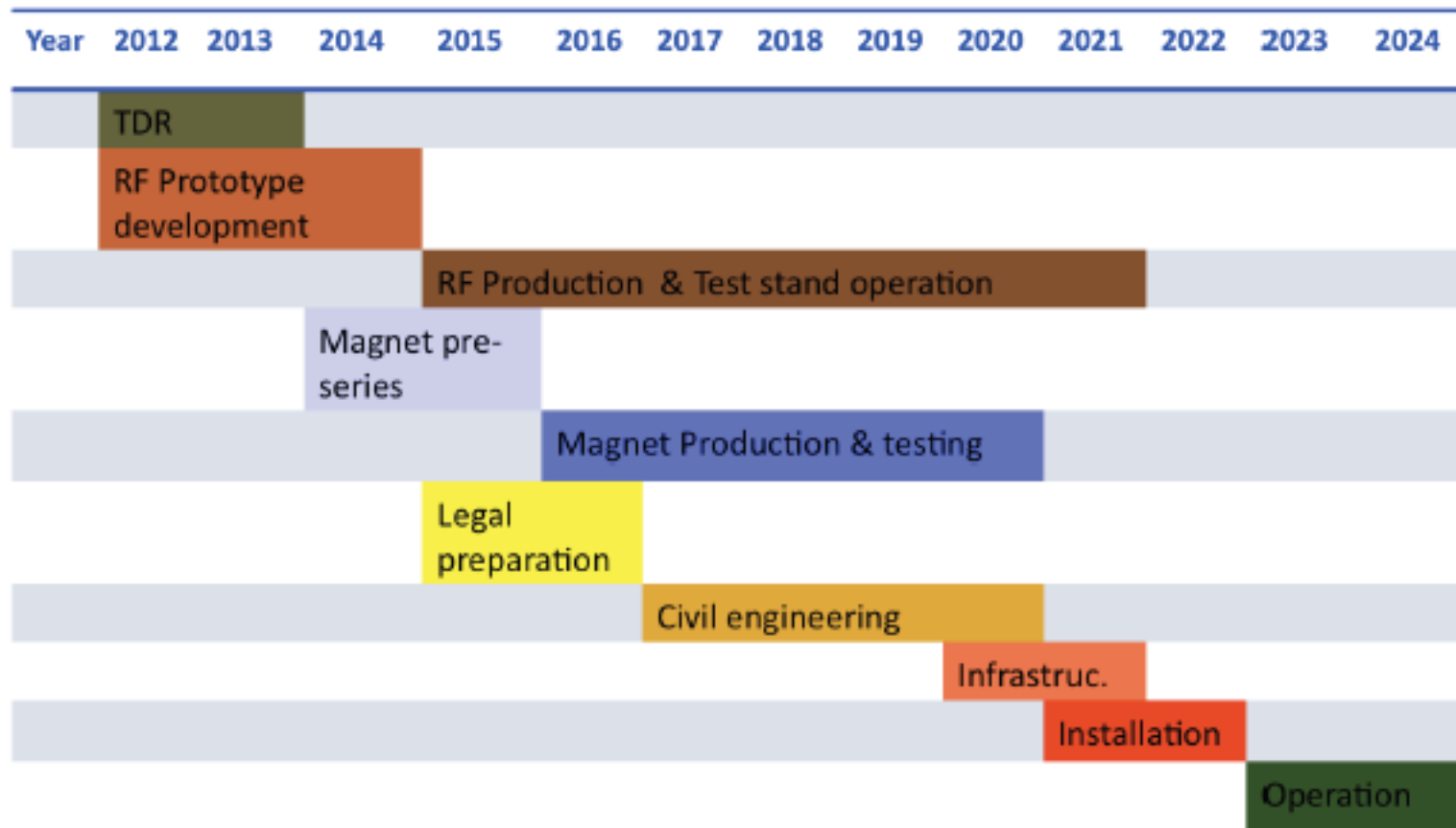
H.Murayama – ICFA11

“We like to see particle physics as driven by experiment...” Burt Richter 2009 at CERN

Liverpool HEP: ATLAS, ATLASupgrade, LHCb, LHCbupgrade, T2K,vFactory, NA62, LHeC + AstroPP

backup

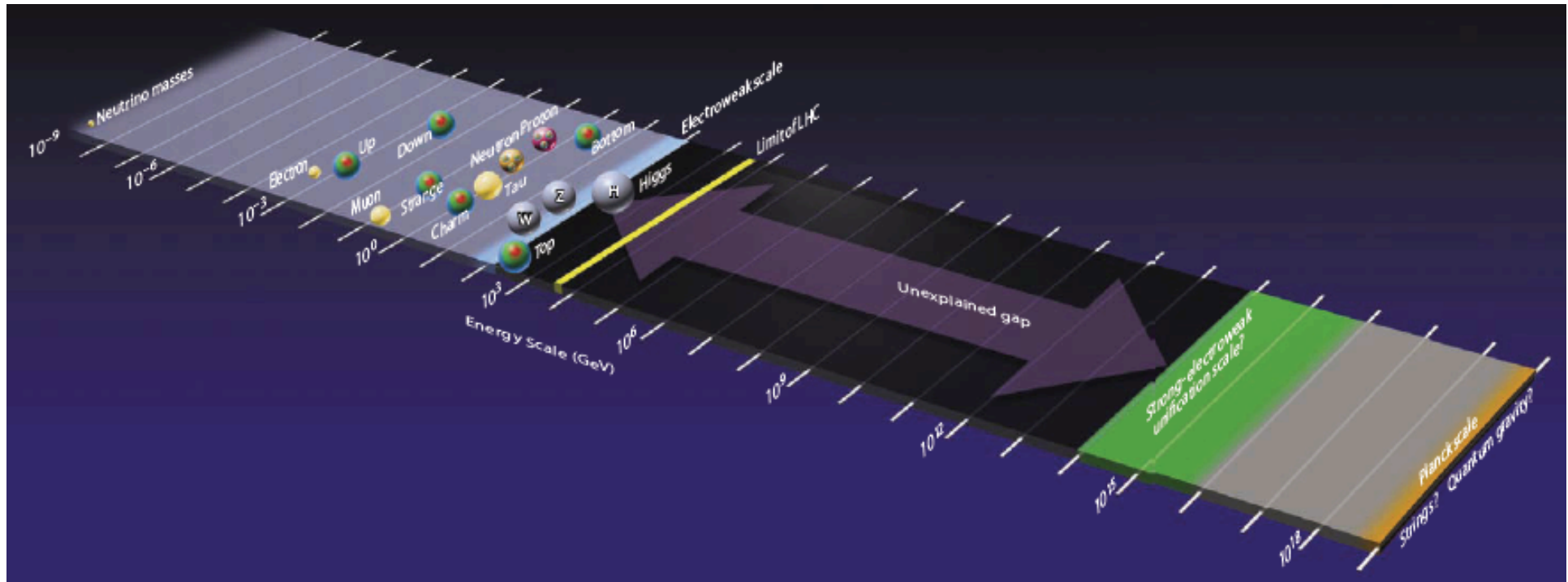
LHeC Tentative Time Schedule



LS3 --- HL LHC →

We base our estimates for the project time line on the experience of other projects, such as (LEP, LHC and LINAC4 at CERN and the European XFEL at DESY and the PSI XFEL)

The energy frontier



C.Quigg

With ingenuitive accelerators of the past century we have used beams of 5 MeV \rightarrow 1 TeV, an increase of six orders of magnitude in 100 years. We have no idea how to achieve a similar gain in this century. Theory is not certain as to what lies above or around the TeV scale. Three decades ago SU(5) predicted 'the desert', but it contradicted experiment

title

title