Magnetic Monopoles Theory and Experiment

Arttu Rajantie HEP Seminar, University of Liverpool 6 November 2013

Maxwell Equations

 $ec
abla \cdot ec E$ $ho_{
m E}$ $ho_{
m M}$ $\vec{\nabla} \cdot \vec{B}$ $\partial \vec{B}$ $\vec{\nabla} \times \vec{E}$ \vec{j}_{M} ∂t $\partial \vec{E}$ $\vec{\nabla} \times \vec{B}$ $ec{j}_{\mathrm{E}}$ ∂t

Duality $\vec{E} \leftrightarrow \vec{B}$

Electromagnetic Potentials

$$\vec{E} = -\vec{\nabla}\phi - \frac{\partial \vec{A}}{\partial t}$$
, $\vec{B} = \vec{\nabla} \times \vec{A}$

• Quantum Mechanics: Complex phase of the wave function couples to \vec{A} $i\hbar \frac{\partial}{\partial t}\psi = -\frac{\left(\hbar \vec{\nabla} + ie\vec{A}\right)^2}{2m}\psi + e\phi\psi$

Sourceless magnetic field: No monopoles? $\vec{\nabla} \cdot \vec{B} = \vec{\nabla} \cdot (\vec{\nabla} \times \vec{A}) = 0$

Dirac Monopole (1931)

• Vector potential $\vec{A}(\vec{r}) = \frac{g}{4\pi |\vec{r}|} \frac{\vec{r} \times \vec{n}}{|\vec{r}| - \vec{r} \cdot \vec{n}}$

Singularity along \vec{n} (Dirac string):

 $\circ~$ Carries magnetic flux $\Phi=g$

Classically observable: Induces current in wire loop

Imperial College

London

Dirac Monopole (1931)

• Vector potential $\vec{A}(\vec{r}) = \frac{g}{4\pi |\vec{r}|} \frac{\vec{r} \times \vec{n}}{|\vec{r}| - \vec{r} \cdot \vec{n}}$

Singularity along \vec{n} (Dirac string):

- QM: Complex phase defined modulo 2π
- String unobservable if $g = g_0 = 2\pi/e$

Dirac Monopole (1931)



- If all electric charges satisfy the Dirac quantisation condition $e \in \frac{2\pi}{g}\mathbb{Z}$, string not detectable: Pointlike monopole
- Explanation of charge quantisation?
- Conversely g must be integer multiple of $g_0 = \frac{2\pi}{e_0}$

Dyons



 \blacktriangleright Particles with electric charge e and magnetic charge g

Any two particles have to satisfy the quantisation condition (Schwinger 1966):

 $e_1g_2 - e_2g_1 \in 2\pi\mathbb{Z}$

Mass Estimate



- Magnetic charge localised at a point
- Divergent energy: $E = \int d^3x \frac{\vec{B}^2}{2} \sim g^2 \Lambda \sim \frac{\Lambda}{e^2}$

QFT of Monopoles



- Full quantum field theory calculation: Monopole loops
- Compare with electron mass correction

$$\delta m = -\frac{e^2}{2\pi^2} m \log \frac{\Lambda}{m} \ll e^2 \Lambda$$

QFT of Monopoles



- Difficult to formulate: Two vector potentials (Schwinger 1975)
- Strong coupling $g = \frac{2\pi}{e} \gg 1$ Non-perturbative!

Georgi-Glashow model: SU(2)+adjoint Higgs

$$\mathcal{L} = -\mathrm{Tr} F^{\mu\nu} F_{\mu\nu} + \mathrm{Tr} [D_{\mu}, \Phi] [D^{\mu}, \Phi] -m^2 \mathrm{Tr} \Phi^2 - \lambda \mathrm{Tr} \Phi^4$$

• $\Phi \neq 0 \Rightarrow$ Symmetry breaking $SU(2) \rightarrow U(1)$

Electrodynamics with magnetic field

$$B_{i} = \frac{1}{2} \epsilon_{ijk} \operatorname{Tr}\widehat{\Phi} \left(F_{jk} - \frac{i}{2e} \left[D_{j}, \widehat{\Phi} \right] \left[D_{k}, \widehat{\Phi} \right] \right)$$

Sourceless ($\vec{\nabla} \cdot \vec{B} = 0$) except when $\Phi = 0$



Smooth "hedgehog" solution:

$$\Phi^a \propto x_a, \quad A_i^a \propto \epsilon_{iaj} x_j$$
• Magnetic charge $g = \int d\vec{S} \cdot \vec{B} = 2g_0 = 4\pi/e$
• Finite mass $M \approx \frac{4\pi\nu}{e} \sim \frac{m}{e^2}$



Montonen-Olive duality conjecture (1977): Electric \leftrightarrow Magnetic $e \leftrightarrow 1/e$



The same solution exists whenever simple Lie group broken to something with a U(1) factor: Grand Unification

Grand Unification

- Standard Model:
 EM & weak forces unified above 100 GeV
- Grand Unified Theory (GUT): Electroweak & strong forces unified above 10¹⁶ GeV

° e.g. $SU(5) \rightarrow SU(3) \times SU(2) \times U(1)$



GUT Monopoles

- Generic prediction of GUTs (and ToEs)
- Mass typically at GUT scale M~10¹⁷ GeV
- Lightest singly-charged $g = g_0 = \frac{2\pi}{e}$
- Catalyse proton decay (Rubakov, Callan 1981)



Simulation of a quantum monopole (AR 2005)

GUT Monopoles

- More complex GUTs, e.g. SO(10)
- Monopoles with different charges
- Can be lighter: $SU(4) \times SU(3) \times SU(3)$ has multiply-charged monopoles with $M \sim 10^7 \, \text{GeV}$ (Kephart et al)



Simulation of a quantum monopole (AR 2005)

String Theory Monopoles

S-duality:

Any superstring theory has magnetic monopoles

► Kaluza-Klein monopole (Gross&Perry, Sorkin): Compactified dimension ↔ U(1) of electrodynamics

Typical mass
$$M \sim \frac{M_{\rm Pl}}{e} \sim 10^{20} {\rm GeV}$$

Cho-Maison Monopole (1996)

- Dirac solution generalised to electroweak theory
- Needs a source particle
- Demonstrates that TeV-scale monopoles are possible
- Twice the Dirac charge

$$g = 2g_0 = \frac{4\pi}{\rho}$$

Mass estimate (Cho&Pinfold 2013) $M \gtrsim \frac{m_W}{e^2} \sim \text{few TeV}$



Probes of New Physics

- Strong, precisely known EM interactions
- ► Electrodynamics ⇒ Stable particles
- Interaction with charged fermions depends on core structure (Rubakov, Callan 1981)
- Properties related to beyond the SM physics:
 - Unification of forces
 - Fundamental properties of electromagnetism

Cosmic Monopoles

- Hot Big Bang: GUT symmetry breaks in a phase transition
- The Higgs field chooses a direction randomly
- ▶ Kibble (1976): Monopoles form, at least one per horizon volume $→ n_{mon} ~ H^{-3}$



Cosmic Monopoles

 Monopoles annihilate until they cannot find partners: Density decreases to

$$n_{\rm mon} \sim 10^{-8} \left(\frac{M}{10^{17} \text{ GeV}} \right) T^3 \sim 10^{-1} \left(\frac{M}{10^{17} \text{ GeV}} \right) \text{m}^{-3}$$

(Zel'dovich & Khlopov 1979, Preskill 1979)

- Mass density higher than observations unless $M \lesssim 10^{10} \text{GeV}$: Monopole problem
- Guth (1981): Inflation wiped monopoles away
 - Monopole production after inflation?

V = 10⁻³ C

Parker Bound (1970)

- Galactic magnetic fields
 B ~ 3µG
- If M ≤ 10¹⁷GeV, this creates a magnetic current, which dissipates the field
- Sets an upper bound on flux

$$F = \frac{nv}{4\pi} \le 10^{-15} \text{ cm}^{-2} \text{ s}^{-1} \text{sr}^{-1}$$

10-10

• Extended Parker bound (Adams et al 1993) $F \lesssim 10^{-16} \left(\frac{M}{10^{17} \text{ GeV}}\right) \text{cm}^{-2} \text{ s}^{-1} \text{sr}^{-1}$



DIRECT SEARCH

Cosmic Rays



(Cabrera 1982)

Early detections:

- Berkeley 1975, Stanford 1982, Imperial 1986
- All turned out to be false

Cosmic Rays



• Upper bound $F \lesssim 10^{-16} \text{ cm}^{-2} \text{ s}^{-1} \text{sr}^{-1}$ over wide mass range

Cosmic Rays



- RICE (South Pole):
 - Intermediate mass monopoles $F \lesssim 10^{-18} \text{ cm}^{-2} \text{ s}^{-1} \text{sr}^{-1}$

Monopoles Trapped in Matter

- Moon rocks (Alvarez et al 1970)
- Deeply buried rocks (Kovalik & Kirschvink 1986)
- Meteorites (Jeon & Longo 1995)
- Polar volcanic rocks (Bendtz et al 2013)
- Bound: Fewer than $\sim 10^{-28}$ per nucleon
- Hard to draw strong conclusions

Monopole-Induced Neutrinos

- Monopoles accumulate in the Sun
- GUT monopoles catalyse p decay: Produces neutrinos (Rubakov 1981)
 Super-K bound (2012):

$$F \lesssim 10^{-23} \left(\frac{\beta_M}{10^{-3}}\right)^2 \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$$

Model-dependent



Accelerator Searches

- Direct searches:
 - OPAL (LEP) 2008: $\sigma < 50 \text{ fb}$ for $M \lesssim 100 \text{ GeV}$
 - CDF (Tevatron) 2006: $\sigma < 200 \text{ fb}$ for $M \leq 700 \text{ GeV}$

• ATLAS (LHC) 2012:



 $\sigma < 16-145~{\rm fb}$ for $M \lesssim 1~{\rm TeV}$

MoEDAL

- Monopole and Exotics Detector at the LHC
- 42 physicists from 11 countries
- Shares LHC intersection point 8 with LHCb
- Will be installed during the long shutdown

monopole

MoEDAL Detectors

- Nuclear track detectors:
 - 10 layers
 - $^\circ\,$ Total area $\sim 10 \times 25\ m^2$
 - Particle tracks revealed by etching
- TimePix radiation monitors:
 - Real time
- Trapping detectors:
 - Analysed with SQUID magnetometer

monopole

MoEDAL



Production Cross Section

- Accelerator experiments can only constrain σ:
 Can we predict it from theory?
- Drell-Yan process:
 q
 r
 q
 q
 m
- Large theoretical uncertainties:
 Perturbation theory not applicable

Exponential Suppression?

- Semiclassical estimate (Drukier&Nussinov 1982):
 - Production cross section $\sim e^{-2/\alpha}$?

- Based on entropy arguments / no of quanta
- Numerical calculation (Demidov&Levkov 2011):
 - Semiclassical Rubakov-Son-Tinyakov method with complex field configurations
 - Confirms exponential suppression for kinks in 1+1D at weak coupling

Beyond Semiclassical

 Lattice calculation of monopole form factor (AR&Weir 2012)

$$\vec{f}(\vec{p}_1, \vec{p}_2) = \left\langle \vec{p}_2 \middle| \vec{B} \middle| \vec{p}_1 \right\rangle$$

- ~ monopole-photon vertex
- Crossing symmetry: Related to production amplitude by analytic continuation
- Always tricky with numerical data

Results for weakish coupling ($\lambda = 0.1$)



Suppression at high k – production rate?

Stronger Coupling?

 1+1D Kinks:
 Suppression disappears at strong coupling

- Monopoles:
 - Two couplings e, λ
 - e fixed
 - Numerical calculation needed
- Singular Cho-Maison monopoles?

f(B)



Summary

- Monopoles are among the best motivated new particles
 - Would open up a window to exciting new physics
- Cosmic monopoles: Stringent bounds
- TeV-scale Cho-Maison monopoles possible
- Theoretical understanding still poor
 - Exponential suppression of production rate?
 - Perturbation theory not applicable
 - Numerical simulations analytic continuation?
- Introductory review: <u>Contemp. Phys. 53 (2012) 195</u>