## Dark photons: bridging the gap between resonant and beam dump searches with LHCb

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UNIVERSITY OF LIVERPOOL Seminar





Dark Photons

## Introduction

- search for dark photons via  $D^{*0} \to D^0 A'$
- still preliminary, so limits will probably change slightly!
- to be released soon (about a month)
- dark matter and dark photons
- dark photon motivation
- current limits
- LHCb detector
- signal rate
- detector response
- search

#### Dark photons from charm mesons at LHCb

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We propose a search for dark photons A' at the LHCs experiment using the charm meson decay  $D^*(2007)^2 \rightarrow D^*A$ . At nominal luminosity,  $D^{-1} \rightarrow D^{-1}$  decays will be produced at about 700 kHz within the LHC5 acceptance, yielding over 5 trillion such decays during Run 3 of the LHC. Replacing the photon with a luniterially-insulad dark photon. LHC5 is then sensitive to dark photons that advantage of the large Lorentz boses of the dark photon, and the excellent vertex resolution of LHCb, yielding a nearly background-free search when the A' decay vertex is significantly displaced from the proton-proton primary vertex. The resonant strategy takes advantage of the large lorentz that decay the proton-proton primary vertex. The resonant strategy takes advantage of the large lorentz instacles are  $D^{+0} - D^{-1} A'$  and the excellent invirant mass resolution of LHCb, yielding a backgroundlimited search that nevertheless corers a significant portion of the A' parameter space. Both search background background to a triggeriser seadout system at LHCb in Run 3, which will photon masses below about 100 MeV. LHCb can explore the entire dark photon parameter space baven existing prompt-A' and background-frame.

## Dark Matter

## Indirect Observation

- $\approx 85\%$  of matter estimated to be dark matter
- large body of experimental evidence, see Phys. Rept. **405** (2005) 279 for a full review
  - spiral galaxy rotation curves
  - strong and weak lensing
  - cosmic microwave background
  - merging clusters and galaxies
  - and many more ...





### MNRAS 249 (1991) 523

Ilten

Dark Photons

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## Dark Matter Types

- hot dark matter
  - light ultra-relativistic particles, e.g. neutrinos
  - not consistent with structure formation
- cold dark matter
  - weakly interacting massive particles, WIMPS
  - possible candidates from SUSY and hidden valley models

$$\Omega_{\chi}h^{2} \propto \frac{1}{\langle \sigma_{a}v \rangle} \implies \langle \sigma_{a}v \rangle \approx 0.9 \text{ pb}$$

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## Dark Sector Portals

• limited number of SM symmetries allow dark matter interactions

type	particles	operator
vector	dark photons	$B_{\mu u}F^{\prime\mu u}$
pseudo-scalar	axions	$F_{\mu\nu}\widetilde{F}^{\mu\nu},\ G_{i\mu\nu}\widetilde{G}_{i}^{\mu\nu},\ \psi^{\dagger}\gamma^{\mu}\gamma^{5}\psi$
scalar	dark Higgs	$H^{\dagger}H$
fermion	sterile neutrinos	LHN

- scalar (Higgs) portal explored via high-energy colliders
- fermion (neutrino) portal via neutrino experiments
- vector and pseudo-scalar (dark photon and axion) portals via high luminosity experiments

## Dark Photons

- minimal model, assume brokenU(1)gauge symmetry in dark sector
- allow mixing between dark and SM hypercharge fields via  $B_{\mu\nu}F'^{\mu\nu}$

$$\mathcal{L} \supset -\frac{1}{4} \widetilde{F}_{\mu\nu} \widetilde{F}^{\mu\nu} - \frac{1}{4} \widetilde{F}'_{\mu\nu} \widetilde{F}'^{\mu\nu} + \frac{\epsilon}{2} \widetilde{F}_{\mu\nu} \widetilde{F}'^{\mu\nu} + \frac{m_{A'}^2}{2} \widetilde{A}'_{\mu} \widetilde{A}'^{\mu} + e J^{\mu}_{\rm EM} \widetilde{A}_{\mu}$$

• transform to eliminate the non-diagonal mixing terms

$$\widetilde{A}'_{\mu} = \frac{1}{1-\epsilon^2} A'_{\mu}, \quad \widetilde{A}_{\mu} = A_{\mu} + \frac{\epsilon}{1-\epsilon^2} A'_{\mu}$$

$$\mathcal{L} \supset -rac{1}{4}F_{\mu
u}F^{\mu
u} - rac{1}{4}F'_{\mu
u}F'^{\mu
u} + rac{m^2_{A'}}{2}A'_{\mu}A'^{\mu} + eJ^{\mu}_{\mathrm{EM}}A_{\mu} + rac{e\epsilon J^{\mu}_{\mathrm{EM}}A'_{\mu}}{2}$$

- mass of the dark photon,  $m_{A'}$  and mixing  $\epsilon$  are free parameters

## Phenomenology



- $\epsilon \approx 10^{-3}$  if heavy lepton with SM and DM charge
- $\epsilon \approx 10^{-5}$  if SM embedded in GUT symmetry
- assume only visible SM decays (invisible also possible)
- search via prompt decays (resonance searches) or displaced decays (usually beam dumps)



## Motivation

## Cosmic Ray Positrons

- PAMELA, Fermi Large Area Telescope, and AMS-02 observe unexpected rise in positron to electron ratio
  - also consistent with complementary measurement from Fermi Gamma-Ray Telescope



Phys. Rev. Lett. 110 (2013) 141102

- possible via modification of cosmic ray propagation model
  - significantly less radiation loss of positrons required
- viable with change in positron production model
  - positrons from proton scattering primarily from cosmic ray origin
- dark photon annihilation also consistent with data

## Lighter Dark Matter

- CDMS observed three events consistent with a WIMP mass between 7 20 GeV
  - compatible with CoGeNT results
  - both CDMS and CoGeNT in tension with XENON100 results
  - SuperCDMS in tension with CoGeNT
  - CoGeNT excess possibly explained via surface background
- possible dark matter annihilation from the galactic center with mass of 10-50 GeV
- results not compatible with Z or H scattering
  - compatible with dark photons



## Self-Interacting Dark Matter

- galaxy structure not perfectly modeled by collisionless dark matter
  - cusps are flatter then expected in data
  - observed sub-halos are not as concentrated as expected
- lag between halo and dark matter halo observed in some merging galaxies
  - non-observation in some sets limits on interactions



ESO/R. Massey

• self-interaction via long range forces can resolve both modeling and lag issues

### Motivation

## Muon Anomalous Magnetic Moment

- the electron and muon anomalous magnetic moments,  $a_e$  and  $a_{\mu}$ , constrain the  $\epsilon^2 m_{A'}$  parameter space
- tension between  $a_{\mu}$  theory and experiment favors a specific parameter space
- further work needed on both experiment and theory, but well in hand here at Liverpool!



Phys. Rev. D 86 (2012) 010001

# Prompt Limits

## BaBar

- BaBar detector at SLAC •
  - Phys. Rev. Lett. **113** (2014) 20, 201801
  - $514 \text{ fb}^{-1}$  of data from 1999 - 2008
  - $e^+$  beam (3.1 GeV),  $e^$ beam (9 GeV)
  - $e^+e^- \rightarrow \gamma A'$  production
  - search for  $A' \to e^+e^-$  and  $A' \rightarrow \mu^+ \mu^-$

MC e'e → γ R



Entries / 0.1 GeV

10 Data/MC

10 m<sub>ee</sub> (GeV)

10-3

NA48/2

## NA48/2

- NA48/2 experiment at CERN
  - Phys. Lett. B 746 (2015) 178
  - 2003 2004 run
  - $K^{\pm}$  beam (60 GeV from 400 GeV proton beam)
  - use  $K^{\pm} \to \pi^{\pm} \pi^0$  and  $K^{\pm} \to \pi^0 \mu^{\pm} \nu$  decays
  - search for  $A' \to e^+ e^-$  from  $\pi^0 \to \gamma A'$
  - irreducible  $\pi^0 \to \gamma e^+ e^-$



 $10^{0}$ 

10<sup>-1</sup>

### Prompt Limits

## APEX and MAMI

- A1 detector at the Mainz Microtron (MAMI)
  - Phys. Rev. Lett. 106 (2011) 251802, Phys. Rev. Lett. 112 (2014) 22, 221802
  - 2011 and 2014 runs
  - electron beam (180 855 MeV), heavy nuclei target
  - search for  $\gamma/A' \rightarrow e^+e^-$
- APEX (A' experiment) at Jefferson Lab
  - Phys. Rev. Lett. 107 (2011) 191804
  - 2011 test run
  - electron beam (2.3 GeV), heavy nuclei target
  - search for  $\gamma \to e^+ e^-$
- future runs expected for both experiments



# Displaced Limits

## E137, E141, and E774

- Phys. Rev. D 80 (2009) 075018, A' interpretation of beam-dump data
- E137 experiment at SLAC
  - Phys. Rev. D **38** (1988) 3375
  - electron beam (20 GeV), aluminum target
  - 200 m through earth, 200 m through air
  - 3 m detector
- E141 experiment at SLAC
  - Phys. Rev. Lett. **59** (1987) 755
  - electron beam (9 GeV), tungsten target (10 cm)
  - 35 m through vacuum
- E774 experiment at FermiLab
  - Phys. Rev. Lett. **67** (1991) 2942
  - electron beam (275 GeV), tungsten target (30 cm, 28 radiation lengths)
  - 7.25 m through vacuum



## Orsay and U70

- Orsay
  - Phys. Lett. B **229** (1989) 150
  - electron beam (1.7 GeV), lead target
- U70 experiment at IHEP Serpukhov
  - interpretation in Phys. Lett. B 701 (2011) 155
  - Z. Phys. C **51** (1991) 341
  - proton beam (70 GeV), iron target (140 cm)
  - 64 m through vacuum
  - A' from  $\pi^0 \to \gamma A'$
  - $A' \rightarrow e^+ e^-$



## Charm and LSND

- can use results from neutrino experiments to place limits
- CHARM experiment at CERN
  - interpretation in Phys. Lett. B 713 (2012) 244
  - Phys. Lett. B **166** (1986) 473
  - proton beam (400 GeV), copper target
  - use  $\eta \to \gamma A'$  and  $\eta' \to \gamma A'$
  - $A' \rightarrow e^+ e^-$
- LSND experiment at Los Alamos
  - Phys. Rev. C 58 (1998) 2489
  - proton beam (800 MeV), water target (30 cm)
  - A' from  $\pi^0 \to \gamma A'$
  - $A' \rightarrow e^+ e^-$



## Supernovae

- can also use results from supernovae to place limits
  - supernovae cooling mechanism well modeled
  - any cooling faster than expected indicative of dark matter
- Phys. Rev. D 89 (2014) 10, 105015
- dark fermions  $\psi'$  and scalars  $\phi$ can be produced from electrons, positrons, and nucleons

$$\begin{array}{ll} e^+e^- \to \bar{\psi}'\psi', & NN \to NN \phi \\ e^+e^- \to \phi'^\dagger \phi', & NN \to NN \phi \end{array}$$





## Collaboration

- Large Hadron Collider beauty experiment on the LHC
- letter of intent 1995, CERN-LHCC-95-5
- a forward collider detector dedicated to the study of CP violation and other are phenomena in the decays of Beauty particles
- 1111 members, 69 institutes (including Liverpool and MIT), 17 countries





### Detector



- fully instrumented between  $2 < \eta < 5$
- momentum resolution between 0.4% at 5 GeV to 0.6% at 100 GeV
- impact parameter resolution of  $13-20\;\mu\mathrm{m}$  for tracks
- secondary vertex precision of 0.01 0.05(0.1 0.3) mm in xy(z)



## Current Datasets

- current data testament to incredible LHCb performance
- 1 fb<sup>-1</sup> pp collisions at  $\sqrt{s} = 7$  TeV (2011)
- 2 fb<sup>-1</sup> pp collisions at  $\sqrt{s} = 8$  TeV (2012)
- 1.1 nb<sup>-1</sup> *p*Pb collisions at  $\sqrt{s} = 5$  TeV (2013)
- $0.5 \text{ nb}^{-1} \text{ Pb}p$  collisions at  $\sqrt{s} = 5 \text{ TeV} (2013)$
- excellent luminosity uncertainty, JINST 9 (2014) 12, P12005
  - 1.71% for 7 TeV dataset
  - 1.16% for 8 TeV dataset



#### Dark Photons

## Future Datasets

• projected luminosity (see V. Vagnoni (2015) HL-LHC)

LHC era			HL-LHC era		
Run 1(a) 2011	Run 1(b) 2012	Run 2 2015 - 2018	Run 3 2020 - 2022	Run 4 2025 - 2028	Run 5 2030 - ?
$1 \text{ fb}^{-1}$	$2 \text{ fb}^{-1}$	$5 \text{ fb}^{-1}$	$15 {\rm ~fb^{-1}}$	$23 { m  fb^{-1}}$	$54 { m  fb^{-1}}$

- LHCb upgrade during LS 2
  - LHCb-PUB-2014-040
  - replacement of ring imaging Cherenkov detectors
  - replacement of tracking detectors
  - huge effort here at Liverpool on VELO upgrade
  - full software trigger, see LHCB-TDR-016
    - currently limited by hardware readout at 1 MHz
    - upgrade will read out entire detector at 40 MHz
- side-note, interesting EWK opportunities outlined in LHCb-TALK-2015-113

# Signal Rate

### Signal Rate

## Production

- generate soft QCD generates with PYTHIA 8.2
  - non-diffractive, elastic, single-diffractive, double-diffractive, central-diffractive
  - SoftQCD:all = on
  - Monash tune
  - check against LHCb tune (minimal difference)
- perform hadron decays with EVTGEN
- final state radiation in decays with PHOTOS
- require  $D^{*0} \to D^0 X$  decay with  $X \in \gamma, \pi^0$ 
  - fiducial requirement of  $p_{\rm T}(D^0)>1$  GeV,  $2<\eta(D^0)<5$

$$\sigma(pp \to D^{*0} \to D^0 X) = 0.95 \text{ mb}$$
  
$$N_{\text{run } 3}(D^{*0} \to D^0 \gamma) = 5.4 \times 10^{12}$$

• cross-sections validated against LHCb measurements

## $D^{*0}$ Decays

- $D^{*0}$  ideal candidate for search
  - copious production within LHCb
  - $I(J^P) = \frac{1}{2}(1^-)$  allows needed decays
  - mass of  $2006.96 \pm 0.1$  MeV provides large range of  $m_{A'}$ ,  $\Delta m_D = 142$  MeV
  - width of 2.1 MeV provides prompt decays

$$\mathcal{B}(D^{*0} \to D^0 \pi^0) = (61.9 \pm 2.9), \ \mathcal{B}(D^{*0} \to D^0 \gamma) = (38.1 \pm 2.9)$$

• use operator analysis to calculate  $\mathcal{B}(D^{*0} \to D^0 e^+ e^-)$  and  $\mathcal{B}(D^{*0} \to D^0 A')$ 

$$\begin{aligned} \frac{\Gamma(D^{*0} \to D^0 e^+ e^-)}{\Gamma(D^{*0} \to D^0 \gamma)} &= 6.4 \times 10^{-3} \\ \frac{\Gamma(D^{*0} \to D^0 A')}{\Gamma(D^{*0} \to D^0 \gamma))} &= \epsilon^2 \left(1 - \frac{m_{A'}^2}{\Delta m_D^2}\right)^{3/2} \end{aligned}$$

## Rare $\pi^0$ Decays

- the rare decay  $\pi^0 \to \gamma A'$  (used by NA48/2, U70, and LSND) also contributes
- use effective  $\mathcal{L}$  to calculate the amplitude

$$\begin{aligned} |\mathcal{M}_{\pi^0 \to \gamma e^+ e^-}|^2 &= \frac{4\alpha_{\rm EM}^3}{\pi f_{\pi}^2 m_{\gamma e^-}^2} \bigg( m_{\pi^0}^4 + 2m_{\gamma e^-}^4 + m_{e^+ e^-}^4 \\ &+ 2m_{\gamma e^-}^2 m_{e^+ e^-}^2 - 2m_{\pi^0}^2 (m_{\gamma e^-}^2 + m_{e^+ e^-}^2) \bigg) \end{aligned}$$

•  $\Gamma(\pi^0 \to \gamma e^+ e^-)$  verified against experimental value  $\frac{\Gamma(\pi^0 \to \gamma e^+ e^-)}{\Gamma(\pi^0 \to \gamma \gamma)} = 0.012$   $\frac{\Gamma(\pi^0 \to \gamma A')}{\Gamma(\pi^0 \to \gamma \gamma)} = 2\epsilon^2 \left(\frac{m_\pi^2 - m_{A'}^2}{m_\pi^2}\right)^3$ 

## Dark Photon Decays

- assume only  $e^+e^-$  final state
  - invisible decays to dark sector modify limits
  - $\mu^+\mu^-$  and heavier decays suppressed by  $\Delta m_D$

$$\Gamma_{A'} = \frac{\epsilon^2 \alpha_{\rm EM}}{3} m_{A'} \left( 1 + 2 \frac{m_e^2}{m_{A'}^2} \right) \sqrt{1 - 4 \frac{m_e^2}{m_{A'}^2}}$$

• mean flight distance can be written in terms of boost

$$\begin{split} \ell_{A'} \simeq & 16 \, \mathrm{mm} \, \left( \frac{\gamma_{\mathrm{boost}}}{10^2} \right) \left( \frac{10^{-8}}{\epsilon^2} \right) \\ & \times \left( \frac{50 \, \, \mathrm{MeV}}{m_{A'}} \right) \end{split}$$



#### Signal Rate

## Dark Photon Flight



## Detector Response

## Track Types

- use simplified model to estimate detector response
- consider simplified LONG tracks and UP tracks
  - $2 < \eta < 5$  and p > 3 is LONG
  - $2 < \eta < 5$  and p > 1 and not LONG is UP





- momentum resolution from IJMPA **30** (2015) 1530022
- angular resolution determined from  $m_J/\psi$  resolution and multiple scattering

$$\left(\frac{\sigma_p}{p}\right)^2 \approx 2\left(\frac{\sigma_{m_{J/\psi}}}{m_{J/\psi}}\right)^2 - 2\left(\frac{p\sigma_\alpha}{m_{J/\psi}}\right)^2$$

## $D^0$ Reconstruction

- consider two types of  $D^0$  reconstruction
  - *F*-type: fully reconstructed, all children are charged and reconstructed with least two LONG tracks
  - *P*-type: pseudo-fully reconstruct, at least two LONG tracks (provides flight direction), mass constraint

decay	${\mathcal B}$	$\mathcal{B}  imes \operatorname{eff}_D^{\mathrm{F}}$	$\mathcal{B} \times \mathrm{eff}_D^\mathrm{P}$
$D^0 \rightarrow \{K^-\pi^+, KK, \pi\pi\}$	4.4%	2.5%	_
$D^0 \rightarrow \{K^- 3\pi, 2K2\pi, 4\pi\}$	9.1%	4.5%	1.0%
$D^0 \rightarrow K \ell( u)$	6.8%	_	2.0%
$D^0 \to K \pi(\pi^0)_{[0,m_{\kappa^0}]}$	22.0%	_	6.6%
$D^0 \rightarrow KK(K^0)_{[all]}$	1.5%	_	0.5%
$D^0 \to K3\pi(\pi^0)_{[0,m_{K^0}]}$	8.5%	_	1.4%
total		7.0%	11.5%

## Dark Photon Reconstruction

- tracks from  $A' \rightarrow e^+ e^-$  are  $\approx 60\%$  UP and 40% LONG
- electrons should be well identified by RICH
- bremsstrahlung and multiple scattering models implemented in fast simulation
  - uses LHCb upgrade material budget
  - low mass tail from bremsstrahlung
- can apply simple mass correction

$$m_{e^+e^-}^{\rm corr} = m_{e^+e^-}^{\rm reco} \left(2 - \frac{\Delta m_D^{\rm reco}}{\Delta m_D}\right)$$

• alternatively, require full energy-momentum conservation



## $D^{*0}$ Reconstruction

- *F*-type or *P*-type  $D^0$
- reconstructed A' candidate
- require mass difference consistent with  $D^{*0}$

$$\Delta m_D^{\text{reco}} = m_{\text{reco}} (D^0 e^+ e^-) - m_{\text{reco}} (D^0)$$
$$-50 \text{MeV} < \Delta m_D^{\text{reco}} - \Delta m_D < 20 \text{MeV}$$

- mass requirement efficiency  $\approx 85\%$
- highly suppresses production from  $D^{*0} \to D^0 \pi^0$

• 
$$\pi^0 \to \gamma e^+ e^- \text{ or } \pi^0 \to A' e^+ e^-$$

# Search

### Search

## Conversion Veto

- large background from conversions in the VELO
- effectively three types of conversions
  - type 1: conversion in the VELO module
  - type 2: conversion in the RF foil
  - type 3: conversion in the RF foil tip



- type 1, require  $e^+$  and  $e^-$  have hit in the first module encountered with separation of one pixel
- type 2, require at least one pixel separation between  $e^+$  and  $e^-$  hits
- type 3, require hits in same module as tip, if intersected

## **Prompt Selection**

- F-type  $D^0$
- effective conversion veto
  - $2.6 < \eta_{A'} < 5.0$
  - $\alpha_{e^+e^-} > 8 \text{ mrad if}$  $\ell_{\mathrm{T}} < 5.1 \text{ mm}$
  - $\alpha_{e^+e^-} > 25 \text{ mrad if}$ 5.1 <  $\ell_{\rm T} < 16.6 \text{ mm}$



#### Search

## Prompt Reach

- primary background from  $D^{*0} \to D^0 e^+ e^-$
- set 95% confidence limits for  $S/\sqrt{B}>2$

$$\begin{split} \frac{S}{\sqrt{B}} &= \frac{\Gamma(D^{*0} \to D^0 A')}{\sqrt{\Gamma(D^{*0} \to D^0 \gamma) \Delta \Gamma}} \\ &\times \sqrt{N \left( D^{*0} \to D^0 \gamma \right) \operatorname{eff}_{\Delta m_D} \operatorname{eff}_D^{\mathrm{F}} \operatorname{eff}_{A'}(m_{A'}, \epsilon^2)} \end{split}$$

$$\Delta \Gamma \equiv \int_{m_{A'}-\Delta m_{A'}}^{m_{A'}+\Delta m_{A'}} \mathrm{d}m_{e^+e^-} \frac{\mathrm{d}\Gamma(D^{*0} \to D^0 e^+e^-)}{\mathrm{d}m_{e^+e^-}}$$
$$\Delta m_{A'} = 2\sigma(m_{e^+e^-})$$

## **Displaced Selection**

- final selection still under consideration
- F or P-type  $D^0$
- A' decay vertex greater than  $7\sigma$  from pp collision
  - roughly equivalent to  $\ell_{\rm T} > 0.1~{\rm mm}$
- effective conversion veto
  - $2.6 < \eta_{A'} < 5.0$
  - $\alpha_{e^+e^-} > 8 \text{ mrad if}$  $\ell_{\mathrm{T}} < 5.1 \text{ mm}$
  - $\alpha_{e^+e^-} > 25 \text{ mrad if}$  $5.1 < \ell_{\mathrm{T}} < 16.6 \text{ mm}$



#### Search

## Displaced Reach

- assume  $\mathcal{O}(100)$  background events over entire region
  - various backgrounds assessed with large PYTHIA sample to obtain estimates
  - set 95% confidence level for  $S\geq 5$

$$S(m_{A'}, \epsilon^2) = N(\Gamma(D^{*0} \to D^0 \gamma) \frac{\Gamma(D^{*0} \to D^0 A')}{\Gamma(D^{*0} \to D^0 \gamma)} \operatorname{eff}_{\Delta m_D} \times \left(\operatorname{eff}_D^{\mathrm{F}} + \operatorname{eff}_D^{\mathrm{P}}\right) \operatorname{eff}_{A'}(m_{A'}, \epsilon^2) \\ \simeq 85 \left(\frac{\epsilon^2}{10^{-10}}\right) \left(1 - \frac{m_{A'}^2}{\Delta m_D^2}\right)^{3/2} \operatorname{eff}_{A'}(m_{A'}, \epsilon^2)$$

# Conclusions

## LHCb Limits



## Future Experiments

### • MESA

- 2017+, Mainz Energy-Recovering Superconducting Accelerator
- Phys. Rev. D 88 (2013) 015032
- APEX
  - 2018, Jefferson Lab
  - JHEP **1102** (2011) 009
- DarkLight
  - 2016, Jefferson Lab
  - arXiv:1307.4432

- VEPP-3
  - Budker, Russia
  - arXiv:1207.5089 [hep-ex]
- Mu3e
  - 2018+, Paul Scherrer Institute, Switzerland
  - JHEP **1501** (2015) 113
- HPS
  - 2016, Jefferson Lab
  - arXiv:1310.2060 [physics.ins-det]
- SHiP
  - 2023, CERN
  - arXiv:1504.04855 [hep-ph]

#### Conclusions

## Future Limits



## Summary

- dark photons are well motivated within the dark matter sector
- a significant effort is being invested through a variety of experiments
- LHCb should be able to unify the coverage gap between 10 100 MeV, something not possible even with current planned experiments
- further reach can be obtained via additional channels, e.g.  $D^0 \to K^{*0} A'$
- additional experimental techniques, *e.g.* bremsstrahlung recovery, improved mass constraint, *etc.* can improve reach
- can begin to validate methods on actual data during Run 2

## Thank you!